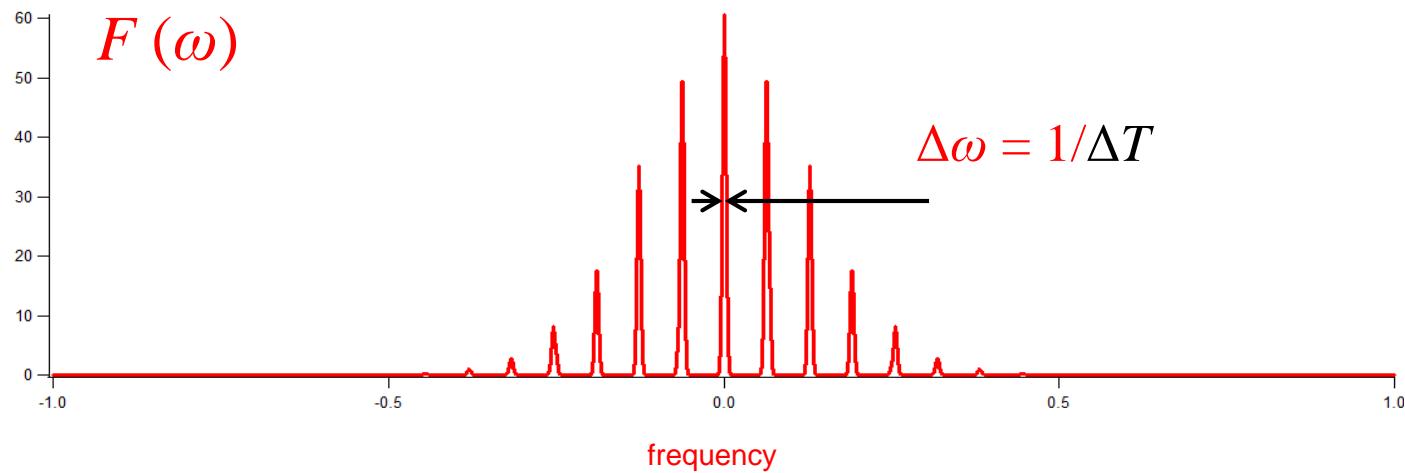
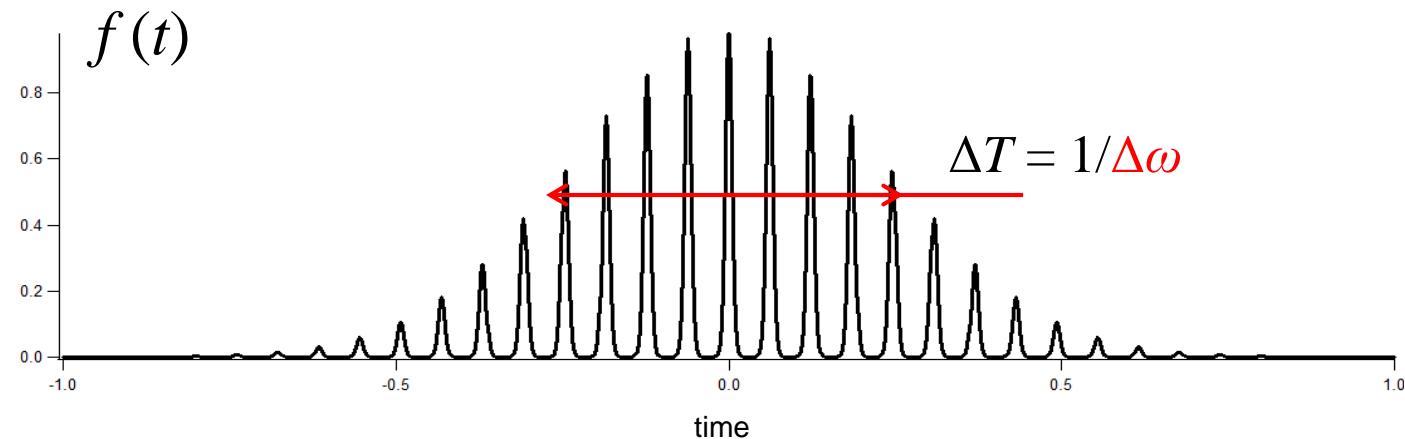
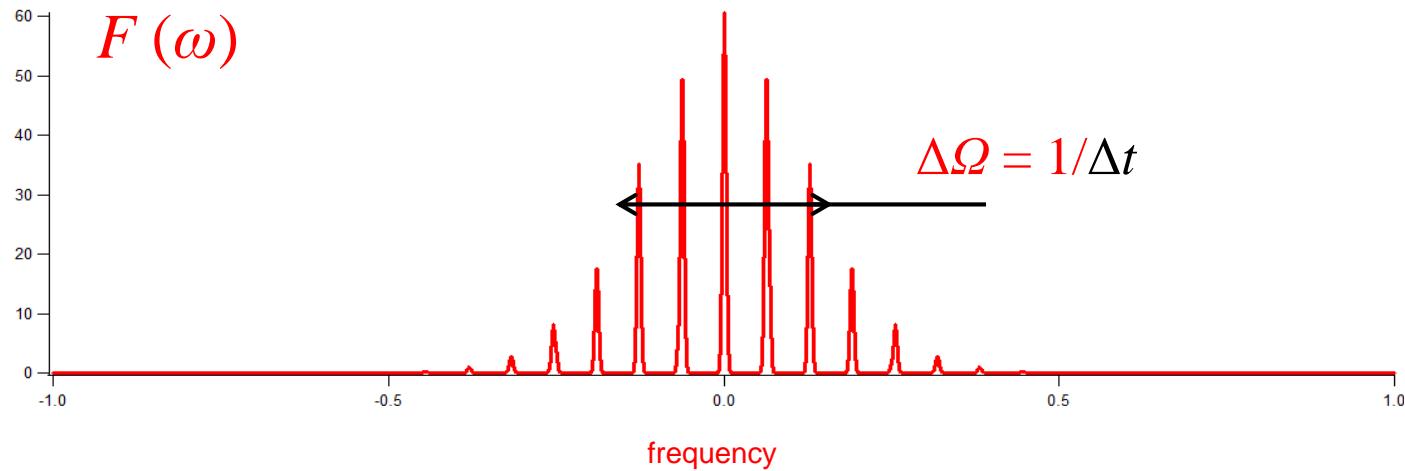
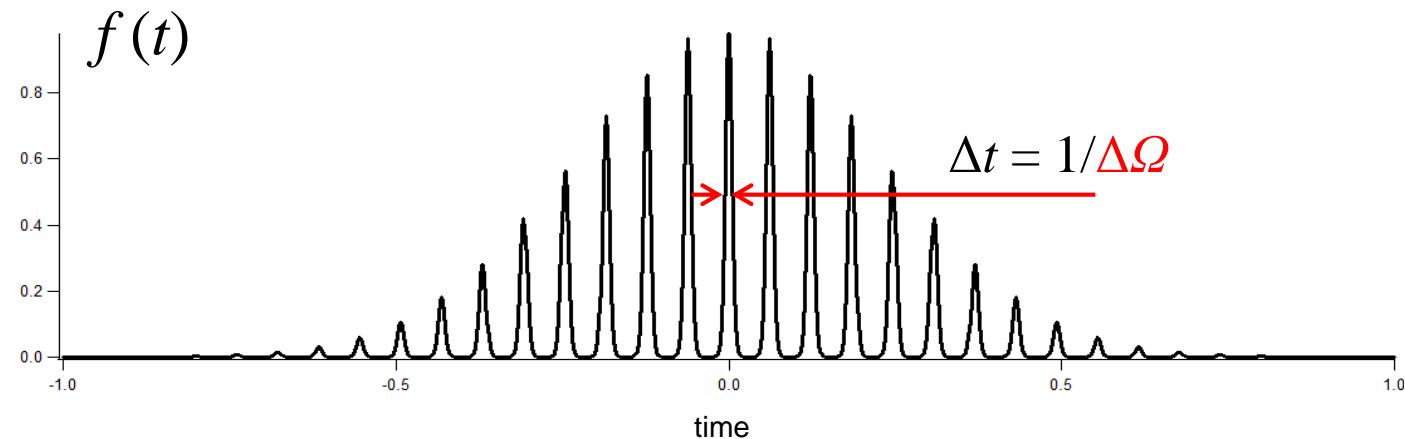
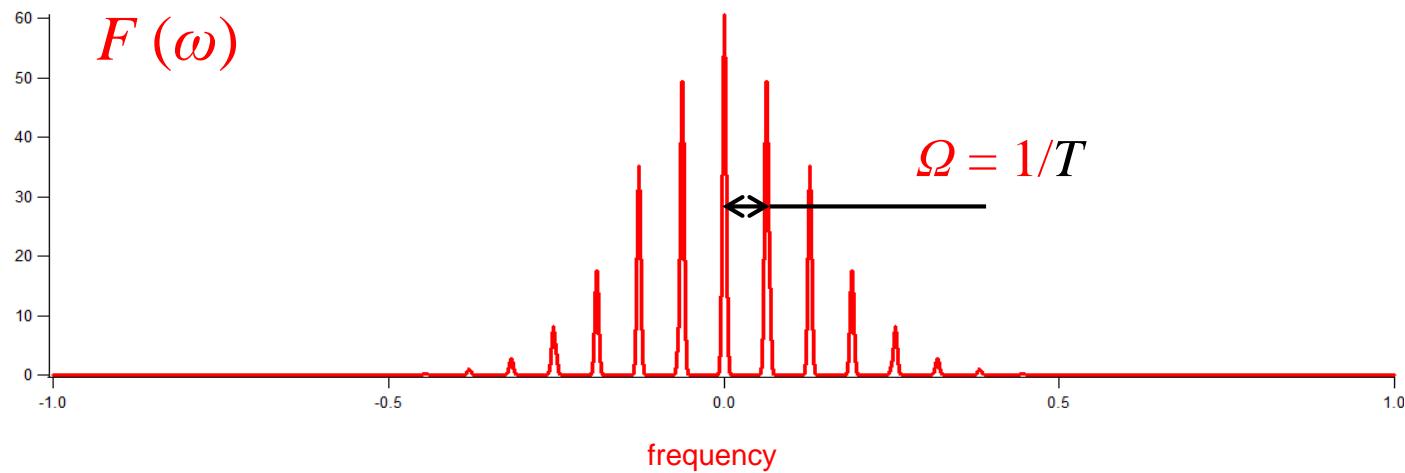
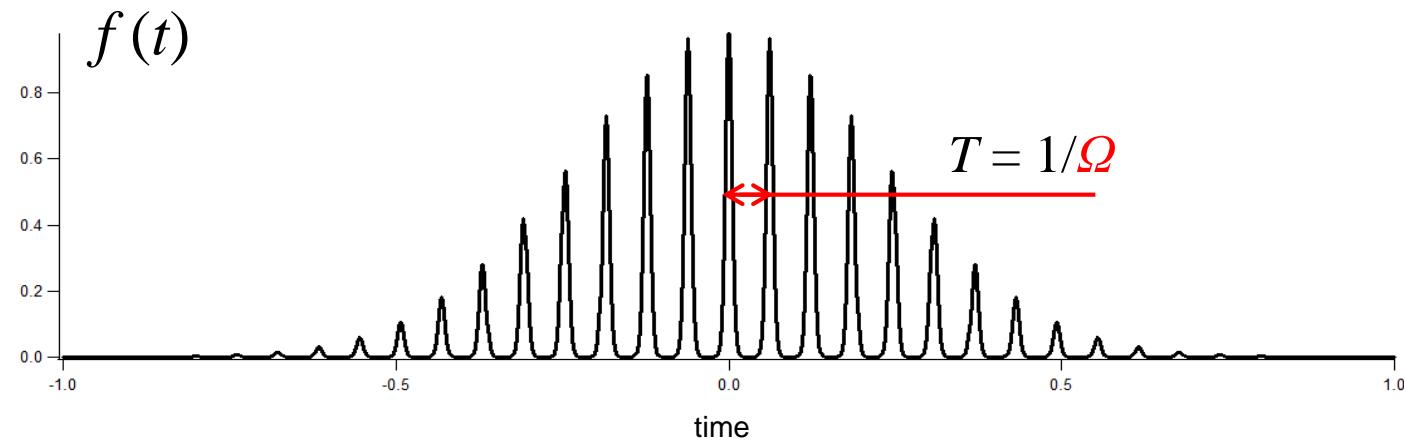


