

Мюонна спінова спектроскопія

Muon spin spectroscopy

**$\mu$ SR**

Muon Spin Rotation/Relaxation

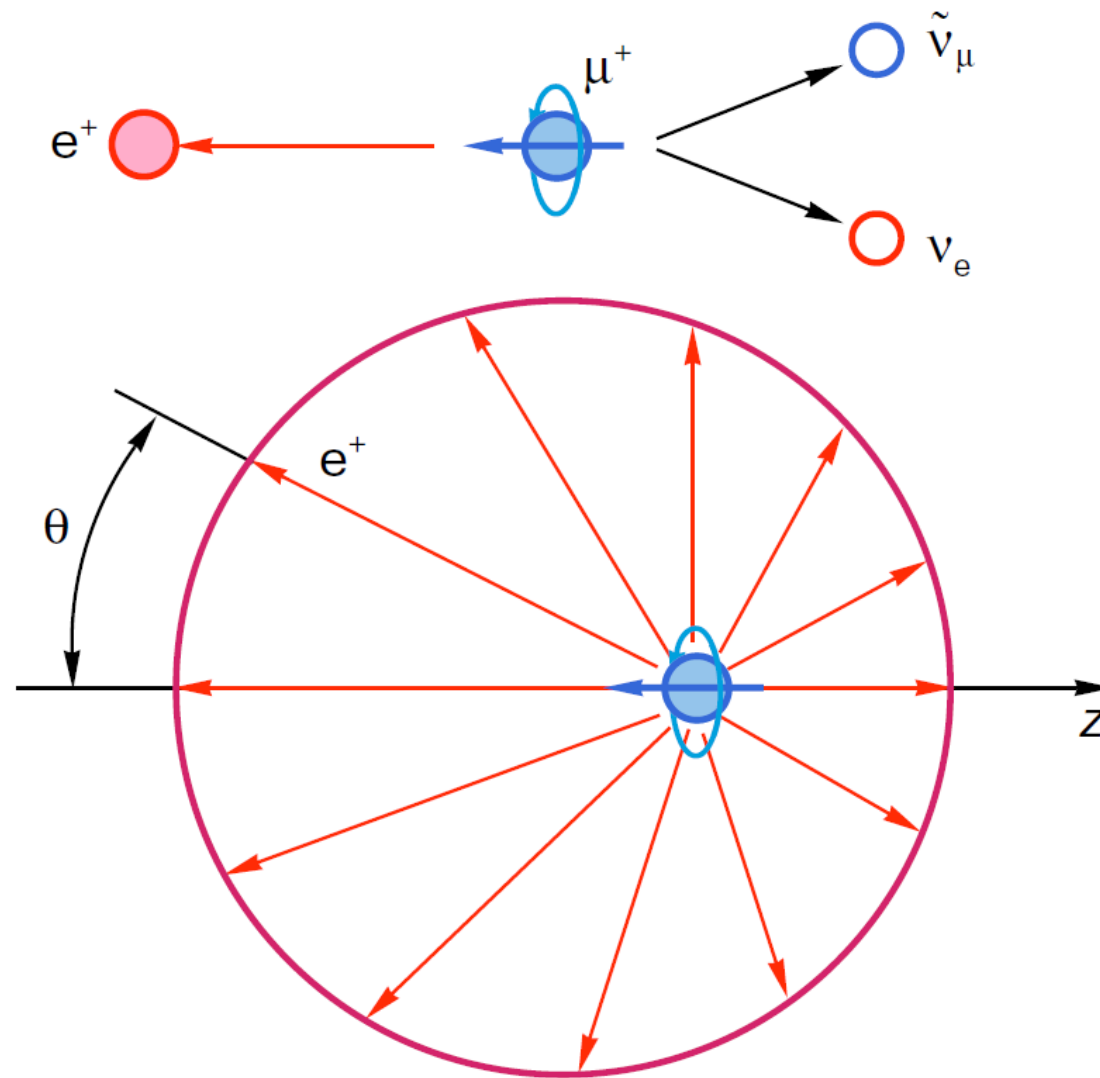
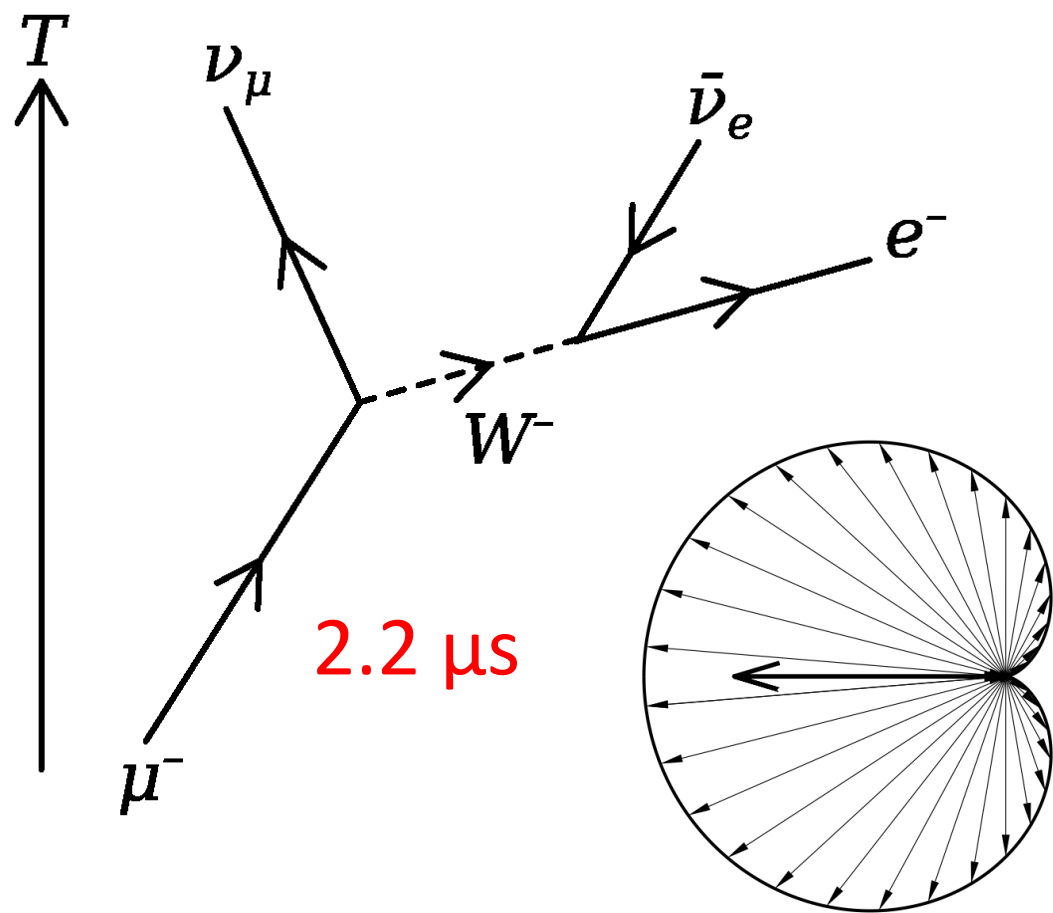
## *To read*

1. S. J. Blundell (1999) Spin-polarized muons in condensed matter physics, *Contemporary Physics* **40**, 175-192 <https://arxiv.org/abs/cond-mat/0207699v1>
2. L. Nuccio *et al.* Muon spin spectroscopy: magnetism, soft matter and the bridge between the two. *J. Phys. D: Appl. Phys.* **47**, 473001 (2014) [pdf](#)
3. Белоусов Ю.М., Смилга В.П. Что такое мюонный метод исследования вещества. *Соросовский образовательный журнал* №1, с. 76-85 (1999) [pdf](#)
4. E. Morenzoni. Muon science with continuous beams at PSI (2014) <https://www.isis.stfc.ac.uk/Pages/2014-morenzoni-psi.pdf>
5. J. E. Sonier *et al.* mSR studies of the vortex state in type-II superconductors. *Reviews of Modern Physics* **72**, 769 (2000) [pdf](#)

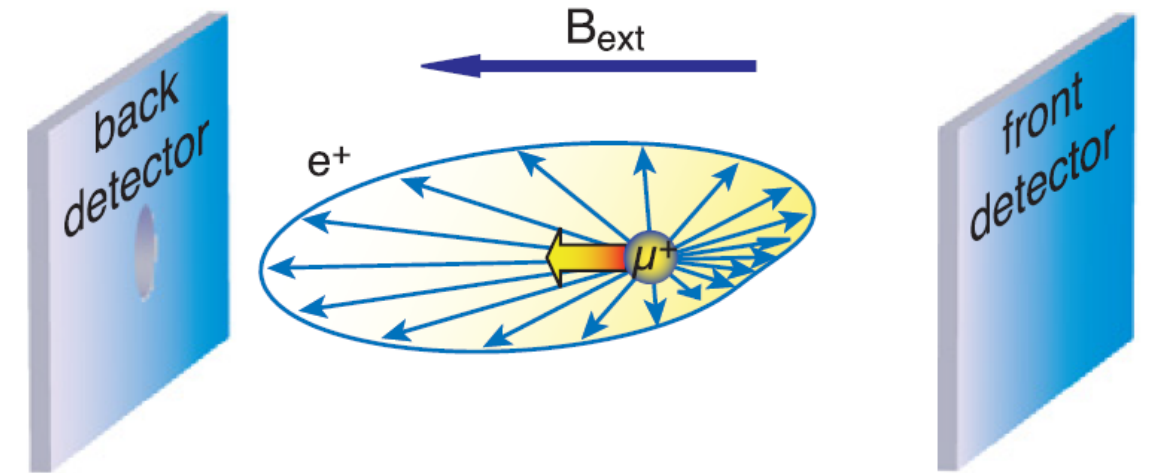
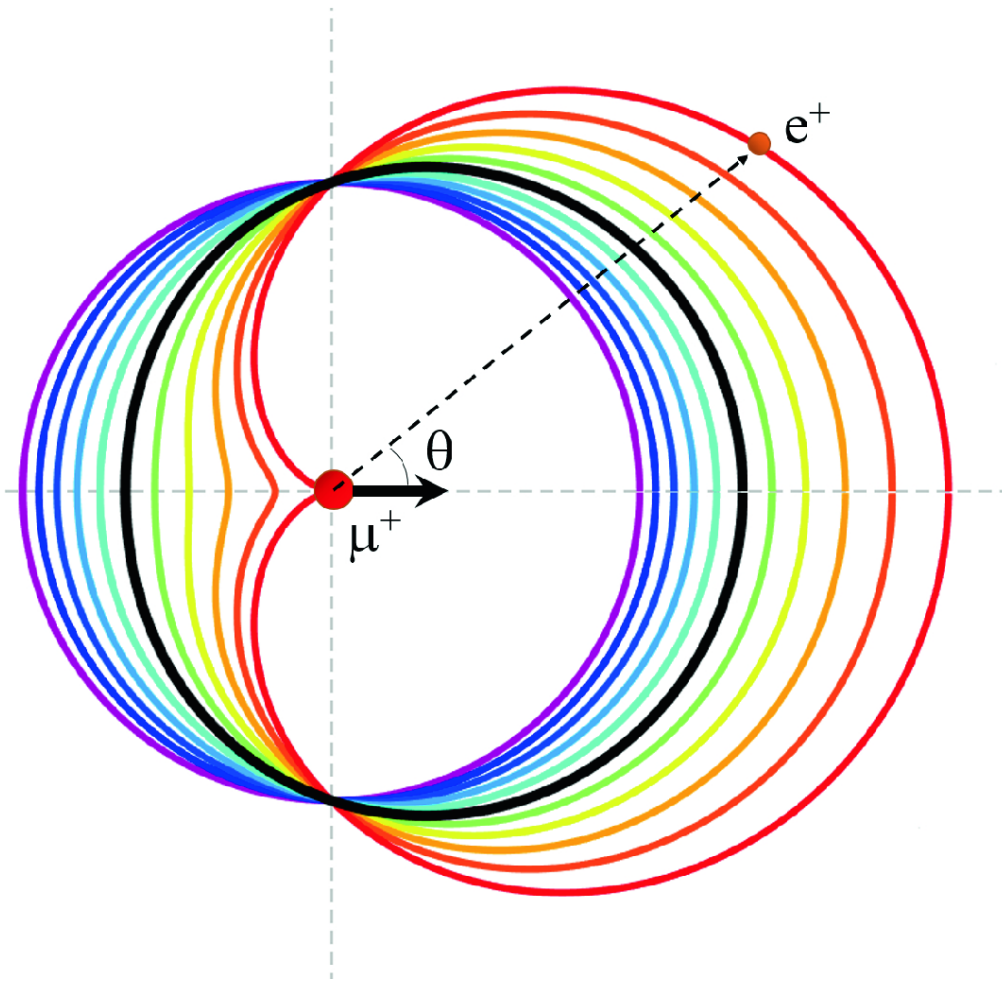
## *To watch*