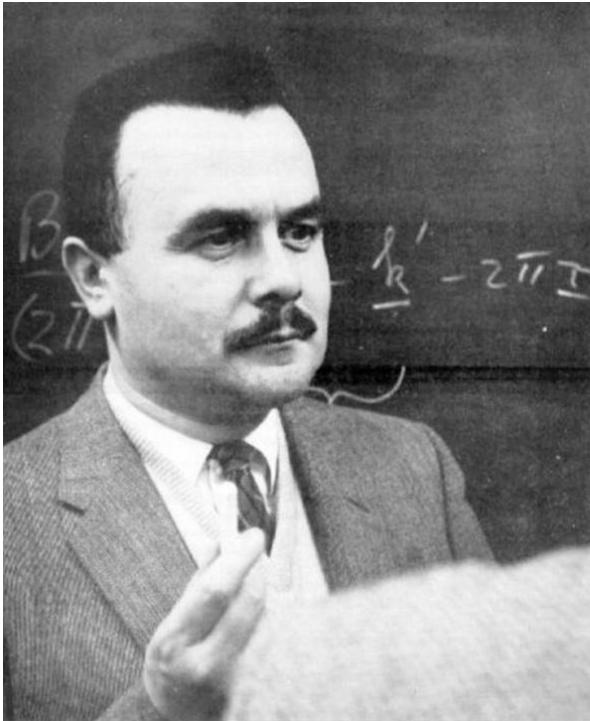
FFT







Inelastic neutron scattering



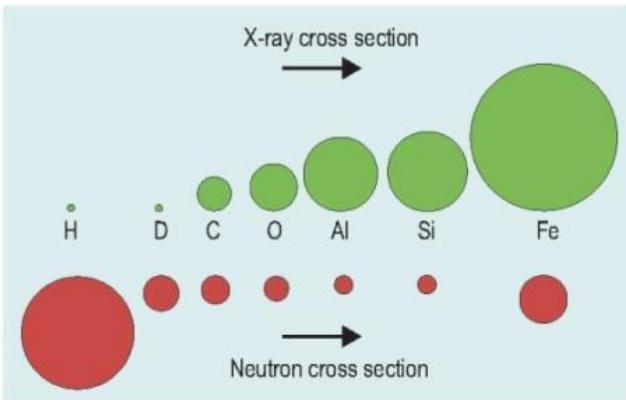
“If the neutron did not exist, it would need to be invented.”

Bertram Brockhouse
1994 Nobel Laureate in Physics

Peter Gehring (NIST)

https://www.ncnr.nist.gov/summerschool/ss12/pdf/SS2012_Gehring_InelasticScattering.pdf

X-rays		EM - wave
Electrons		Charged particle
Neutrons		Neutral particle



Advantages

Wavelengths easily varied to match atomic spacings

Zero charge → not strongly attenuated by furnaces, etc.

Magnetic dipole moment → can study magnetic structures

Nuclear interaction → can see low-Z elements easily like H → good for the study of biomolecules and polymers.

Nuclear interaction is simple → scattering is easy to model Low energies → Non-destructive probe

Disadvantages

Neutrons are expensive to produce → access is not as easy

Interact weakly with matter → often require large samples

Available fluxes are low compared to those for x-rays

Let's consider neutrons ...

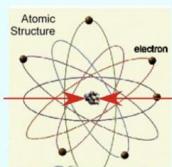
Peter Gehring (NIST)

https://www.ncnr.nist.gov/summerschool/ss12/pdf/SS2012_Gehring_InelasticScattering.pdf

Fast Neutron,

$V \sim 20,000,000 \text{ m/sec}$

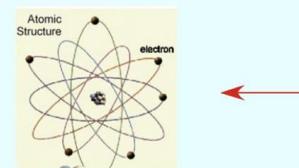
$\sim 0.00002 \text{ nm}$



$$m_n = 1.675 \times 10^{-27} \text{ kg}$$
$$Q = 0$$
$$S = \frac{1}{2} \text{ h}$$
$$\mu_n = -1.913 \mu_N$$

Thermal Neutron,
 $V \sim 2,000 \text{ m/sec}$

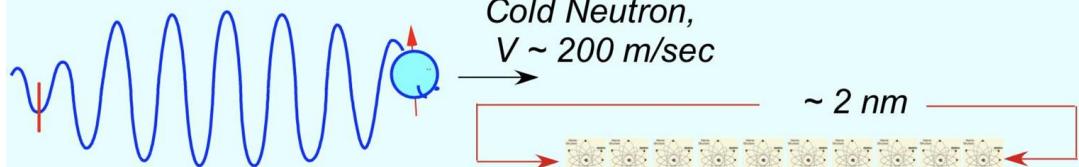
$\sim 0.2 \text{ nm}$



$$\lambda = 1 \text{ \AA}^{\circ}$$
$$v = 4000 \text{ m/s}$$
$$E = 82 \text{ meV}$$

Cold Neutron,
 $V \sim 200 \text{ m/sec}$

$\sim 2 \text{ nm}$



$$\lambda = 9 \text{ \AA}^{\circ}$$
$$v = 440 \text{ m/s}$$
$$E = 1 \text{ meV}$$

Peter Gehring (NIST)

https://www.ncnr.nist.gov/summerschool/ss12/pdf/SS2012_Gehring_InelasticScattering.pdf

Neutron Scattering Cross Sections

What are the physical meanings
of these three cross sections?

σ Total # of neutrons scattered per second / Φ_i .

$\frac{d\sigma}{d\Omega}$ Total # of neutrons scattered per second into $d\Omega$ / $d\Omega \Phi_i$.
(Diffraction → structure.

$\frac{d^2\sigma}{d\Omega dE_f}$ Total # of neutrons scattered per second into $d\Omega$
with a final energy between E_f and dE_f / $d\Omega dE_f \Phi_i$.
(Inelastic scattering → dynamics.

Peter Gehring (NIST)

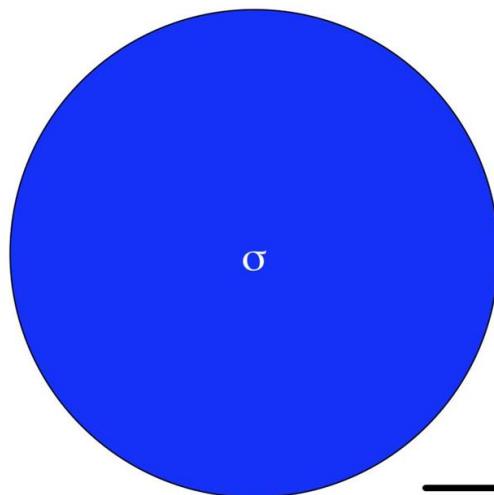
https://www.ncnr.nist.gov/summerschool/ss12/pdf/SS2012_Gehring_InelasticScattering.pdf

Neutron Scattering Cross Sections

What are the relative
sizes of the cross sections?

Clearly: $\sigma = \int \frac{d\sigma}{d\Omega} d\Omega = \int \frac{d^2\sigma}{d\Omega dE_f} d\Omega dE_f$

Thus: $\sigma >> \frac{d\sigma}{d\Omega} >> \frac{d^2\sigma}{d\Omega dE_f}$



$$\frac{d\sigma}{d\Omega}$$

$$\cdot \xleftarrow{\frac{d^2\sigma}{d\Omega dE_f}}$$

Typically, $\frac{d\sigma}{d\Omega} \sim 10^6 \times \frac{d^2\sigma}{d\Omega dE_f}$

Peter Gehring (NIST)

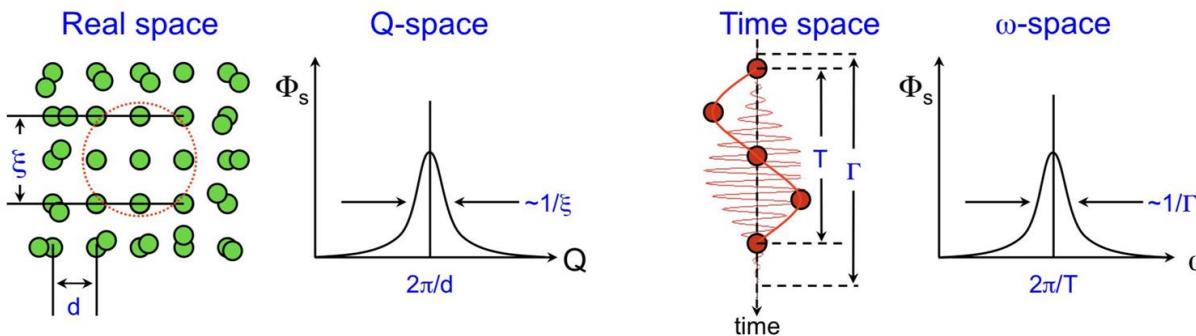
https://www.ncnr.nist.gov/summerschool/ss12/pdf/SS2012_Gehring_InelasticScattering.pdf

Pair Correlation Functions

KEY IDEA – Neutron interactions are weak →
Scattering only probes two-particle correlations
in space and time, but does so simultaneously!

The scattered neutron flux $\Phi_s(\vec{Q}, \hbar\omega)$ is proportional to the space (\vec{r}) and time (t) Fourier transform of the probability $G(\vec{r}, t)$ of finding an atom at (\vec{r}, t) given that there is another atom at $r = 0$ at time $t = 0$.

$$\Phi_s \propto \frac{\partial^2 \sigma}{\partial \Omega \partial \omega} \propto \iint e^{i(\vec{Q} \cdot \vec{r} - \omega t)} G(\vec{r}, t) d^3 \vec{r} dt$$



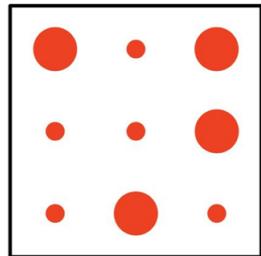
Peter Gehring (NIST)

https://www.ncnr.nist.gov/summerschool/ss12/pdf/SS2012_Gehring_InelasticScattering.pdf

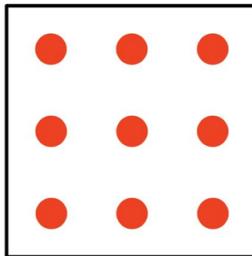
Neutron Coherent and Incoherent Scattering

Consider a system composed of two different scattering lengths, b_1 and b_2 .

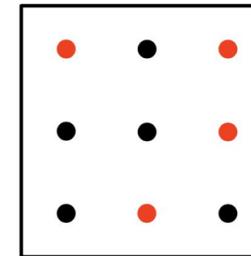
$$b_1 = \bullet$$
$$b_2 = \textcolor{red}{\bullet}$$



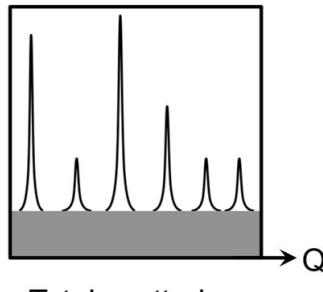
=



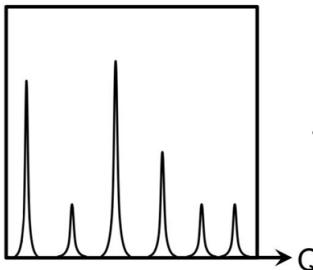
+



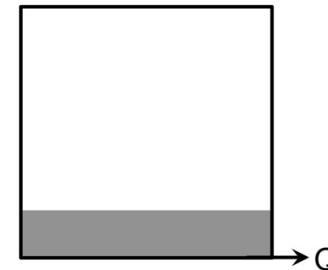
The two isotopes are randomly distributed.



=



+



Total scattering

Coherent scattering

Incoherent scattering

$$\frac{1}{2}(b_1 + b_2) = \bar{b}$$

Deviations δb

Peter Gehring (NIST)

https://www.ncnr.nist.gov/summerschool/ss12/pdf/SS2012_Gehring_InelasticScattering.pdf

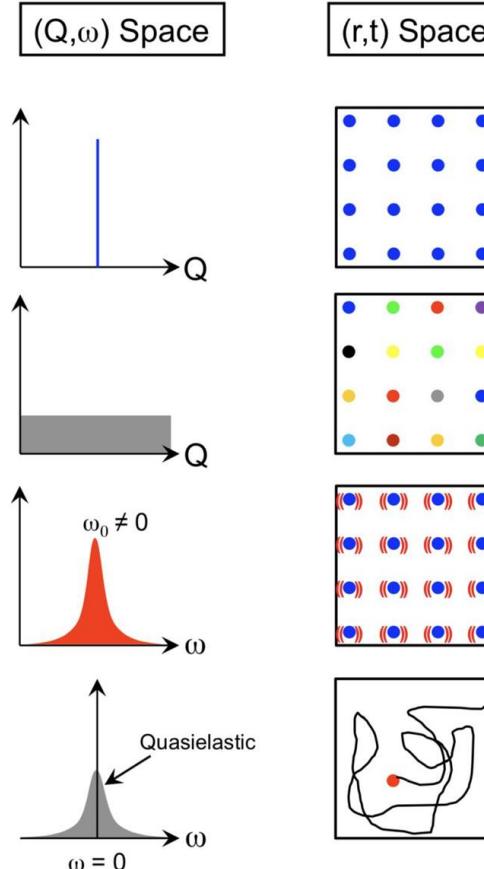
Summary of Cross Sections

$$\frac{d\sigma}{d\Omega} \Big|_{coh} = \frac{\sigma_{coh}}{4\pi} S(Q)$$

$$\frac{d\sigma}{d\Omega} \Big|_{inc} = \frac{\sigma_{inc}}{4\pi}$$

$$\frac{d^2\sigma}{d\Omega dE_f} \Big|_{coh} = \frac{k_f}{k_i} \frac{\sigma_{coh}}{4\pi} S_{coh}(Q, \omega)$$

$$\frac{d^2\sigma}{d\Omega dE_f} \Big|_{inc} = \frac{k_f}{k_i} \frac{\sigma_{inc}}{4\pi} S_{inc}(Q, \omega)$$



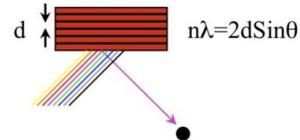
Peter Gehring (NIST)

https://www.ncnr.nist.gov/summerschool/ss12/pdf/SS2012_Gehring_InelasticScattering.pdf