- Мюони – частинки, які підтверджують теорію відносності [MinutePhysics] <https://youtu.be/nusyrYcoBjM>
- Распад элементарных частиц — Дмитрий Казаков / ПостНаука <https://youtu.be/YXBMSwqw6cA>
- Analyzing  $\mu$ SR Spectra, Stephen Blundell (2018) <https://youtu.be/T90ShkoysQ>
- Muon-spin rotation - Hugo Keller <https://youtu.be/mwcSp2PEcOw>

# Muon decay

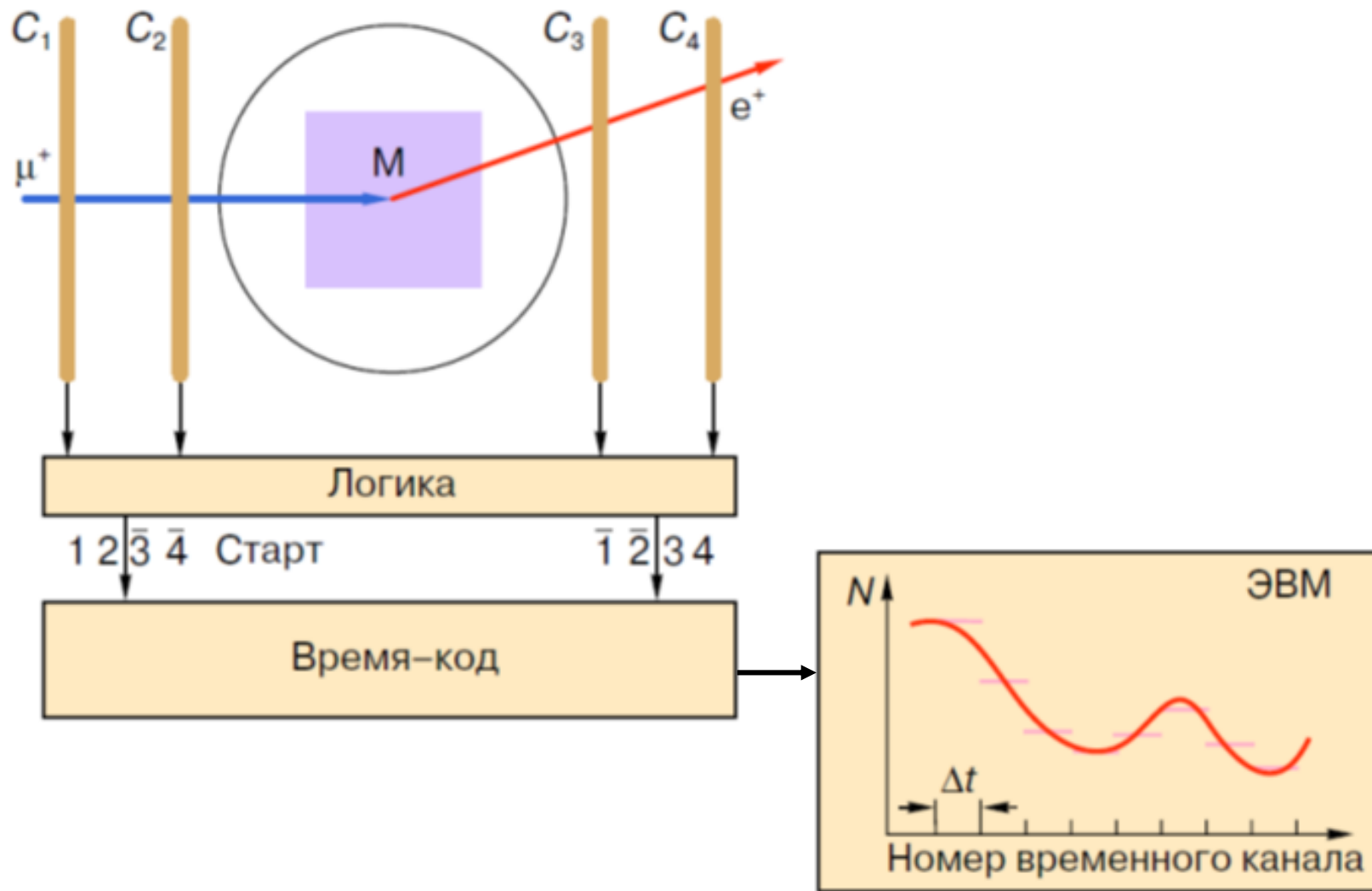


# Muon decay

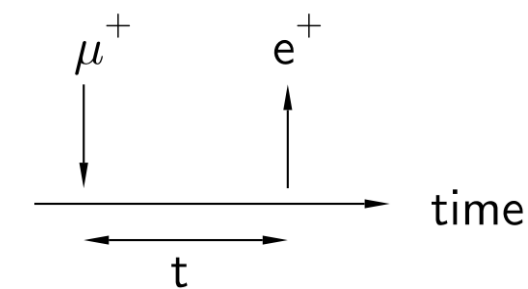


$$dW(\epsilon, \theta) = \frac{e^{t/\tau_\mu}}{\tau_\mu} [1 + a(\epsilon) \cos \theta] n(\epsilon) d\epsilon d\cos\theta dt,$$

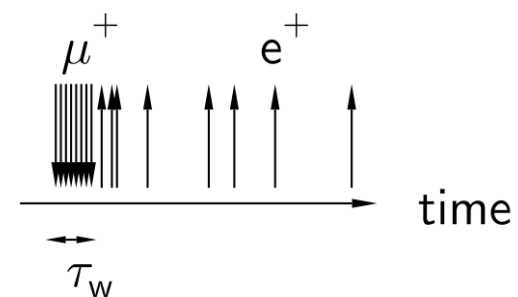
where  $a(\epsilon) = (2\epsilon - 1)/(3 - 2\epsilon)$ ,  $n(\epsilon) = 2\epsilon^2(3 - 2\epsilon)$  and the reduced positron energy  $\epsilon$  is defined as  $\epsilon = E/E_{\max}$ , where  $E_{\max}$  is the maximum positron energy  $E_{\max} = 52.83 \text{ MeV}$ .



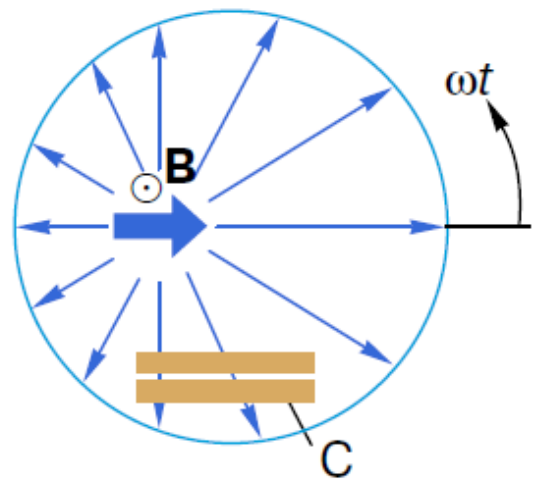
continuous wave



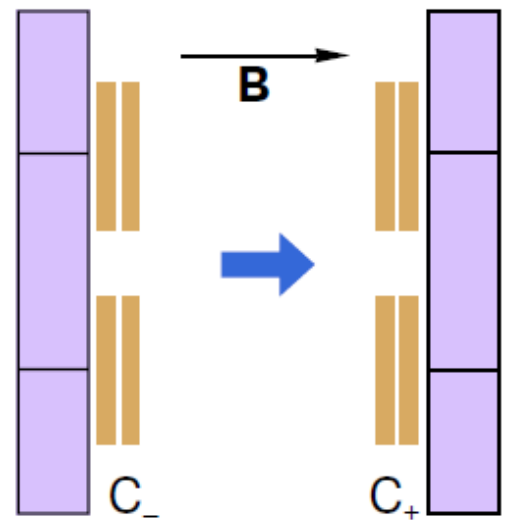
pulsed



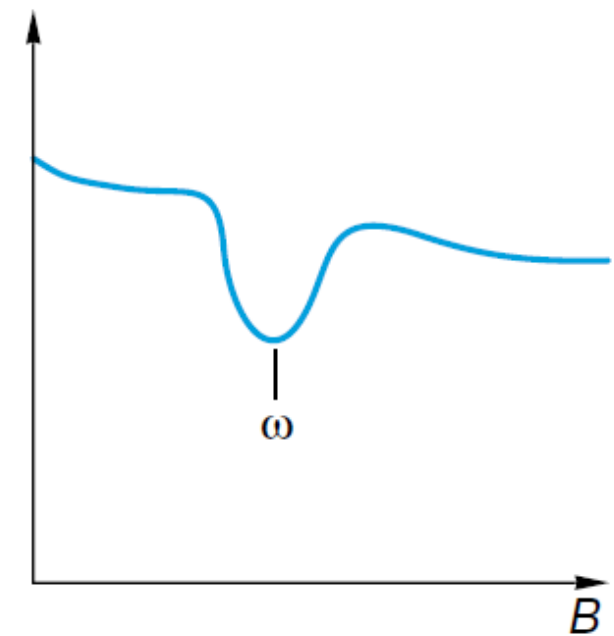
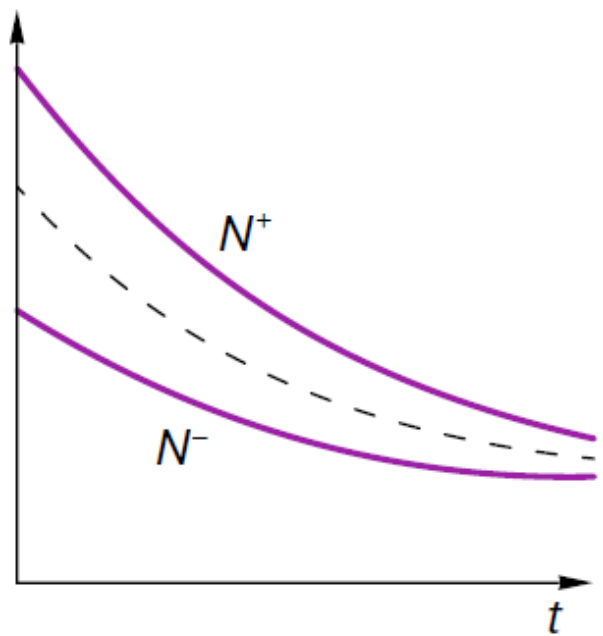
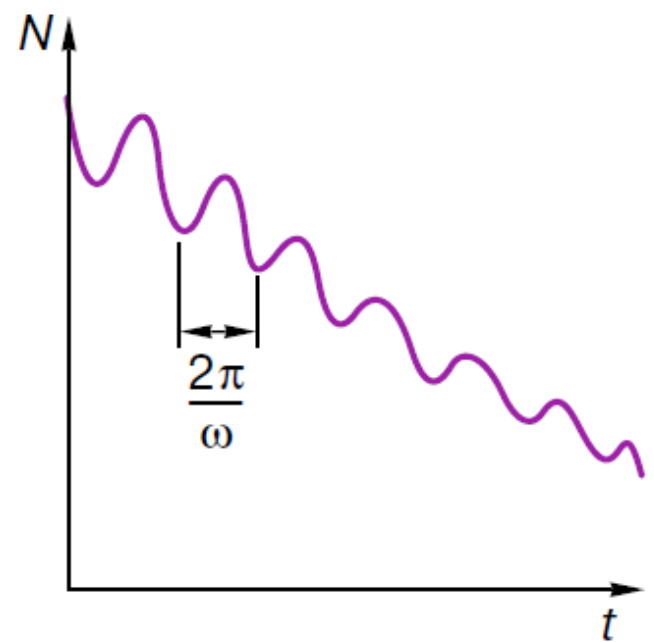
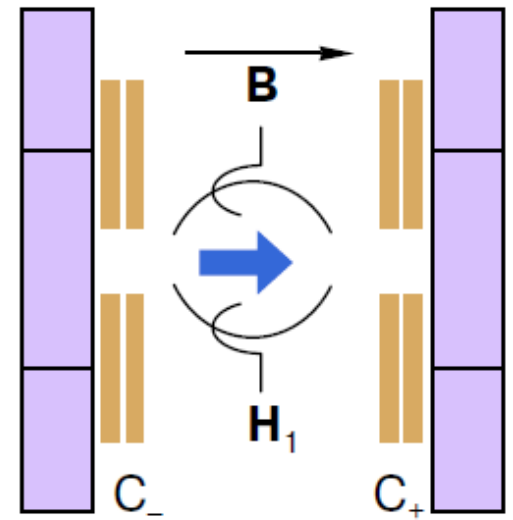
Вращение