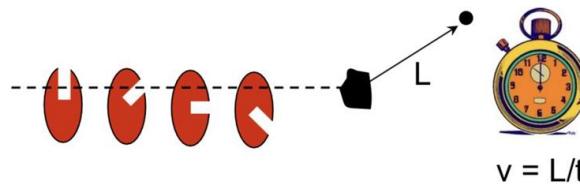
Basics of Neutron Scattering Methods

Methods of specifying and measuring \vec{k}_i and \vec{k}_f

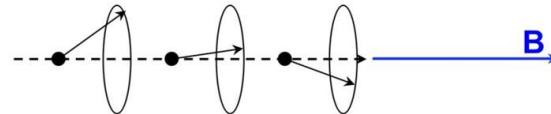
1. Bragg Diffraction



2. Time-of-Flight (TOF)

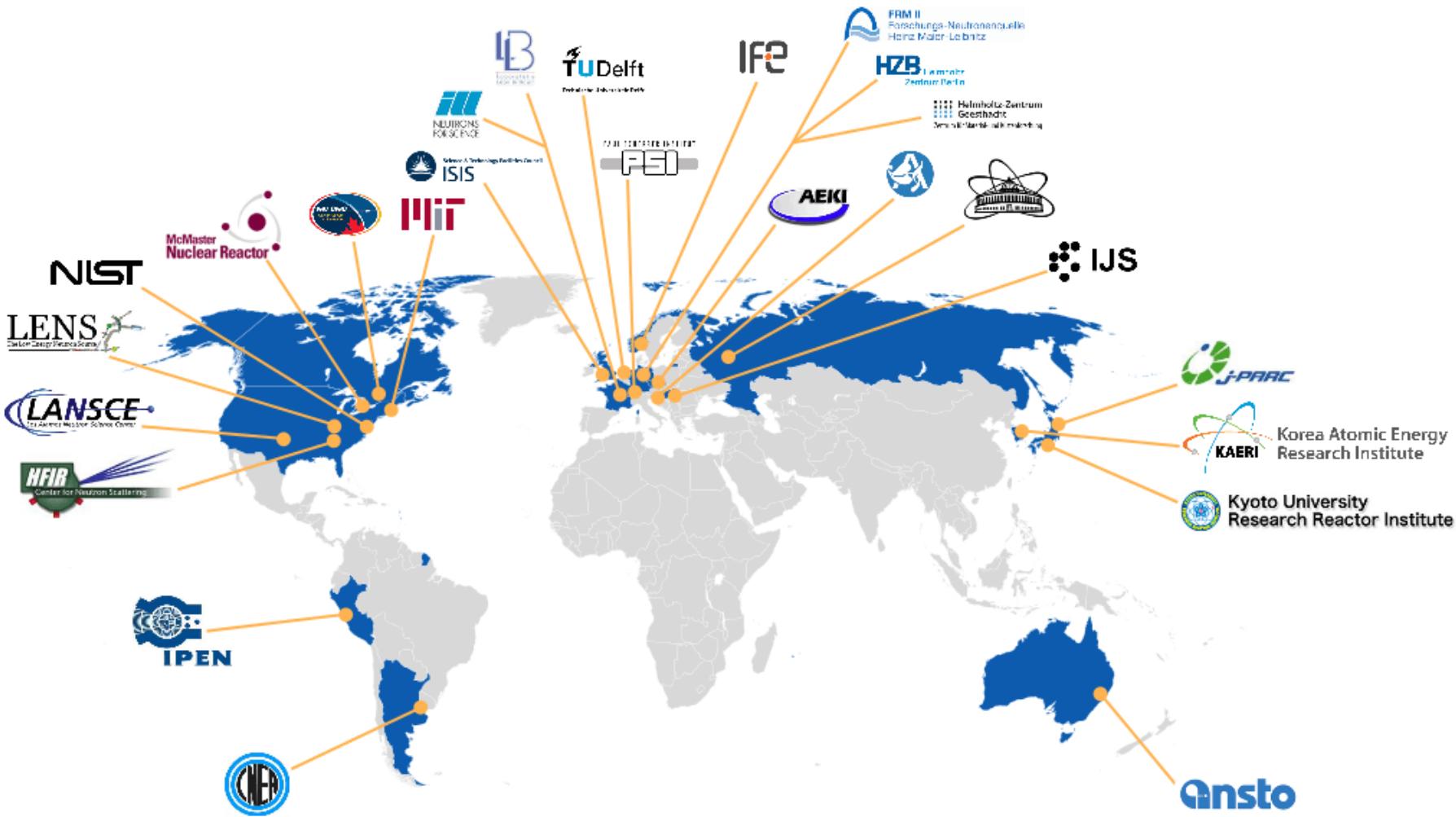


3. Larmor Precession

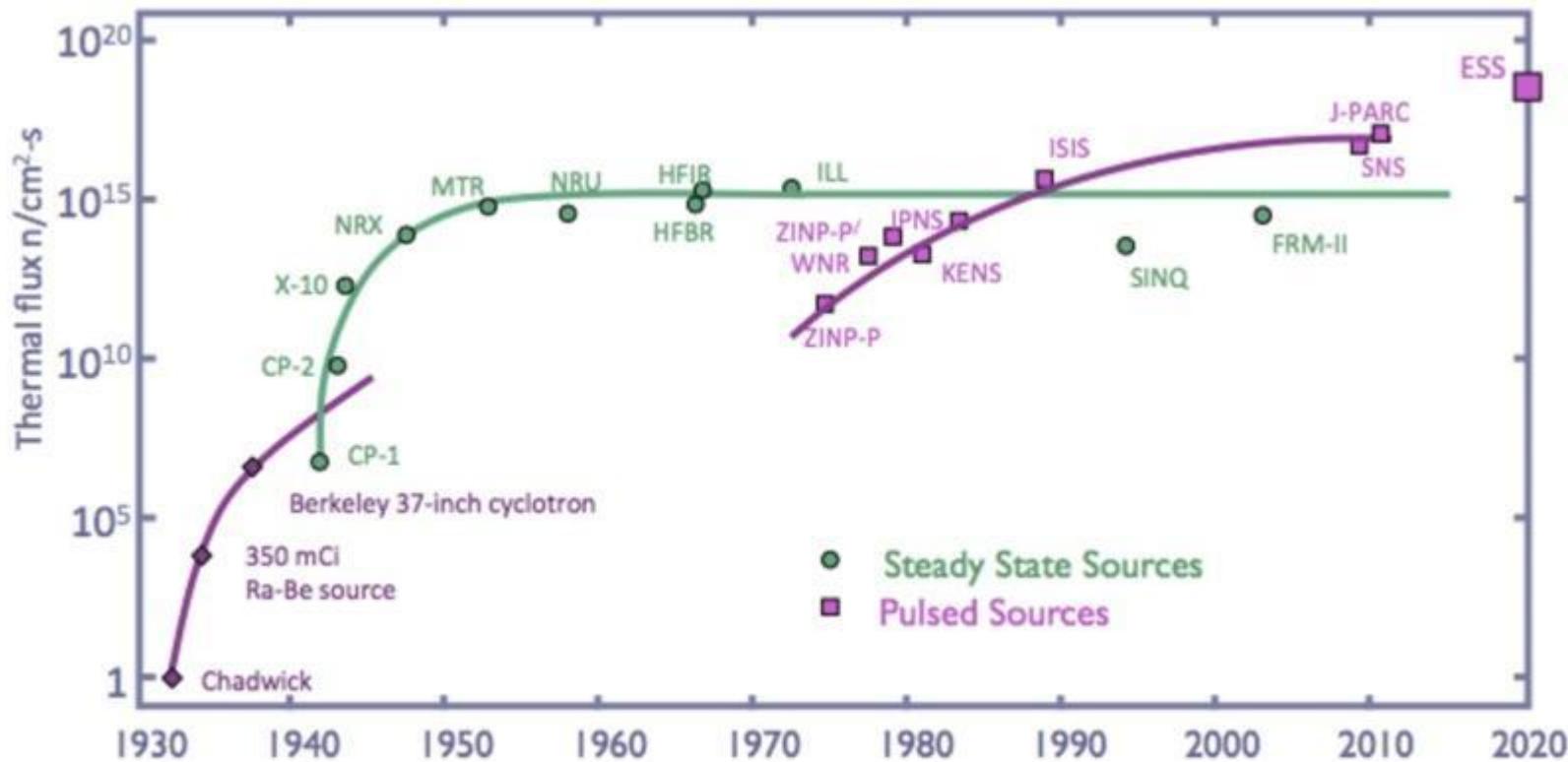


Peter Gehring (NIST)

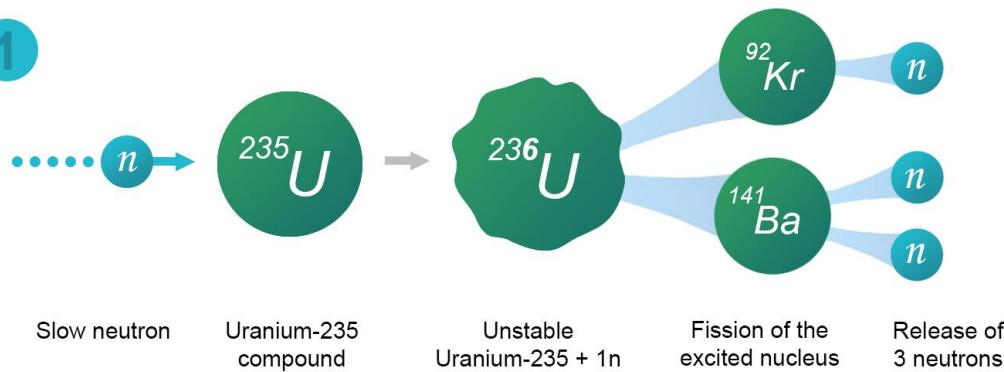
https://www.ncnr.nist.gov/summerschool/ss12/pdf/SS2012_Gehring_InelasticScattering.pdf



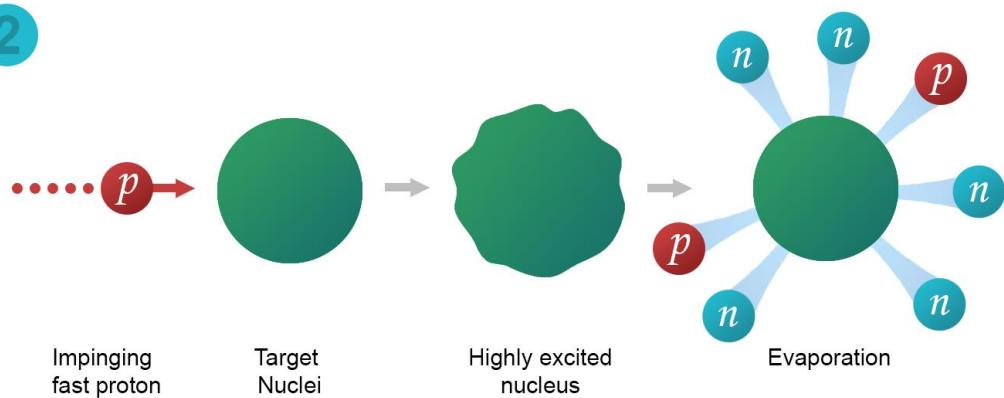
Evolution of neutron sources



1



2



- **Fission:** A high continuous flux of neutrons is produced in the core of a conventional fission reactor.