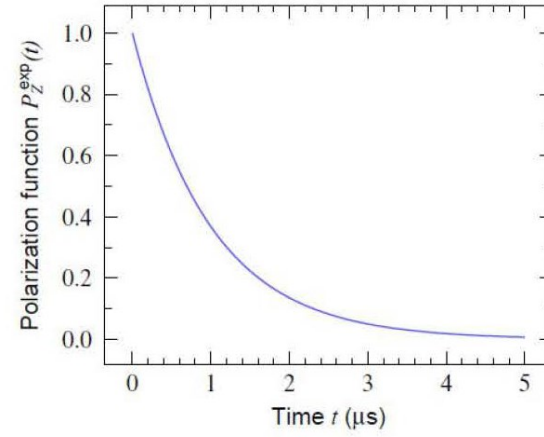
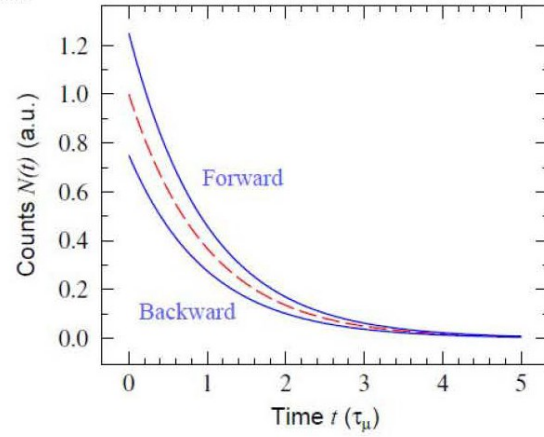
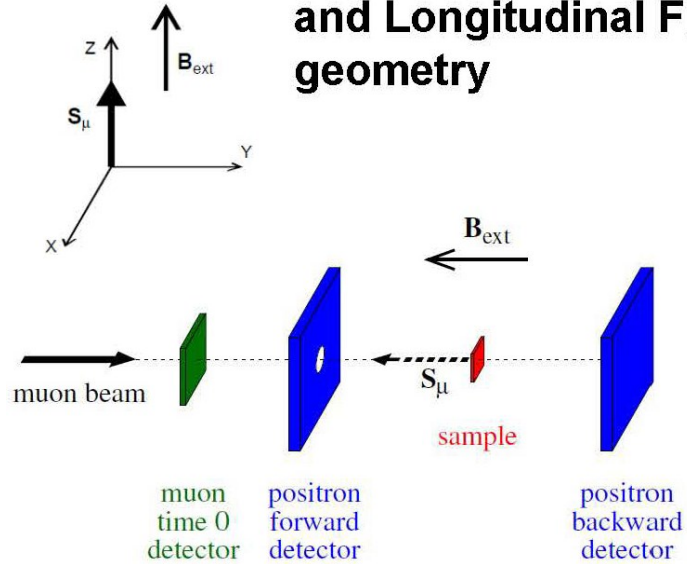
Релаксация



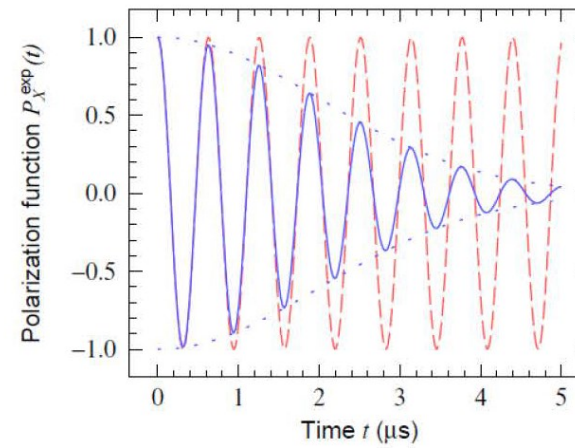
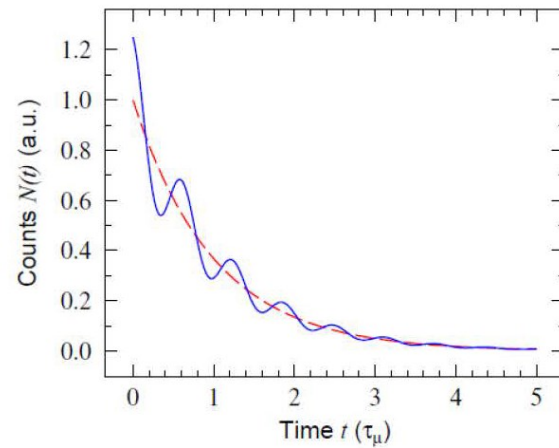
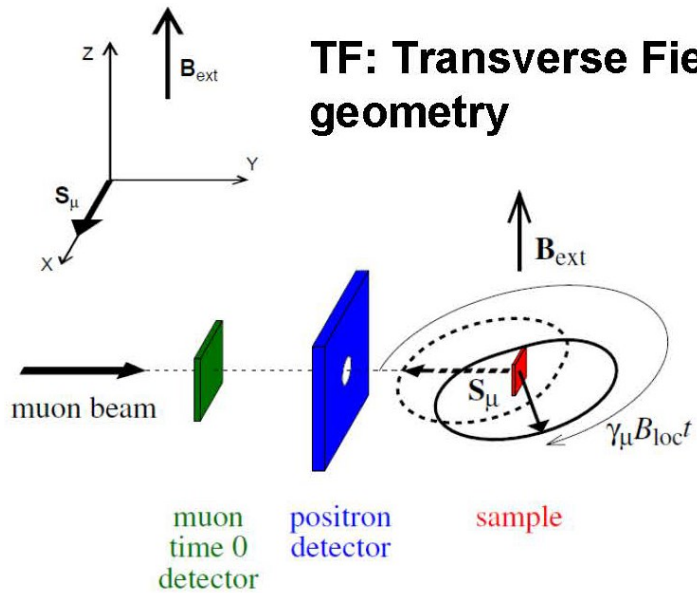
Резонанс