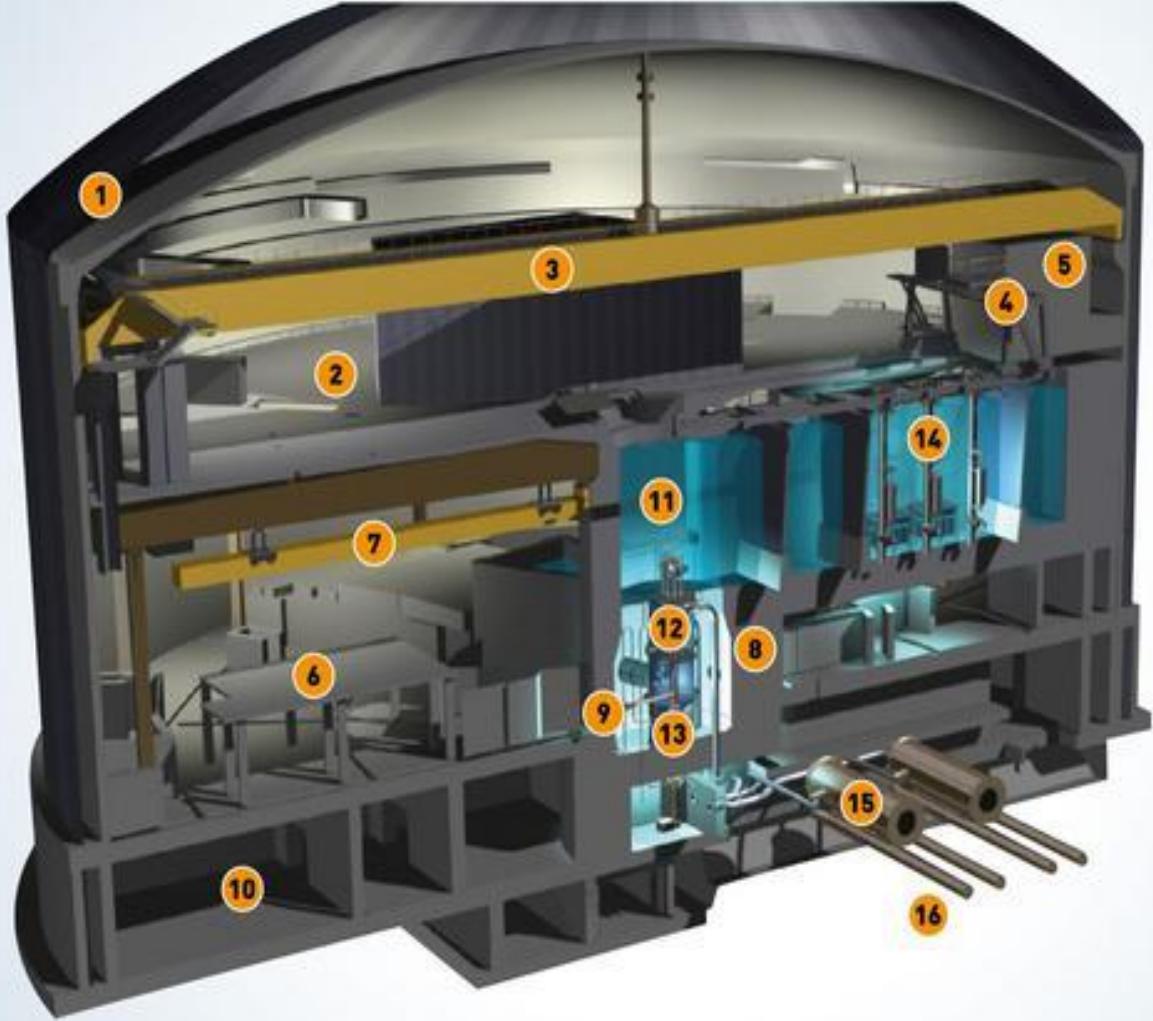
- **Spallation:** A pulsed production of neutrons is obtained by bombarding a target of heavy elements with high-energy particles, typically accelerated protons.



Institut Laue-Langevin

Institut Laue-Langevin <https://youtu.be/xrQgURmcAlc>





DOUBLE-WALLED REACTOR BUILDING

1

LEVEL D - REACTOR HALL

2

CRANE FOR REACTOR OPERATIONS LEVEL D

3

GANTRY FOR HANDLING OF FUEL ELEMENTS

4

HOT CELL

5

LEVEL C - EXPERIMENTAL HALL

6

CRANE FOR EXPERIMENTAL OPERATIONS

7

BIOLOGICAL SHIELDING (CONCRETE)

8

COLLIMATED NEUTRON EXIT POINT

9

LEVEL B - REACTOR AUXILIARY EQUIPMENT

10

REACTOR POOL (LIGHT WATER)

11

HEAVY WATER (MODERATOR & FUEL ELEMENT COOLING)