## ZF and LF: Zero field and Longitudinal Field geometry

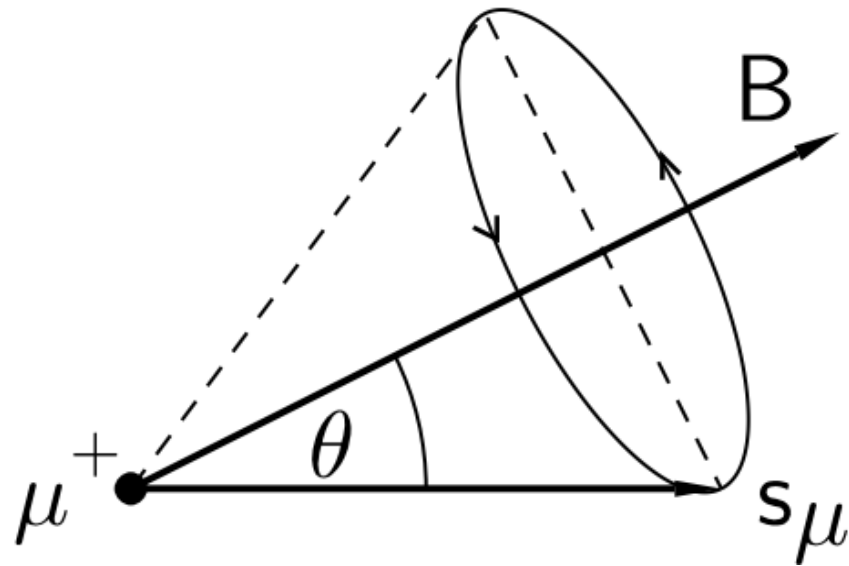


## TF: Transverse Field geometry





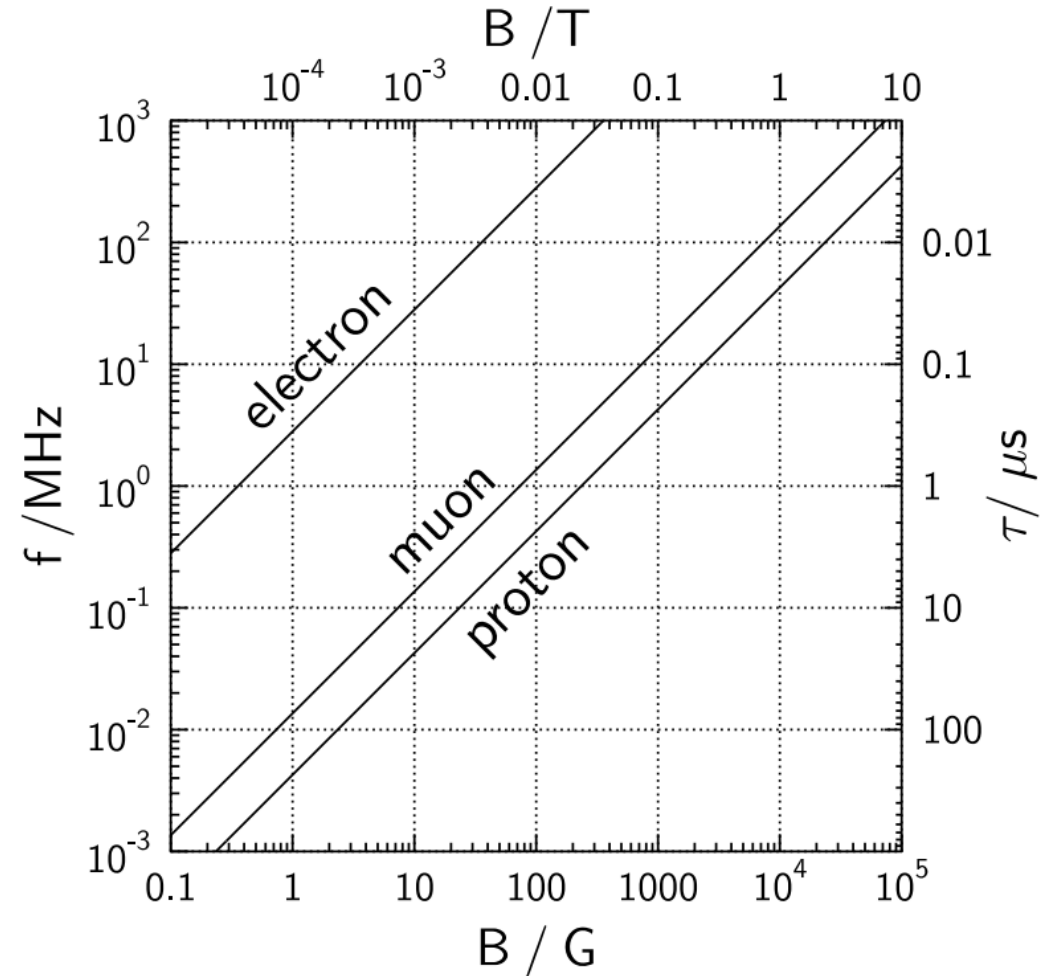
# Muon-spin precession

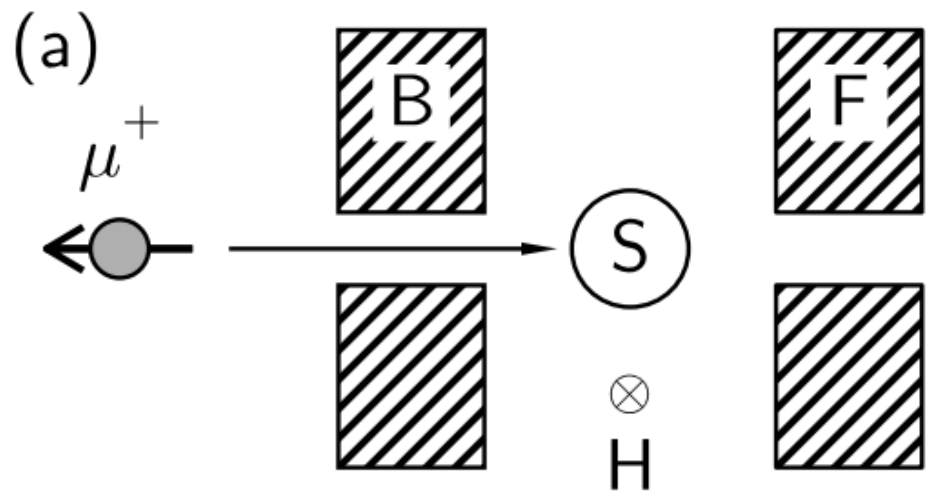


$$P_z(t) = \cos^2 \theta + \sin^2 \theta \cos(\gamma_\mu |B|t)$$

$|B|$  is the *modulus* of the local **dipolar** field

$$\gamma_\mu = ge/2m_\mu$$





Schematic illustration of a  $\mu$ SR experiment. (a) A spin-polarized beam of muons is implanted in a sample S. Following decay, positrons are detected in either a forward detector F or a backward detector B. If a transverse magnetic field H is applied to the sample as shown then the muons will precess. (b) The number of positrons detected in the forward and backward detectors.  $A(t)$  - the asymmetry function.

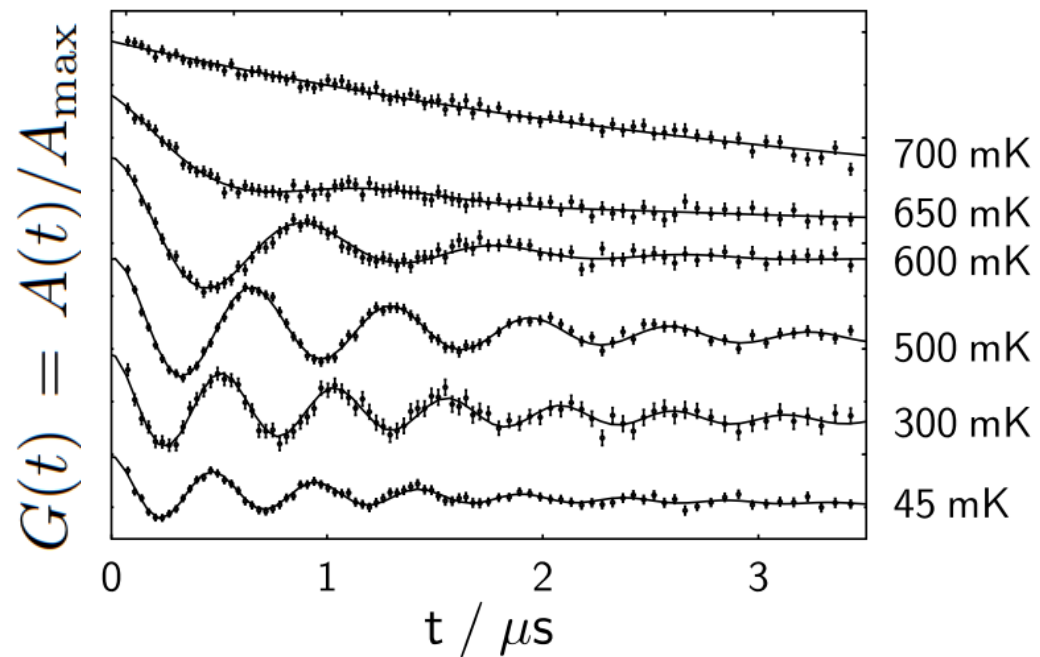
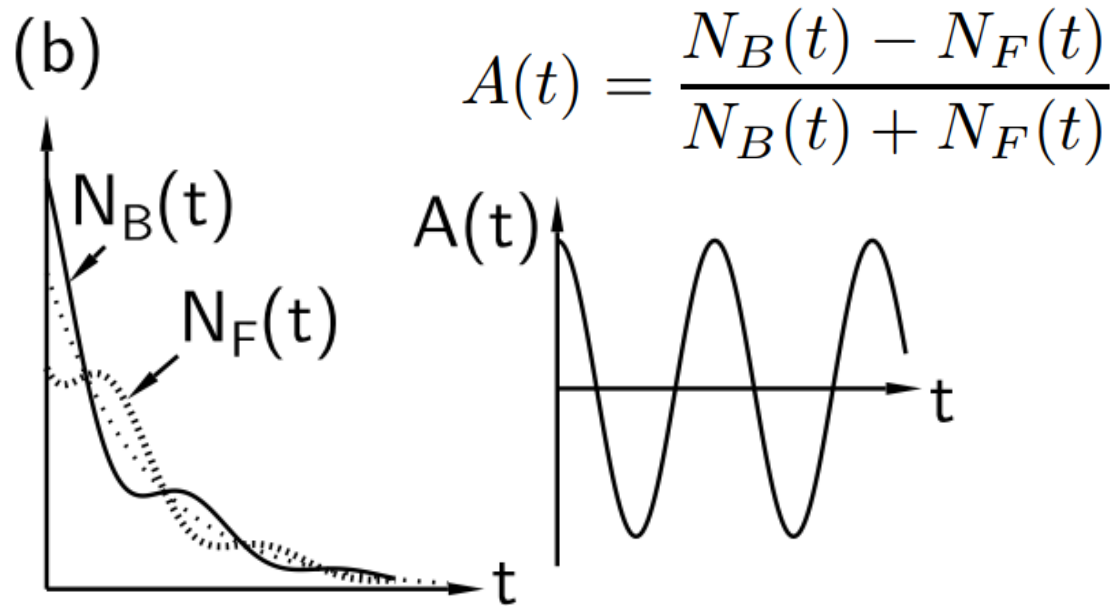


Figure 8: Zero-field muon spin rotation frequency in the organic ferromagnet  $p$ -NPNN (Blundell *et al.* 1995).

# Kubo-Toyabe relaxation

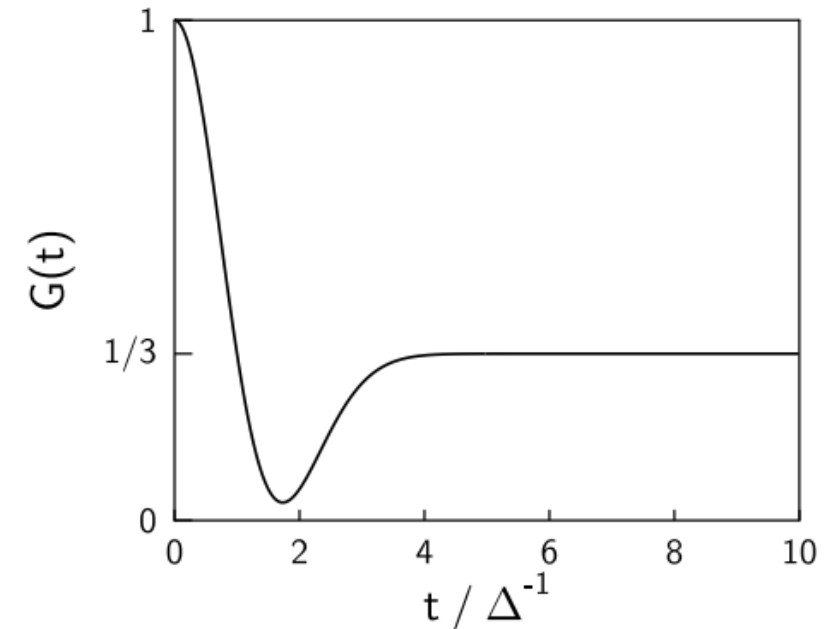
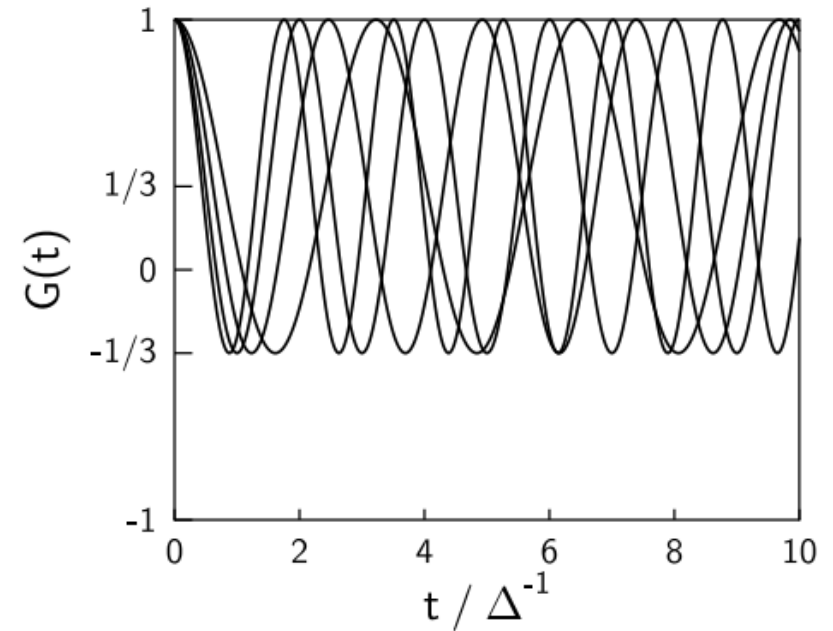
$$G(t) = \cos^2 \theta + \sin^2 \theta \cos(\gamma_\mu B t)$$

If the direction of the local magnetic field is entirely random then averaging over all directions would yield

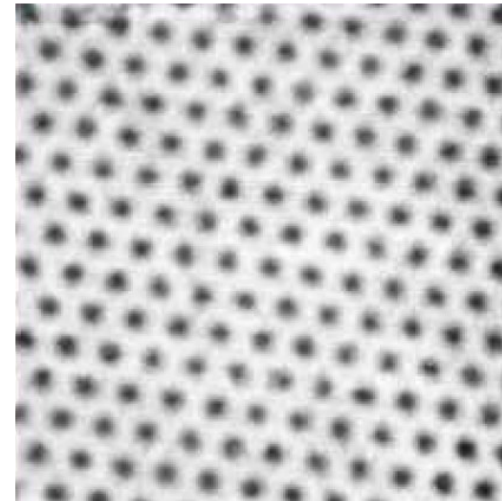
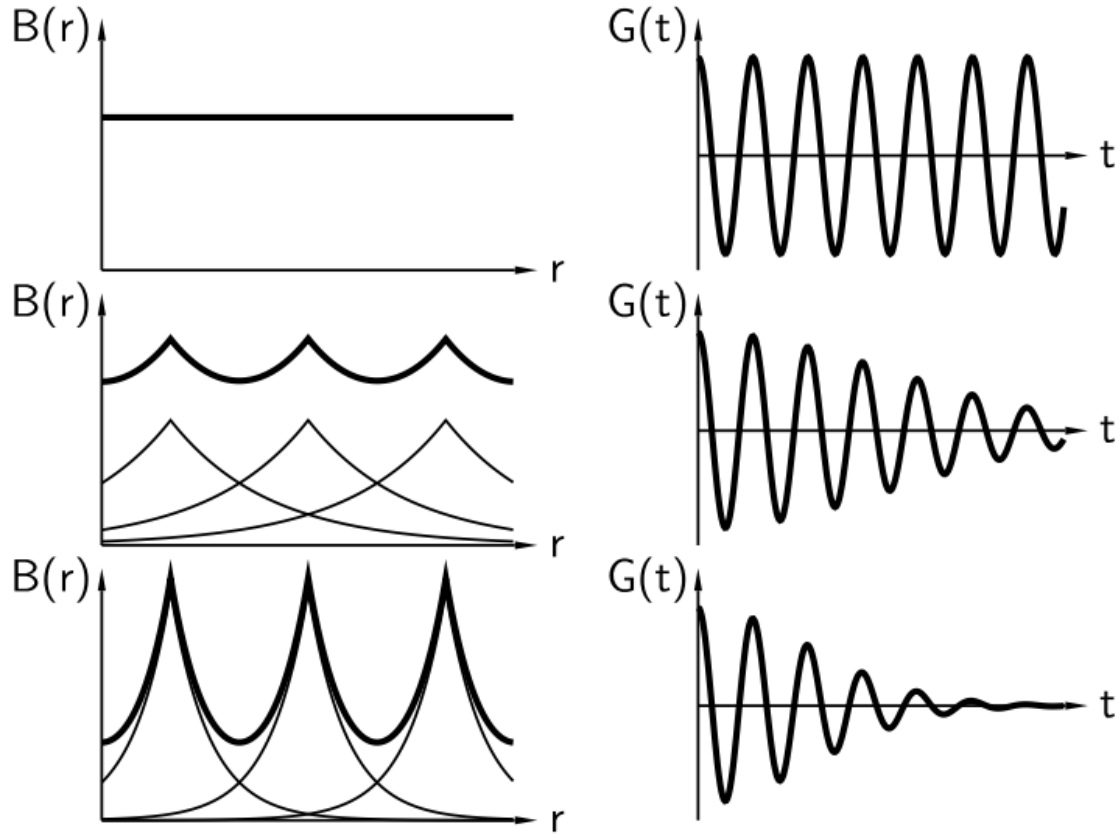
$$G(t) = \frac{1}{3} + \frac{2}{3} \cos(\gamma_\mu B t)$$