12

FUEL ELEMENT

13

SPENT FUEL ELEMENTS STORAGE

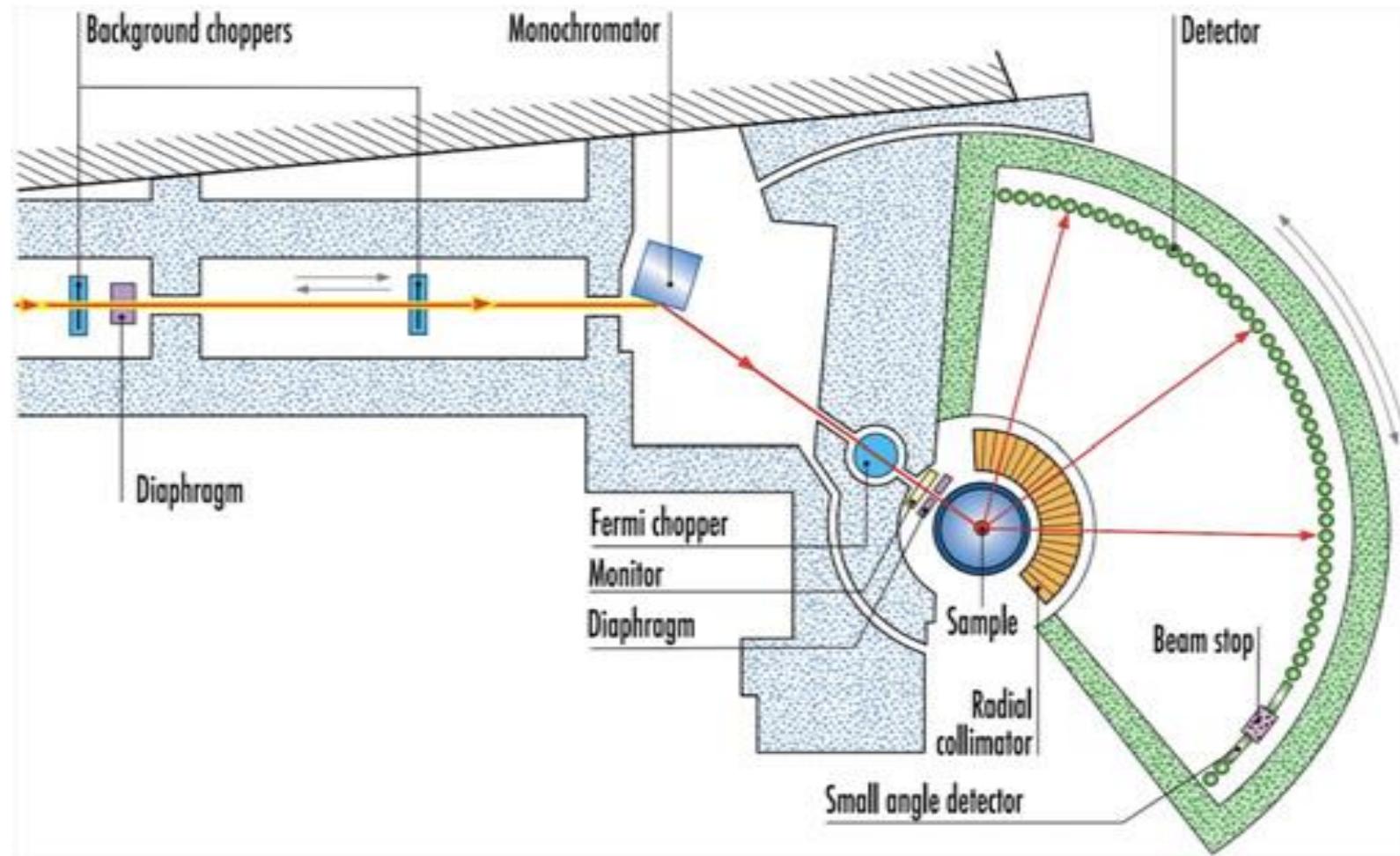
14

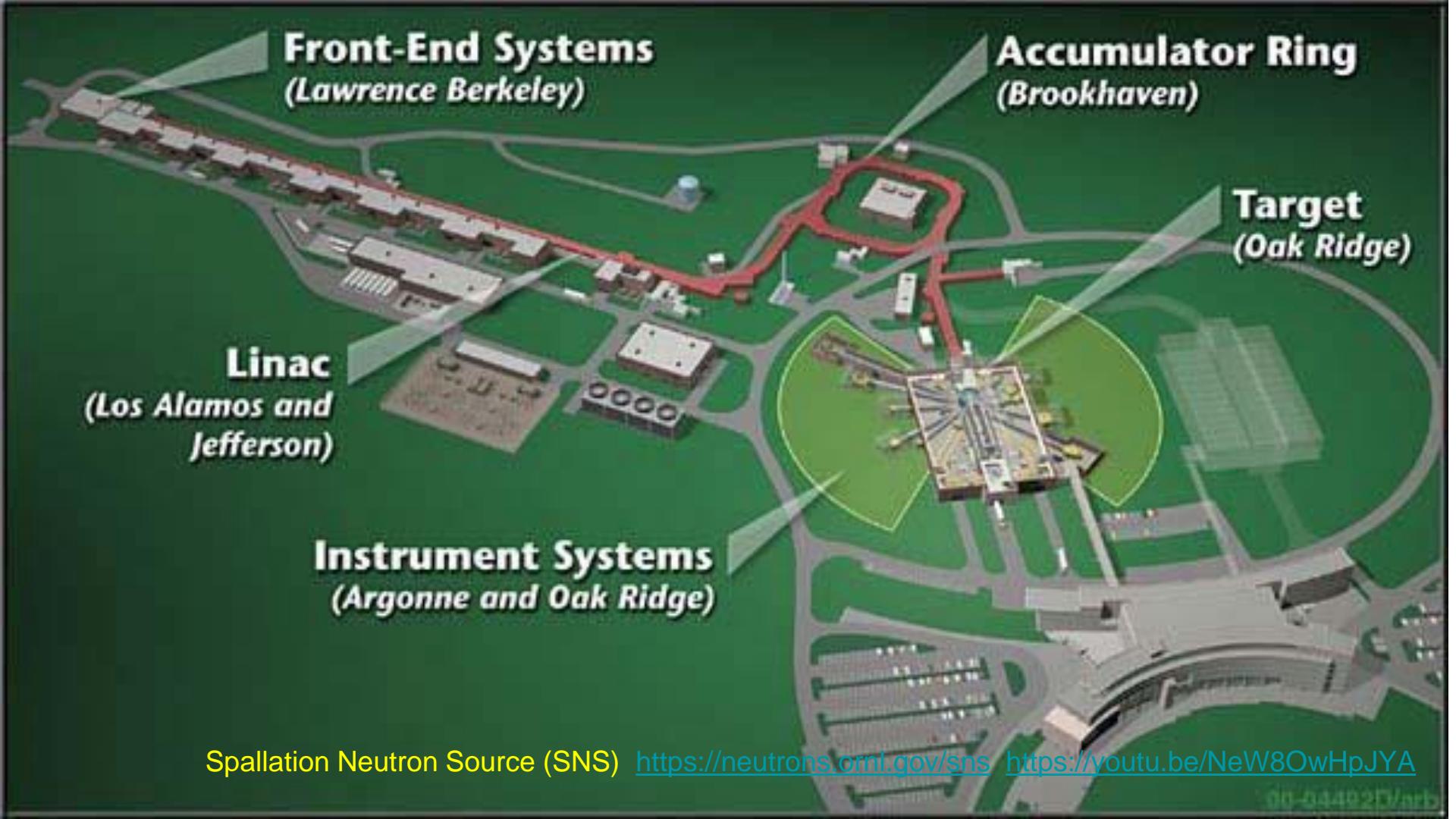
HEAT EXCHANGERS (PRIMARY/SECONDARY)

15

SECONDARY COOLING CIRCUIT 'DRAC RIVER'

16





Front-End Systems
(Lawrence Berkeley)

Accumulator Ring
(Brookhaven)

Linac
(Los Alamos and Jefferson)

Instrument Systems
(Argonne and Oak Ridge)

Target
(Oak Ridge)

Spallation Neutron Source (SNS) <https://neutrons.ornl.gov/sns> <https://youtu.be/NeW8OwHpJYA>



Oak Ridge,
Tennessee



2006

LOCATION
START OF OPERATIONS



759 (FY 2019)

NUMBER OF USERS



Spallation Neutron Source (SNS)