If the strength of the local magnetic field is taken from a Gaussian distribution of width  $\Delta/\gamma_\mu$  centred around zero, then a straightforward averaging over this distribution gives

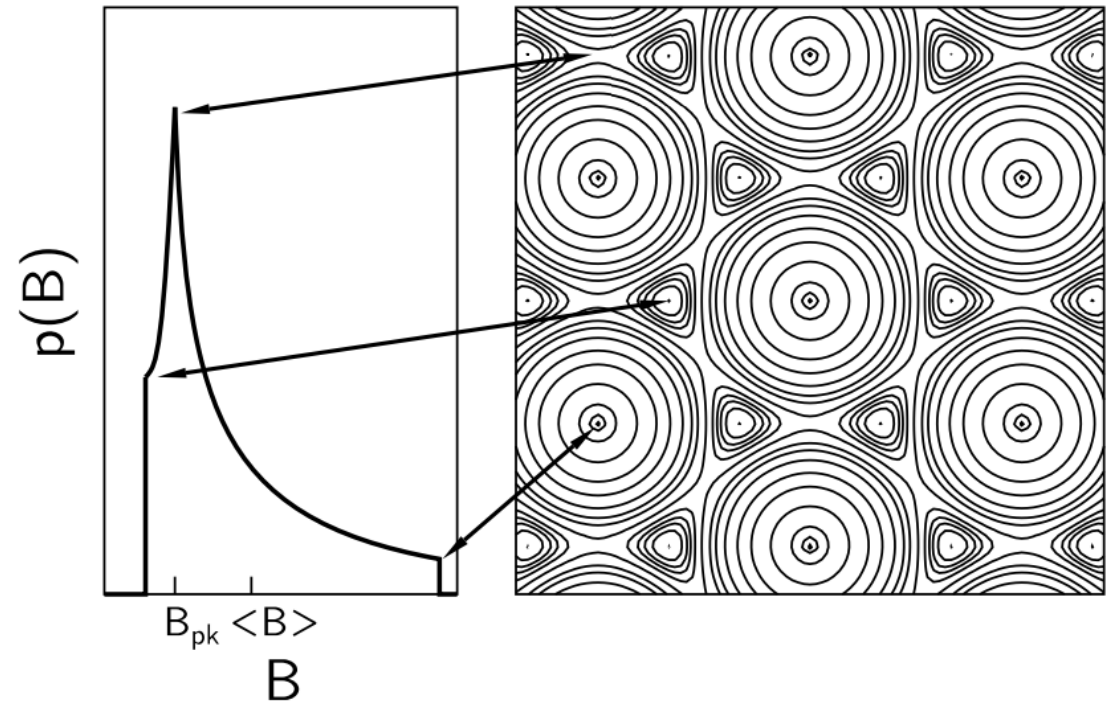
$$G(t) = \frac{1}{3} + \frac{2}{3} e^{-\Delta^2 t^2 / 2} (1 - \Delta^2 t^2)$$



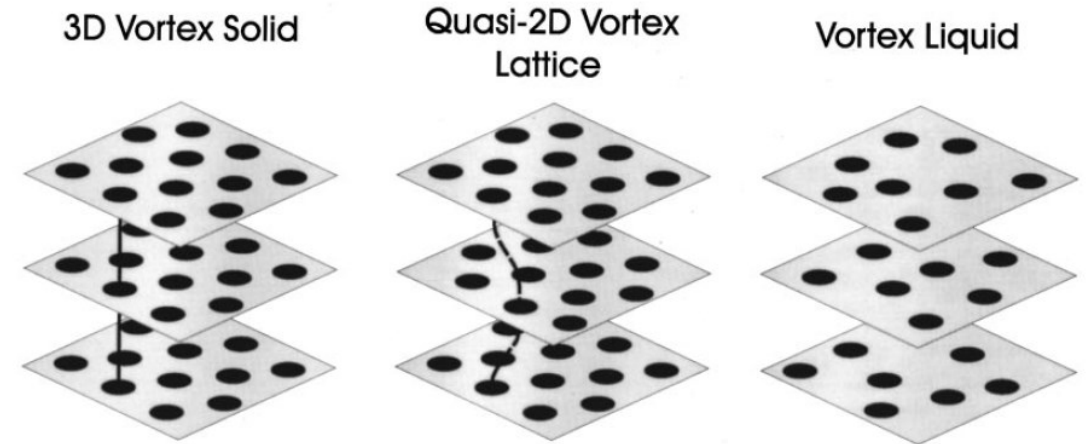
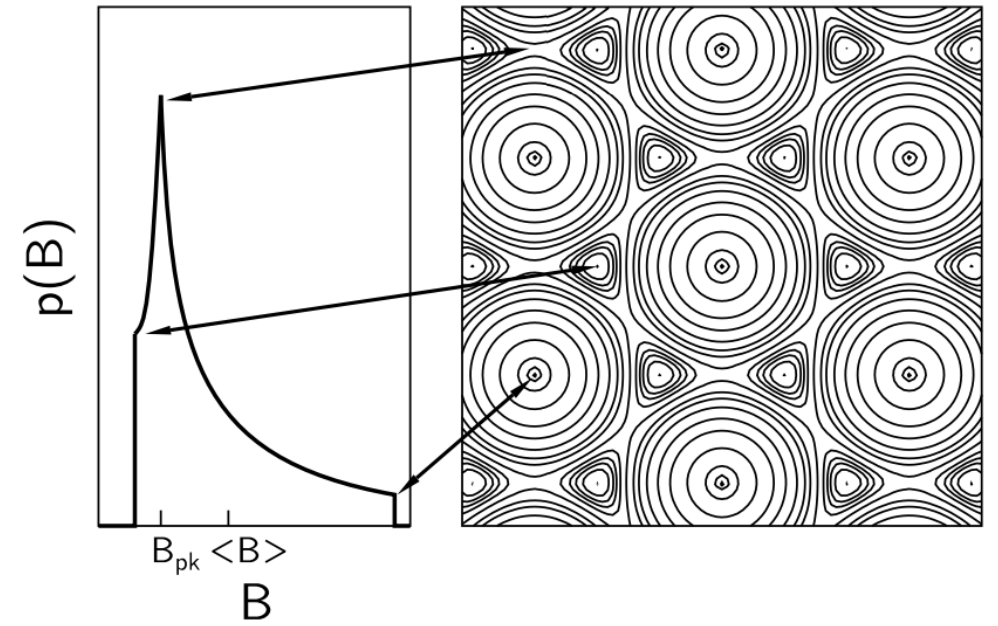
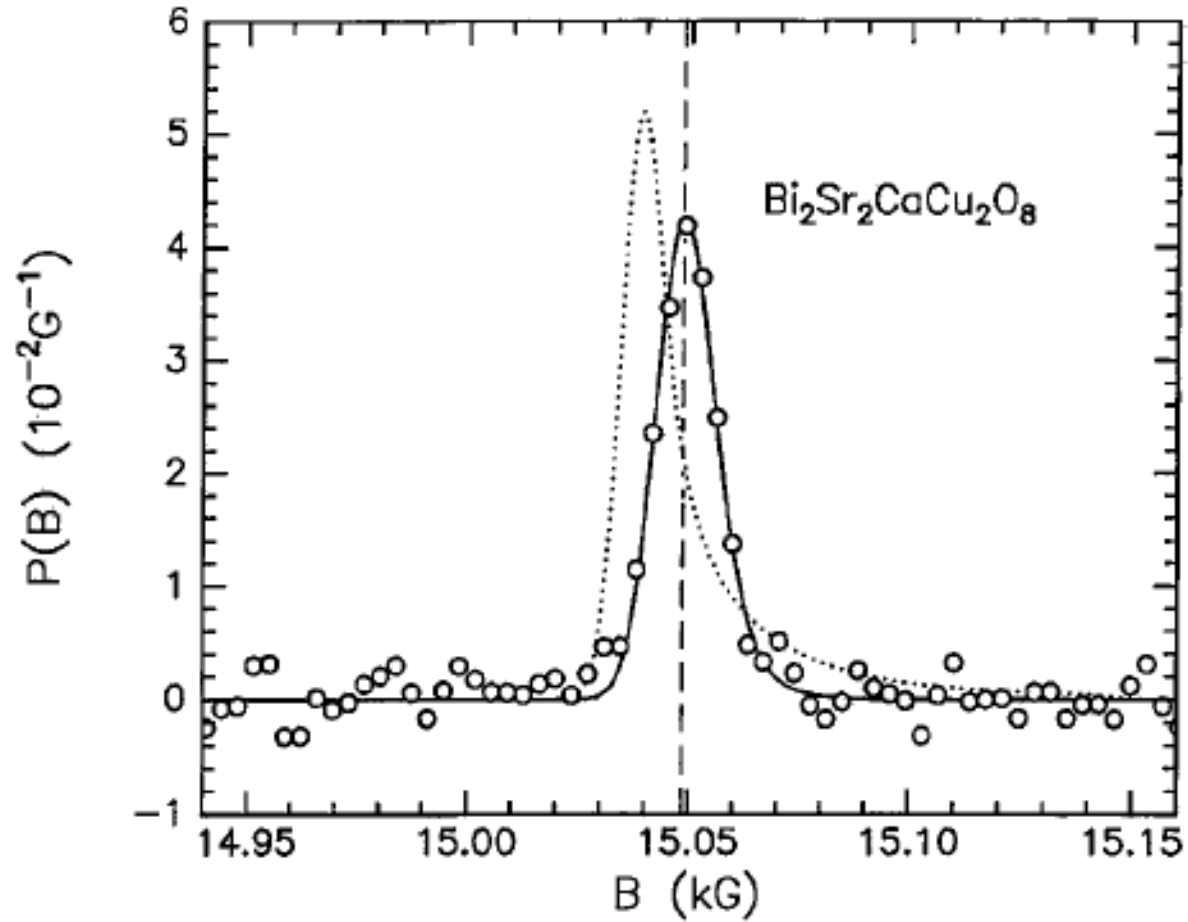
# Abrikosov vortices



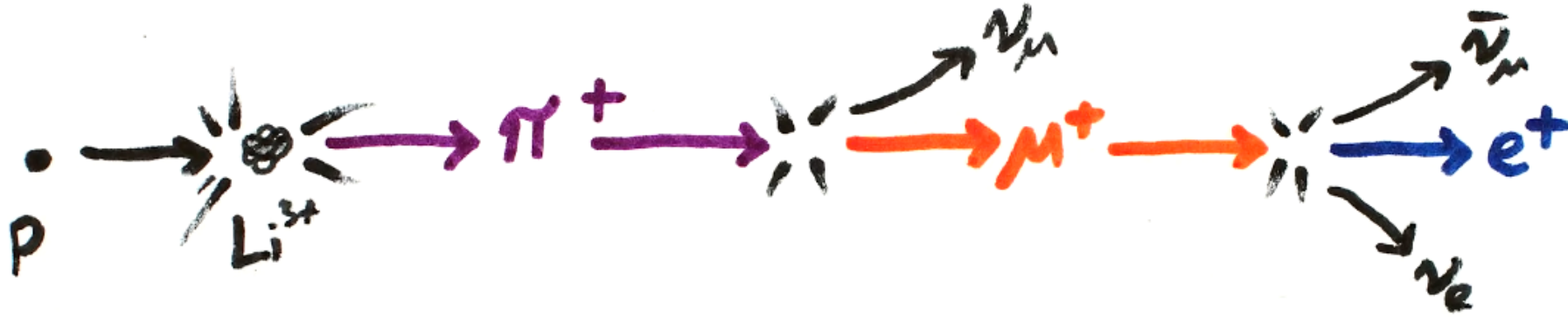
$$\sigma = \gamma_{\mu} \langle B(\mathbf{r}) - \langle B(\mathbf{r}) \rangle_{\mathbf{r}} \rangle_{\mathbf{r}}^{1/2} \approx 0.0609 \gamma_{\mu} \Phi_0 / \lambda^2$$



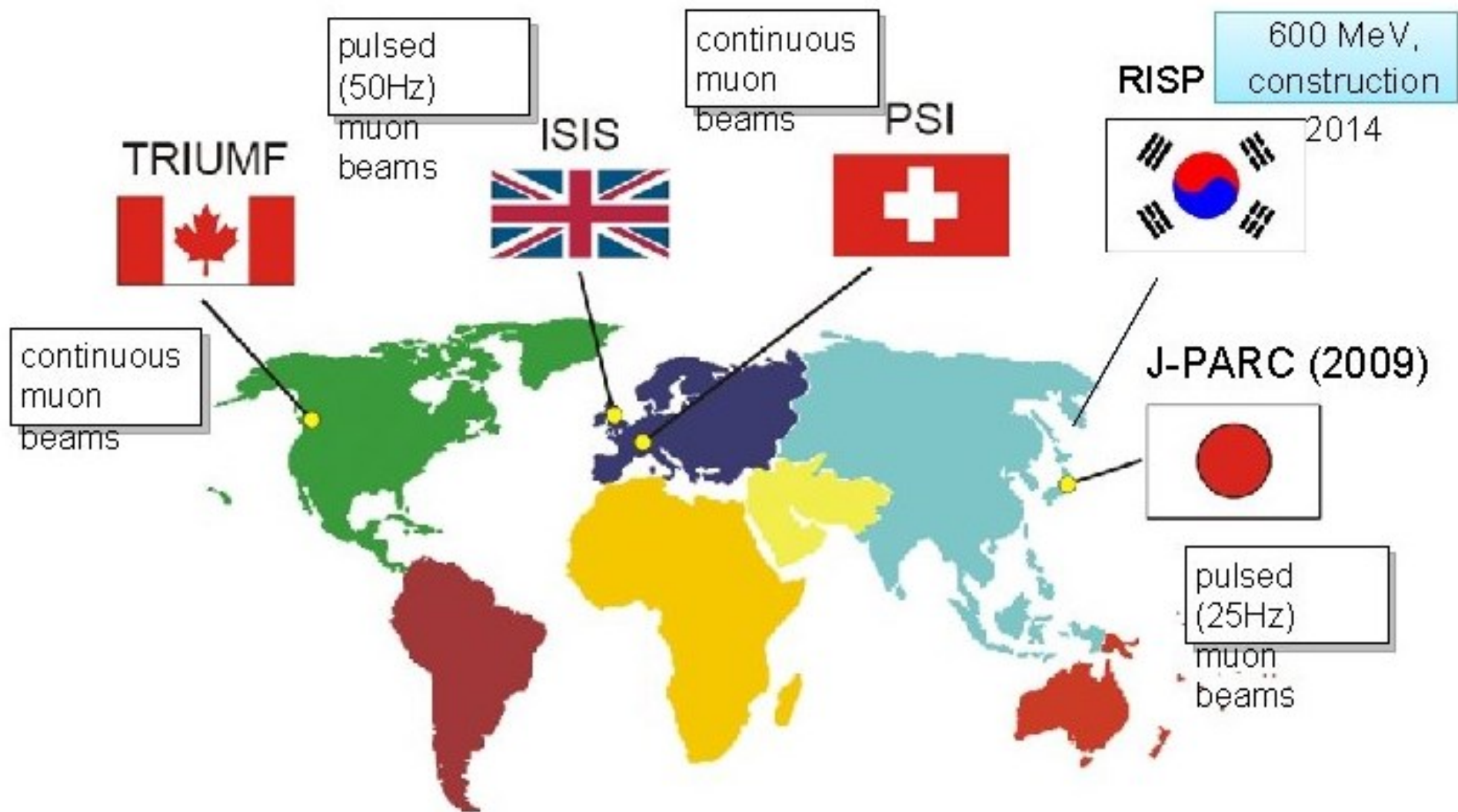
# Abrikosov vortices



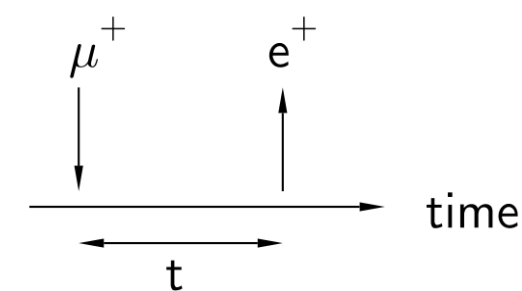
# Muon decay



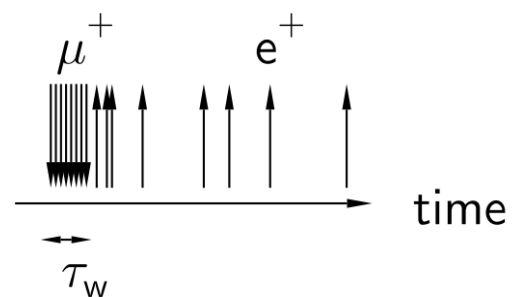
# $\mu$ SR Facilities around the World

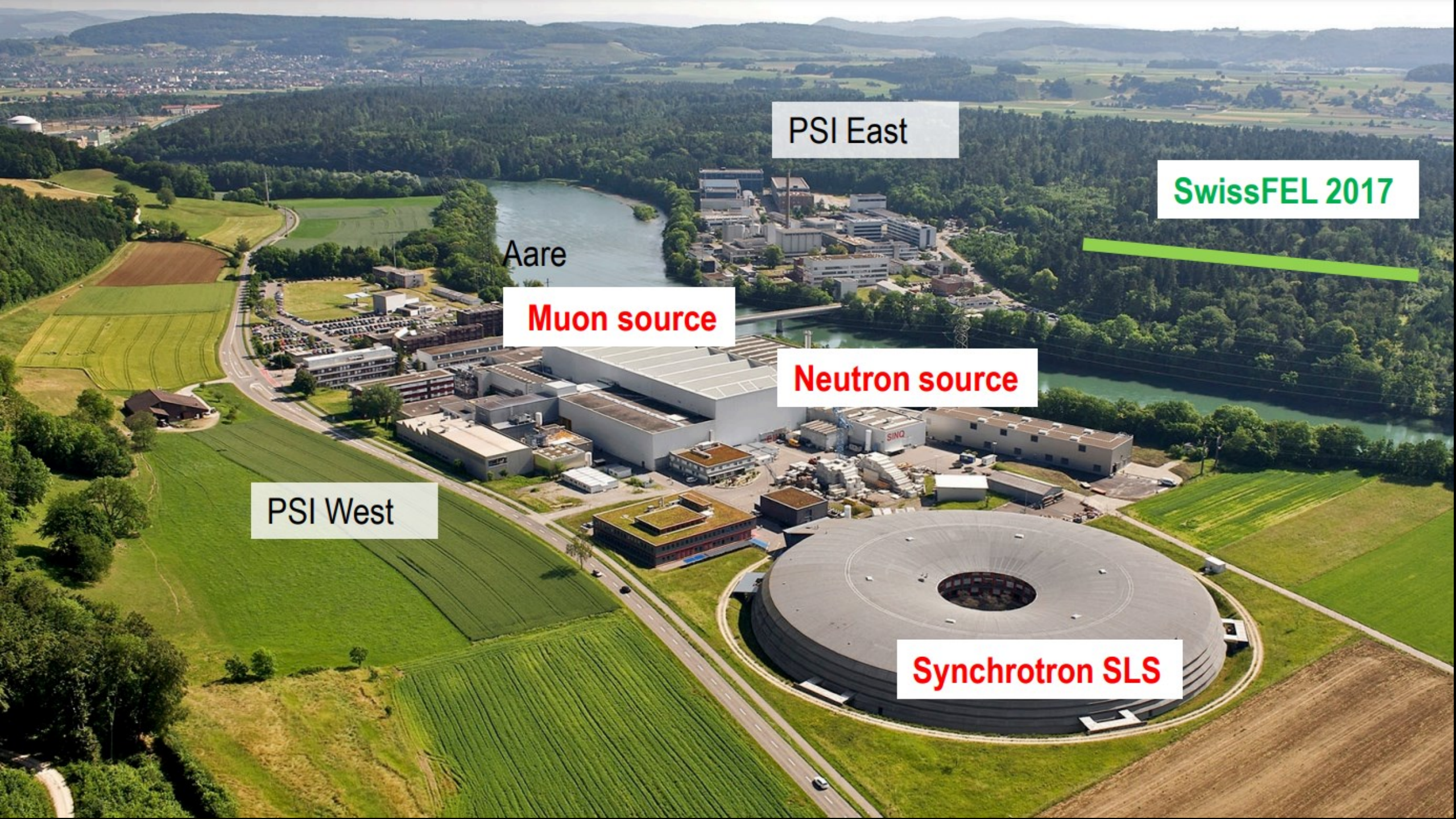


continuous wave



pulsed





PSI East

SwissFEL 2017

Aare

Muon source

Neutron source

PSI West

Synchrotron SLS



# Muon Instruments at PSI : S $\mu$ S (Swiss Muon Source)

## HAL-9500

High Field and Low Temperature

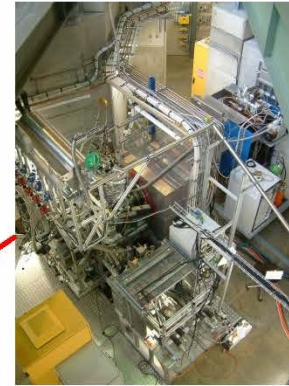
9.5 T

< 20 mK



590 MeV

2.2 mA



## LEM

Low-energy muon beam and instrument, tunable energy (0.5-30 keV,  $\mu^+$ ), thin-film, near-surface and multi-layer studies (1-300 nm)

0.3 T

2.5 K

## DOLLY

General Purpose Surface Muon Instrument  $\mu^+$  energy: 4.2 MeV

0.5 T

250 mK

## GPS

General Purpose Surface Muon Instrument  
Muon energy: 4.2 MeV ( $\mu^+$ )