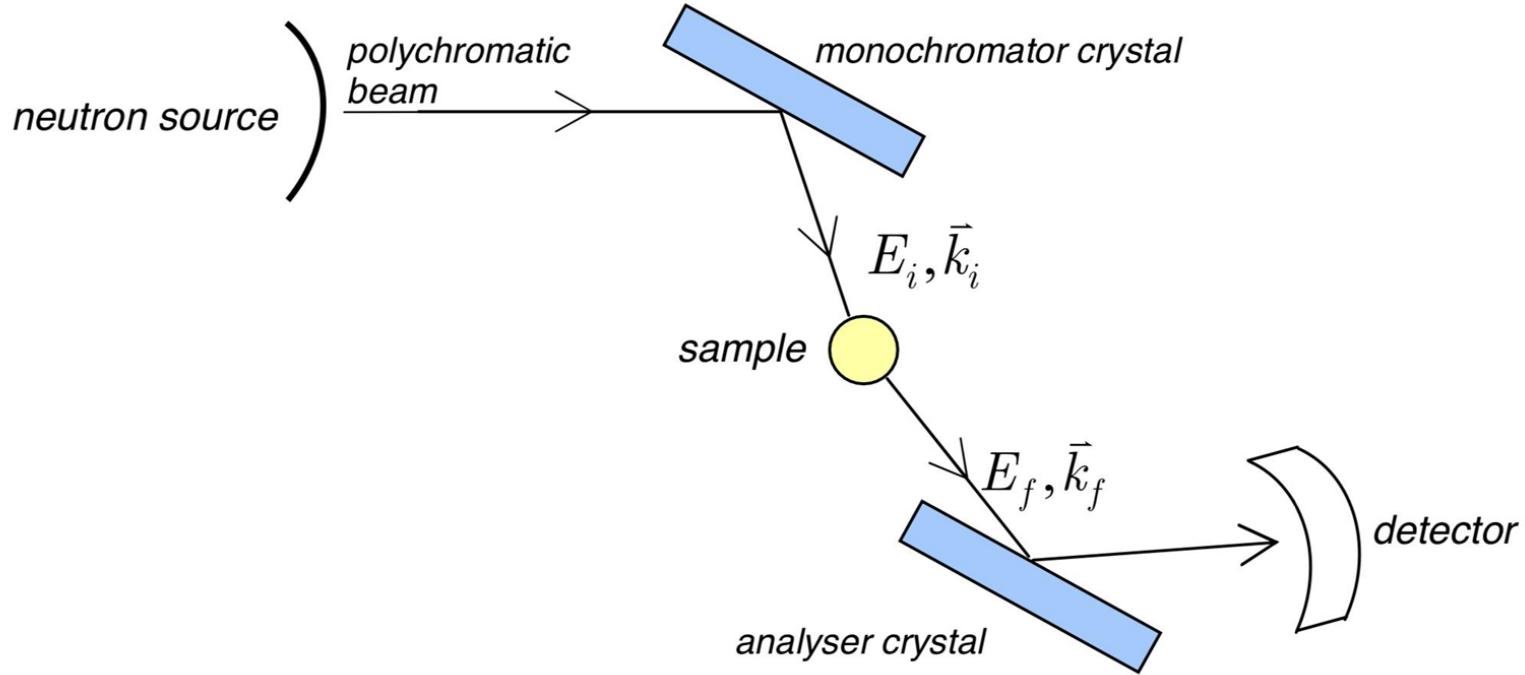


EUROPEAN
SPALLATION
SOURCE

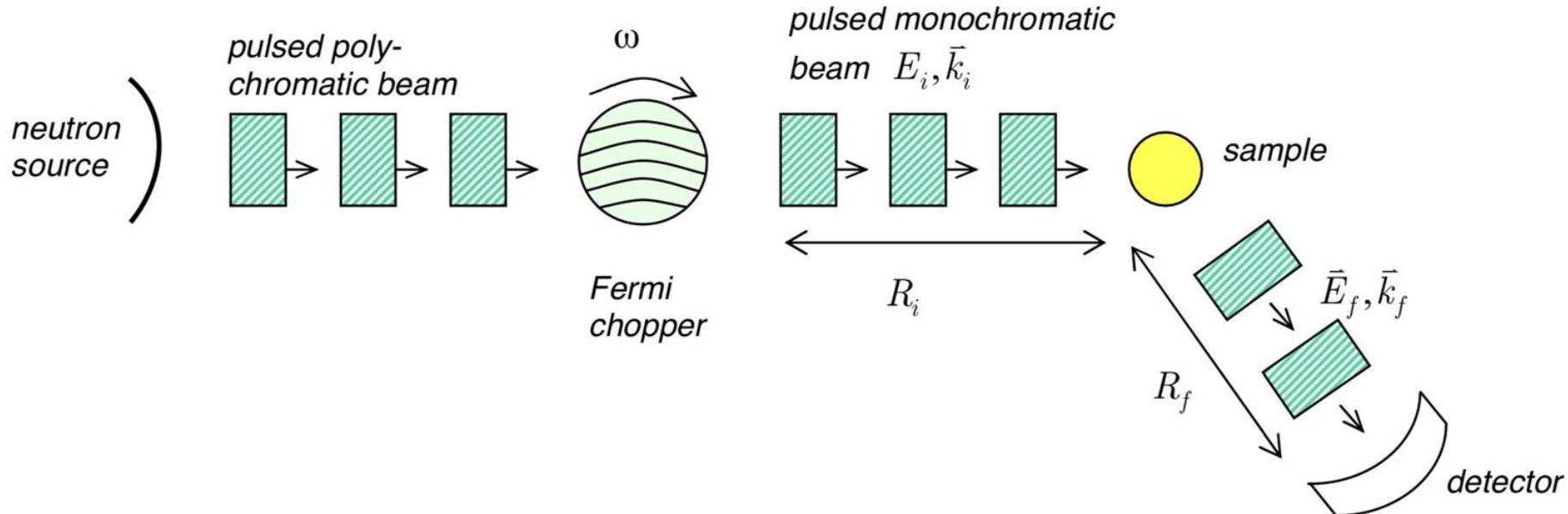
<https://europeanspallationsource.se/webcams>

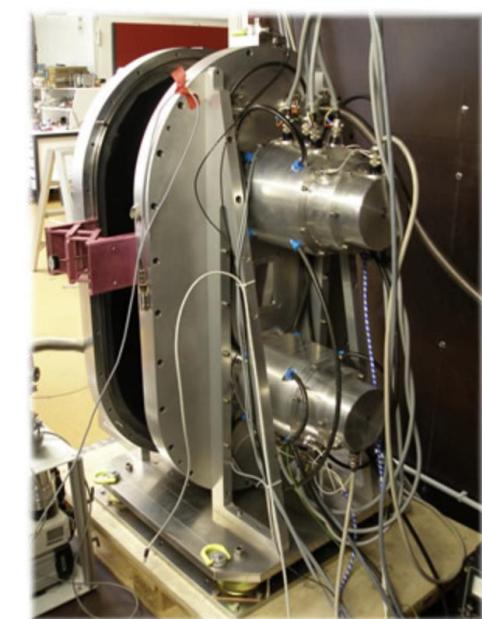
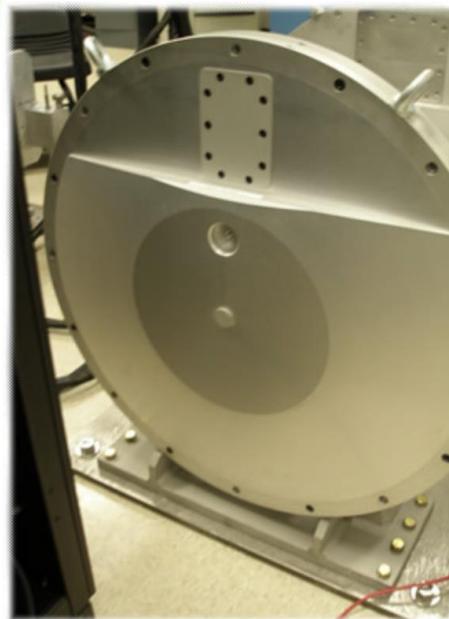
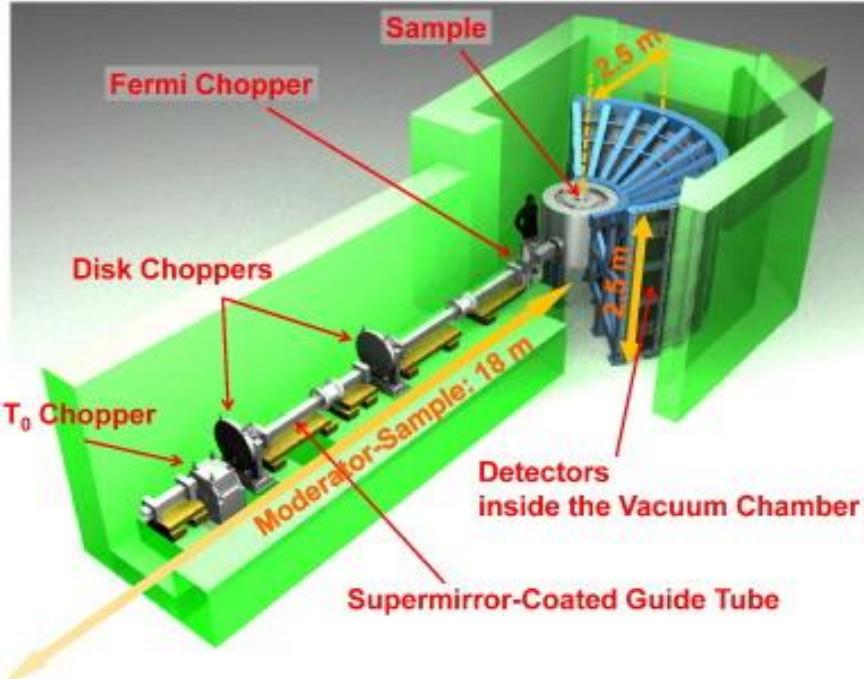


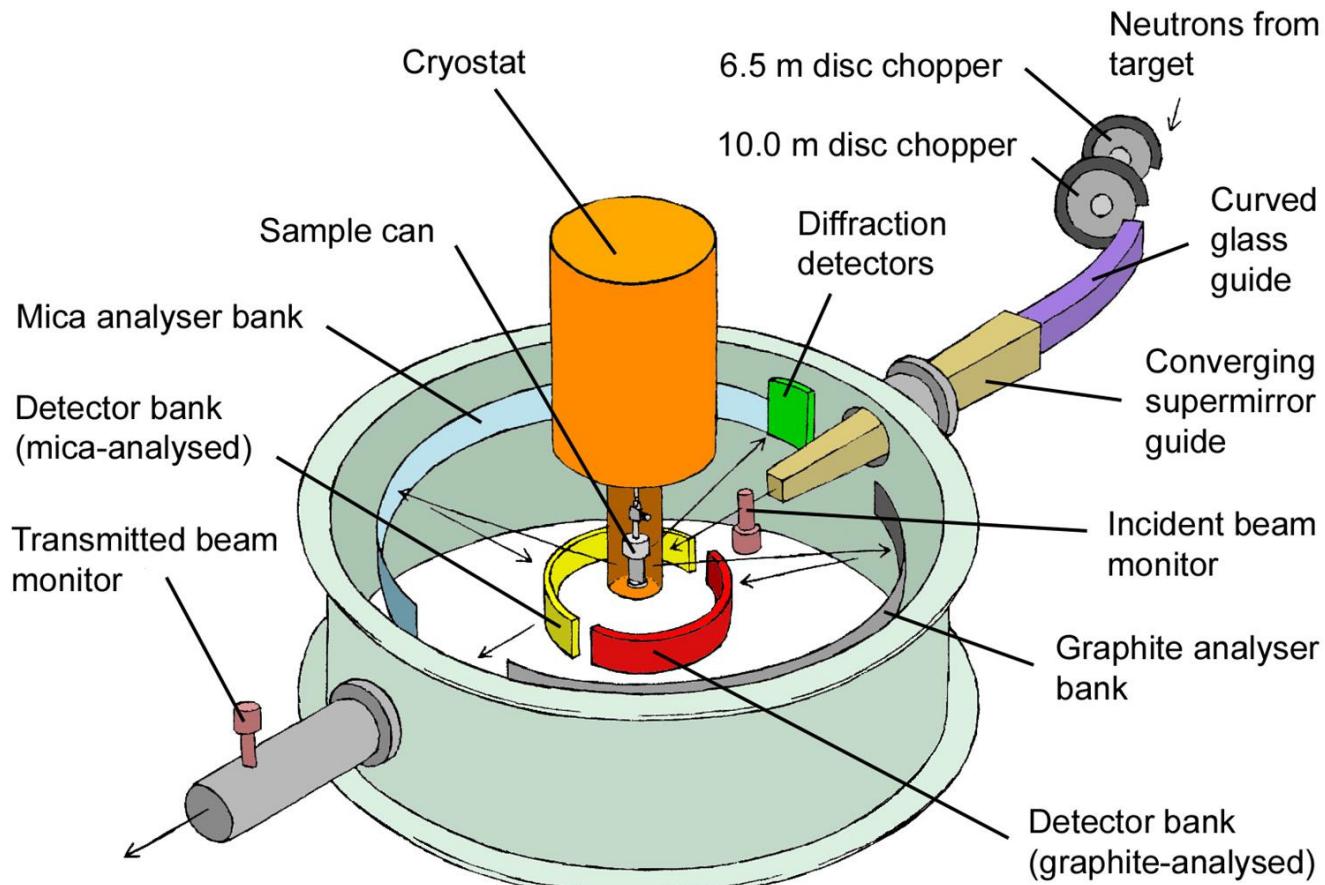
1 – Triple-axis spectrometer (commonly used at steady-state neutron sources)