0.6 T, 1.6 K



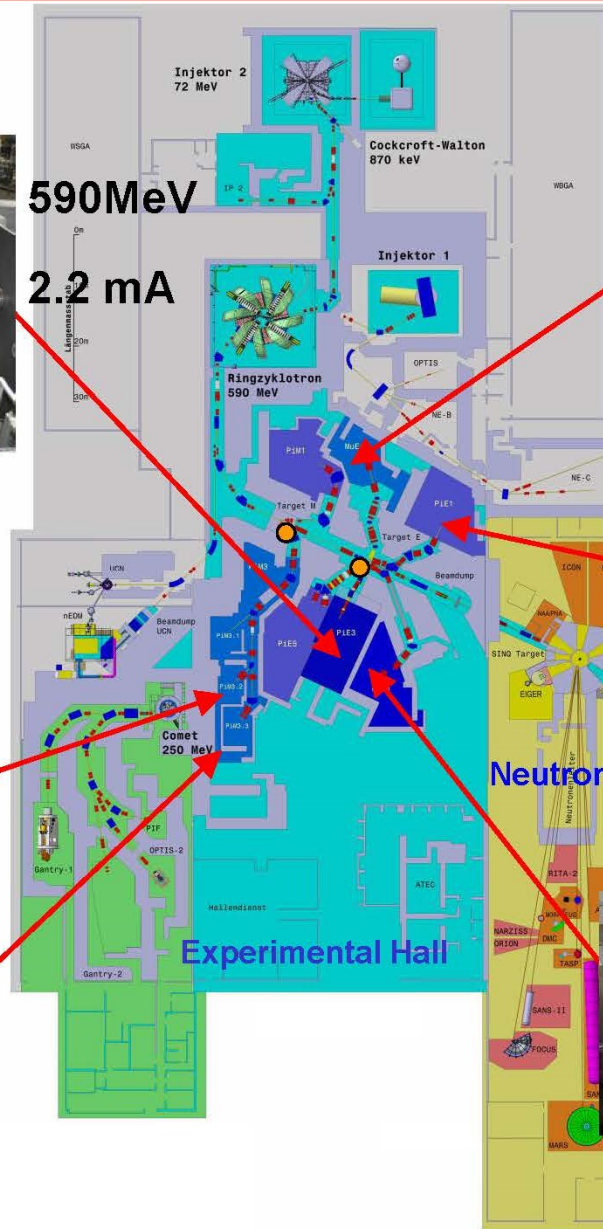
Shared Beam Surface Muon Facility (Muon On REquest)

## LTF

Low Temperature Facility  
Muon energy: 4.2 MeV ( $\mu^+$ )

3 T,

20 mK- 4 K



Neutron Hall

Experimental Hall



## GPD

General Purpose Decay Channel Instrument  
Pressure studies

Muon energy: 5 - 60 MeV ( $\mu^+$  or  $\mu^-$ )

0.5 T,

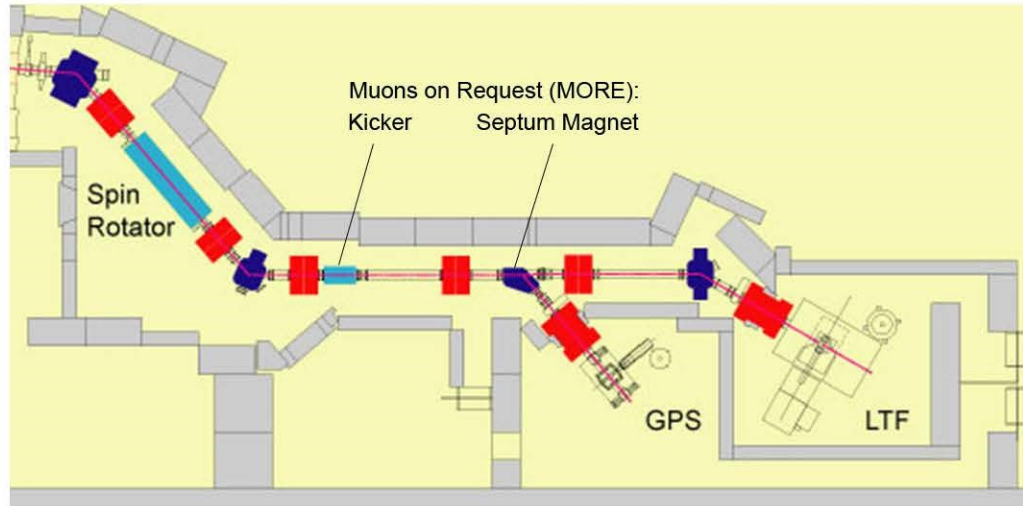
300 mK

2.8 GPa



# Surface Muon Instruments – GPS/LTF/Dolly

$\pi$ M3 beam line: shared by: General Purpose Spectrometer and Low Temperature Facility



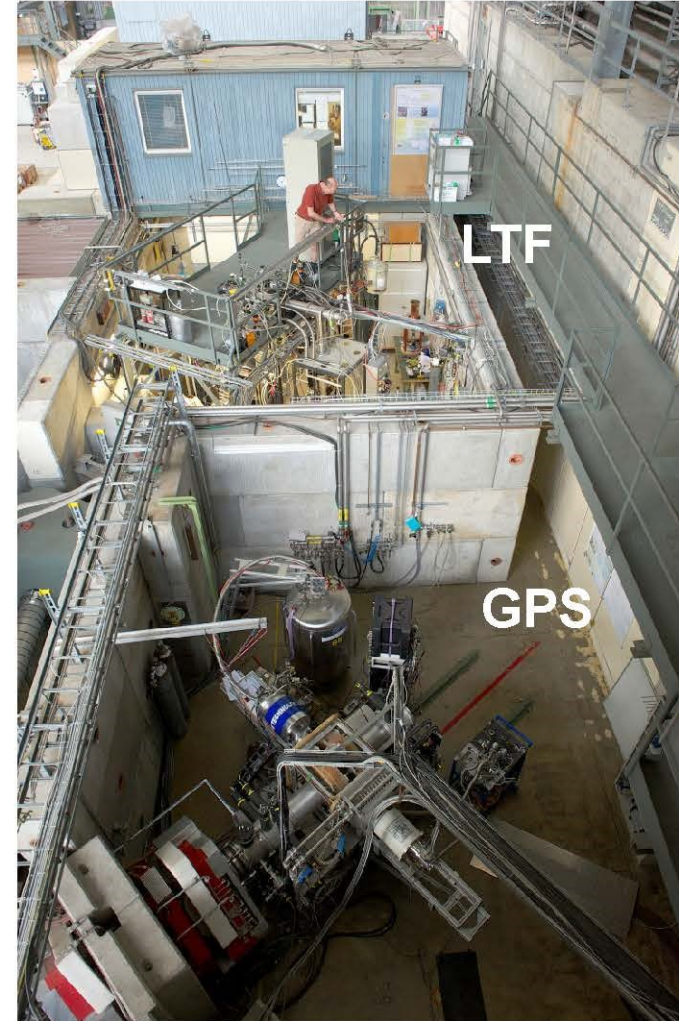
- 4 MeV  $\mu^+$ , 100% polarized
- $B_{\text{ext}}$ 
  - GPS: 0 - 0.6 T
  - Dolly: 0 - 0.6 T
  - LTF: 0 - 3 T
- T
  - GPS: 1.8 - 1200 K
  - Dolly: 0.3 – 300 K
  - LTF(DR): 0.02 - 4.2 K

Veto system for low background and small samples:

Sample size:

~ 2 mm DIA

or ~ **30 mg**



# HAL-9500: High field And Low temperature $\mu$ SR (9.5 T, < 20 mK)

---

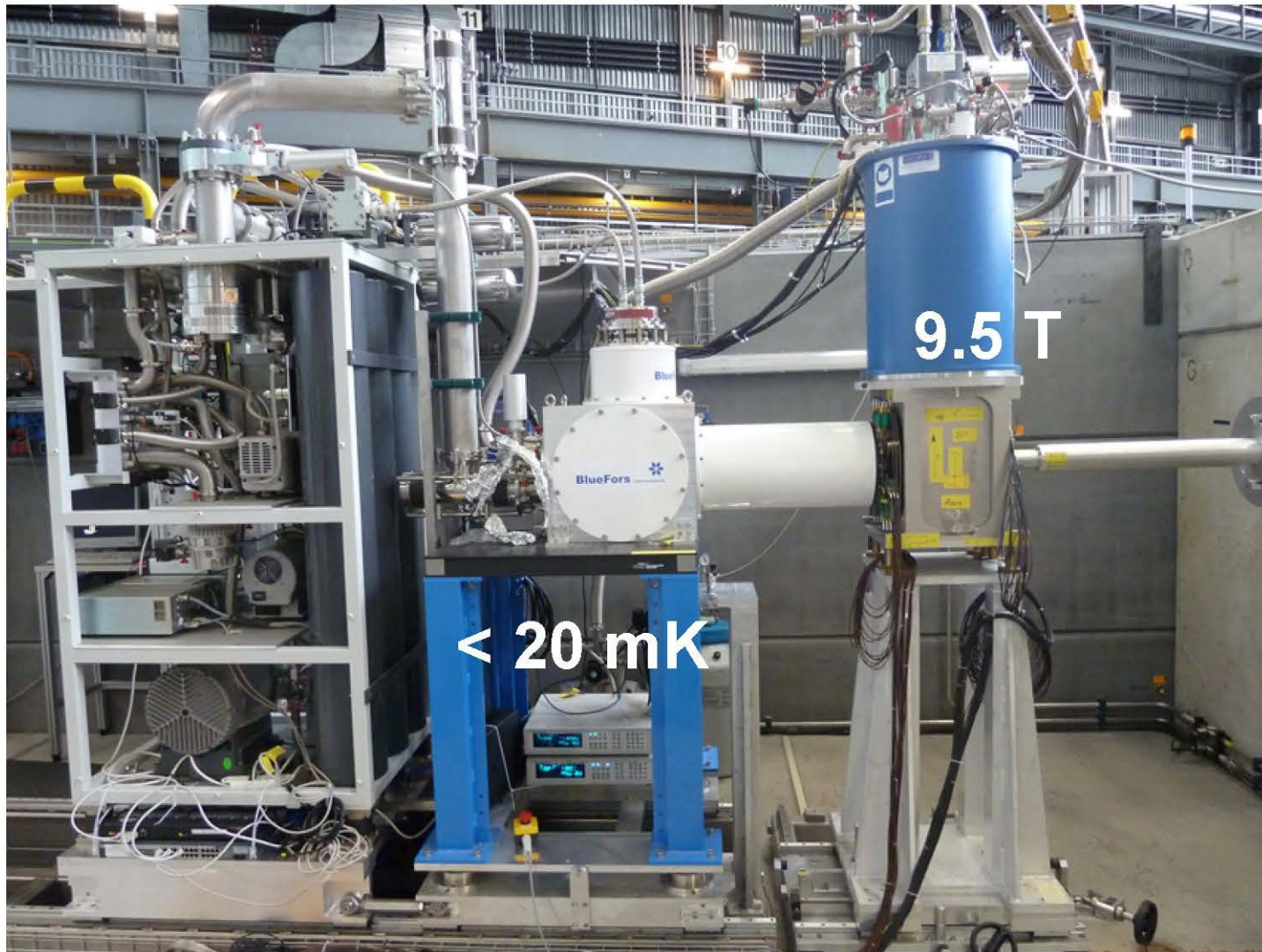


Instrument



Beam line with 90° spin rotator

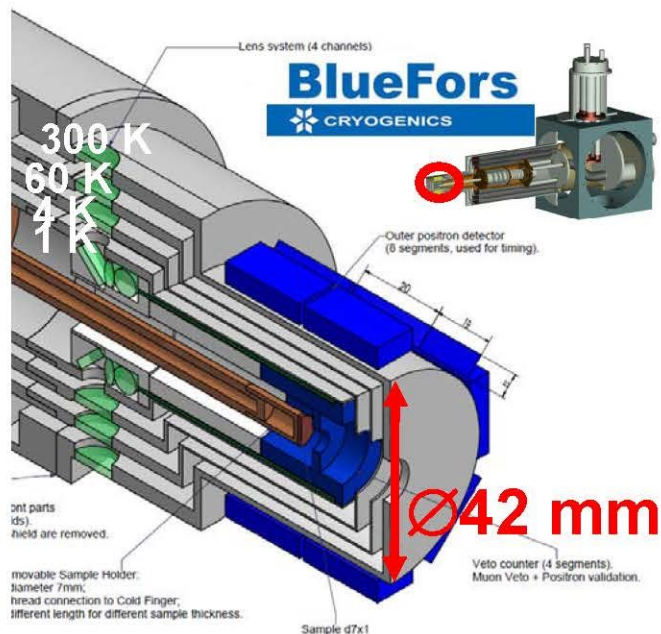
## Horizontal dilution refrigerator, high field homogeneity ( $10^{-4}$ T)