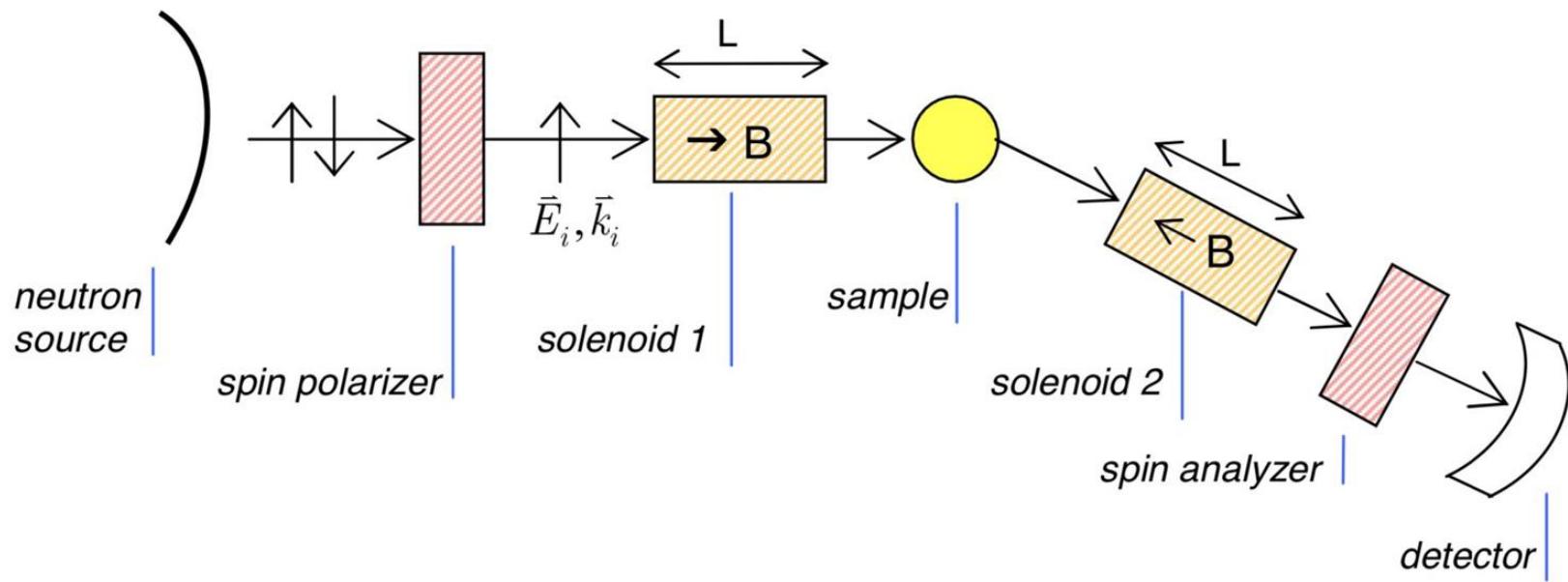
2 – Time-of-flight spectrometer (commonly used at pulsed neutron sources)



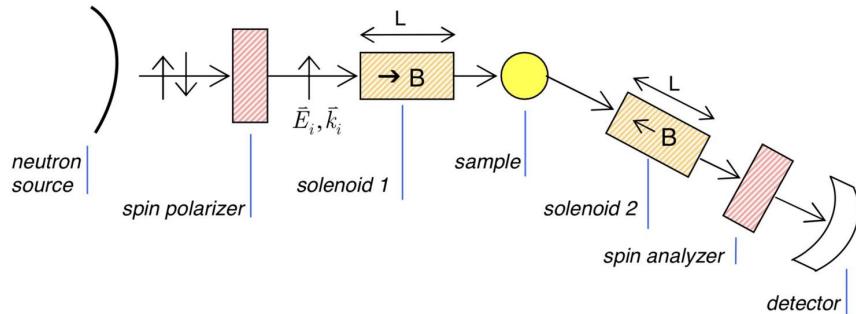




3 – Spin-echo spectrometer



3 – Spin-echo spectrometer

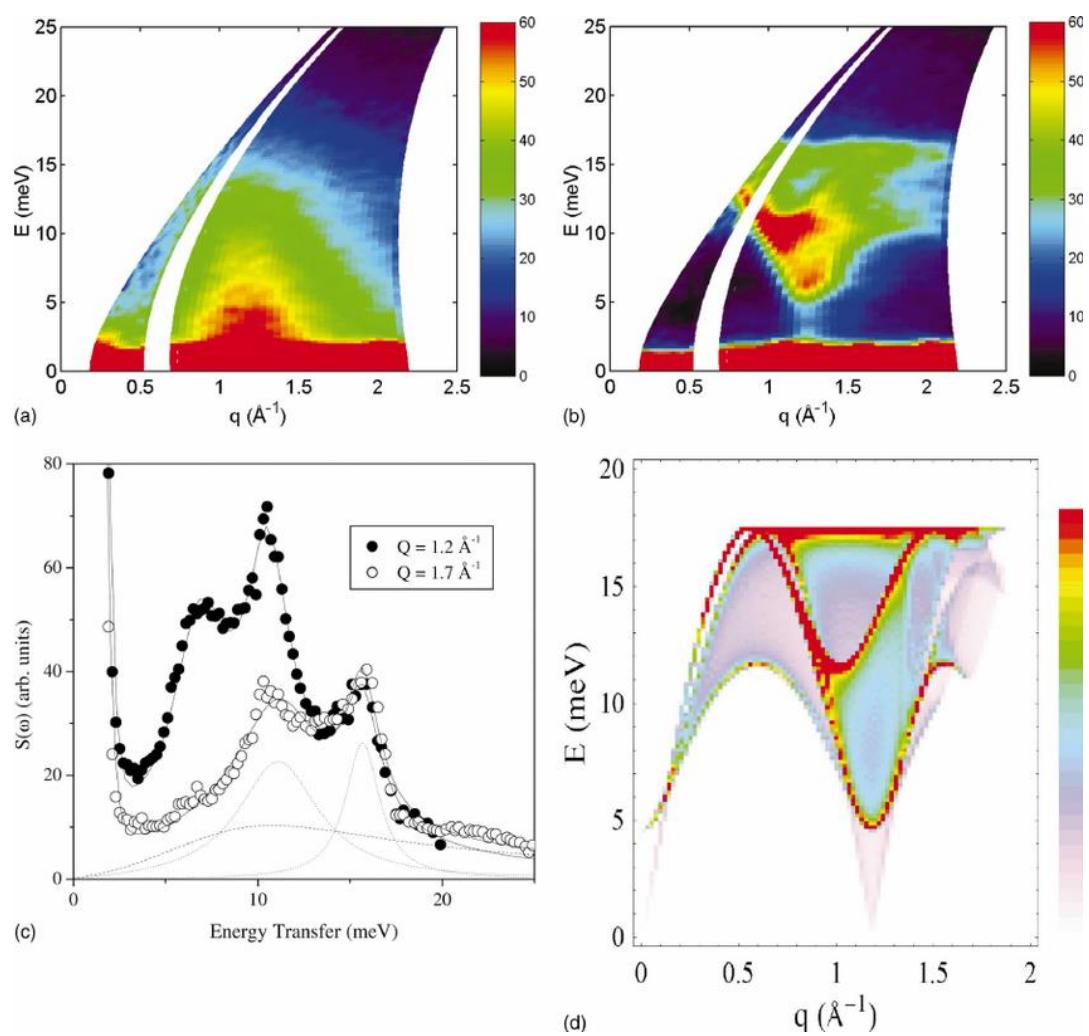


$$\omega_L = \frac{\gamma \mu_N B}{\hbar} \quad \phi_i = \frac{\omega_L L}{v_i}$$

$$\Delta\Phi = \Phi_i - \Phi_f = \omega_L L \left(\frac{1}{v_i} - \frac{1}{v_f} \right) = \omega_L L \left(\frac{1}{v_i} - \frac{1}{v_i + \Delta v} \right) \approx \frac{\omega_L L}{v_i^2} \Delta v \text{ for } \Delta v \ll v_i$$

$$\hbar\omega = \frac{m_n}{2} \left(v_f^2 - v_i^2 \right) \approx m_n v_i \Delta v \Rightarrow \Delta\Phi \approx \frac{\hbar\omega_L L}{m_n v_i^3} \omega \equiv \omega \tau_{SE}$$

τ_{SE} = spin echo time



Magnetic ordering and spin-liquid state of YMnO₃ - Park *PRB* 2003