Option with He-Flow Cryostat 2K- 320 K

*R. Scheuermann et al.*

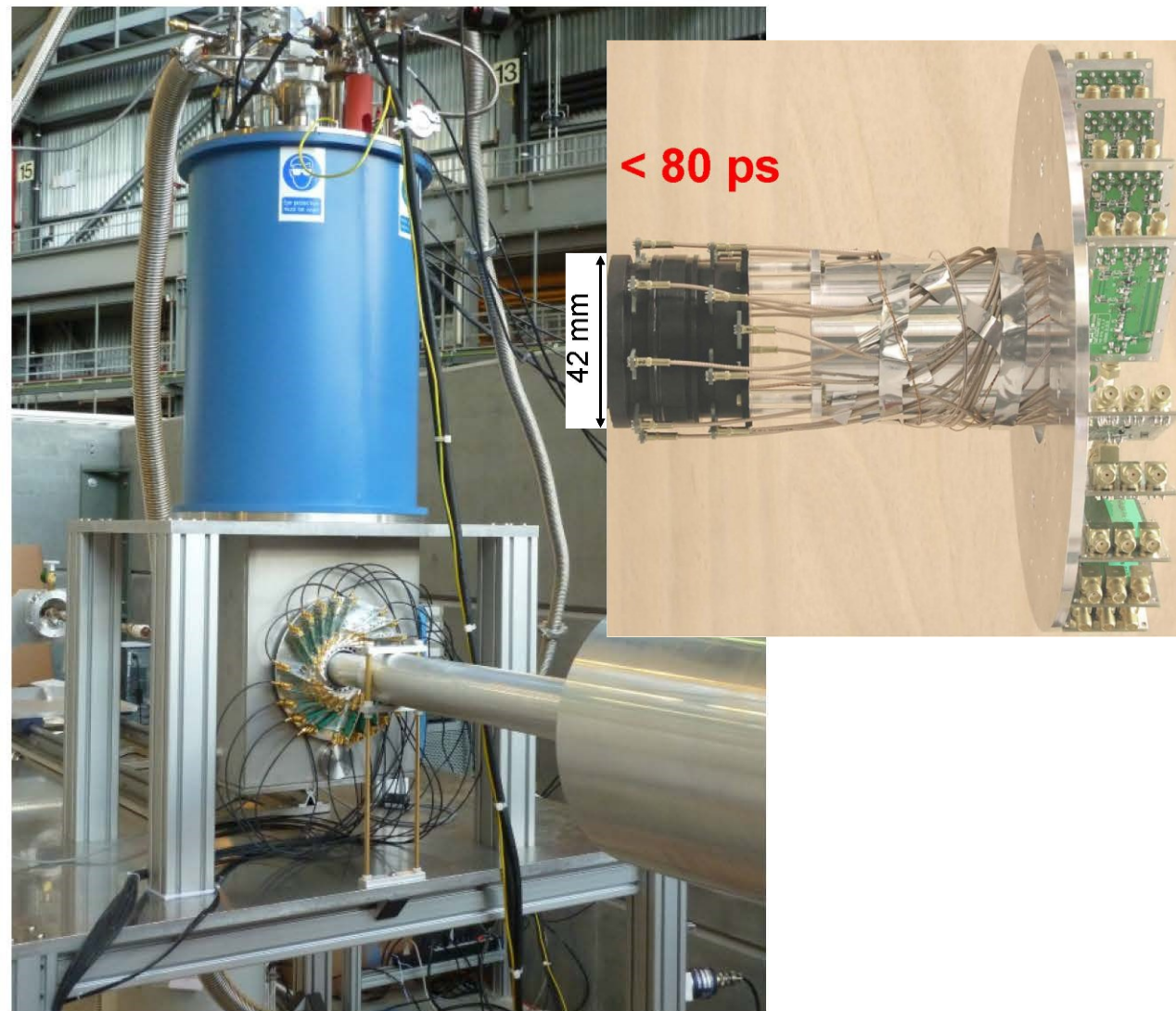
# A compact and fast $\mu$ SR spectrometer



Need spectrometer which is:

- **very compact** (spiraling radius of 30 MeV decay positron: 1 cm in 10 T)
- **very fast** Larmor precession frequency at 10 T  $\nu_L \sim 1.35$  GHz (remember  $A(\nu_L)$ )
- **field insensitive**

→ Solid state detectors: Avalanche Photo Diodes (APDs)

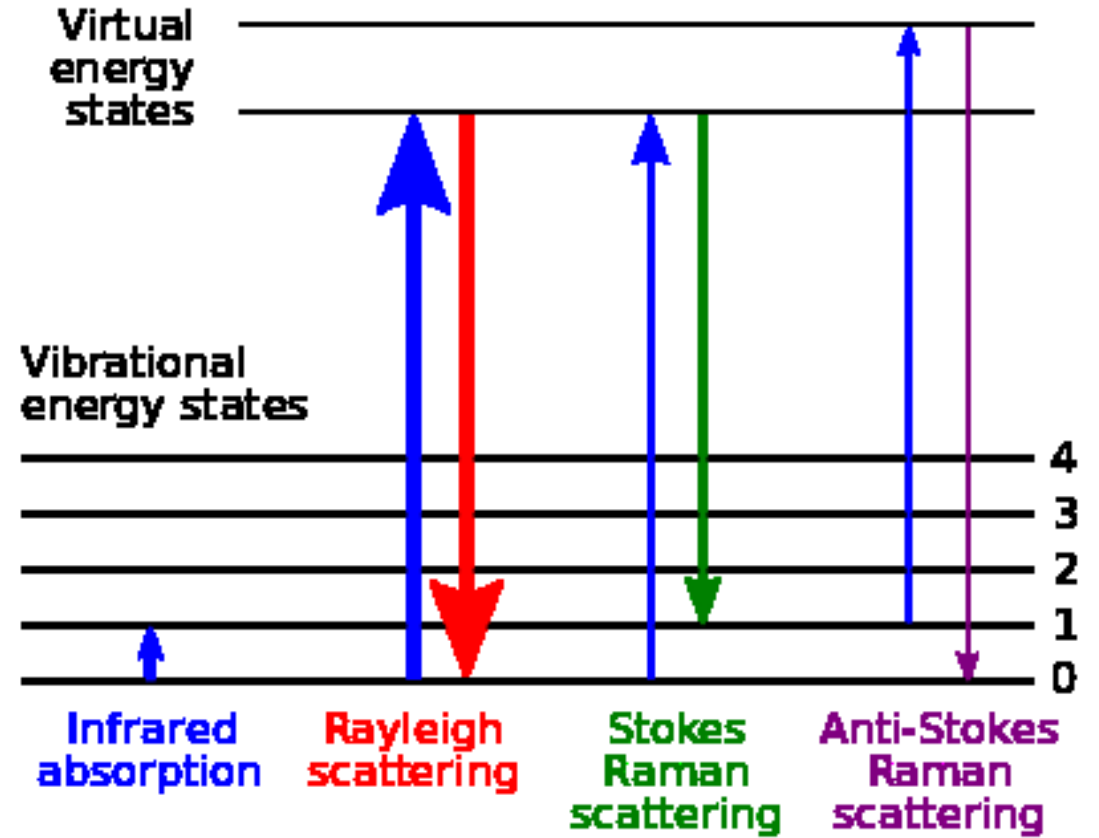
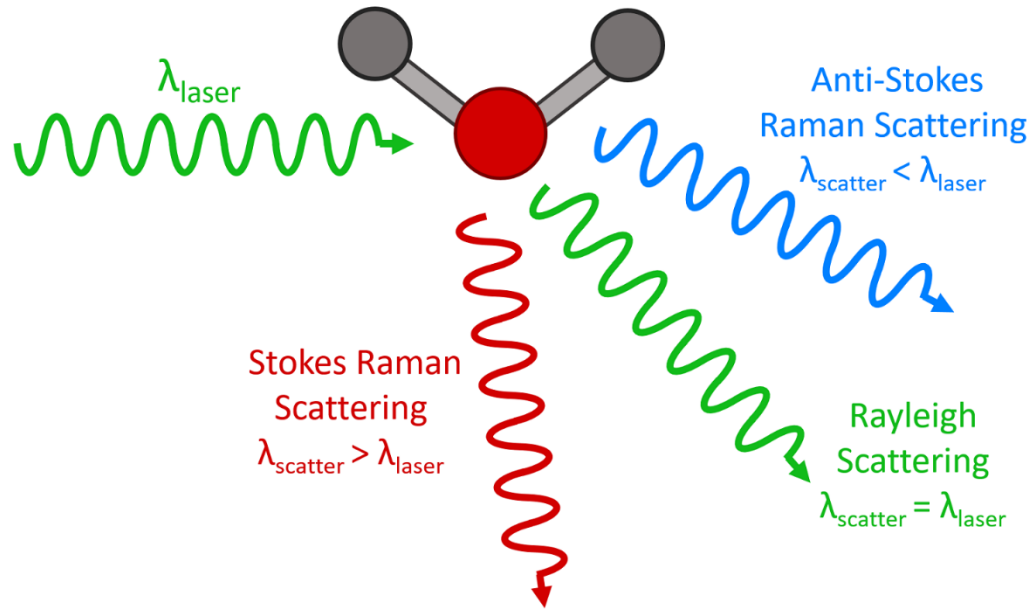


# Raman spectroscopy

## *To read*

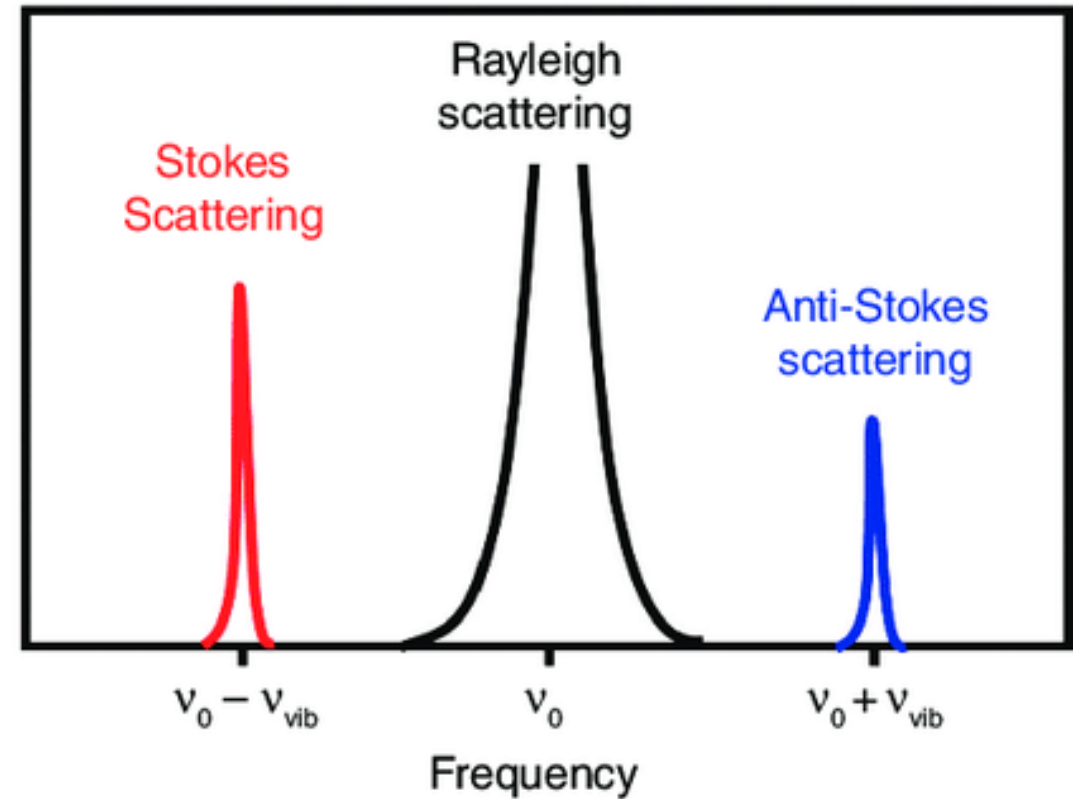
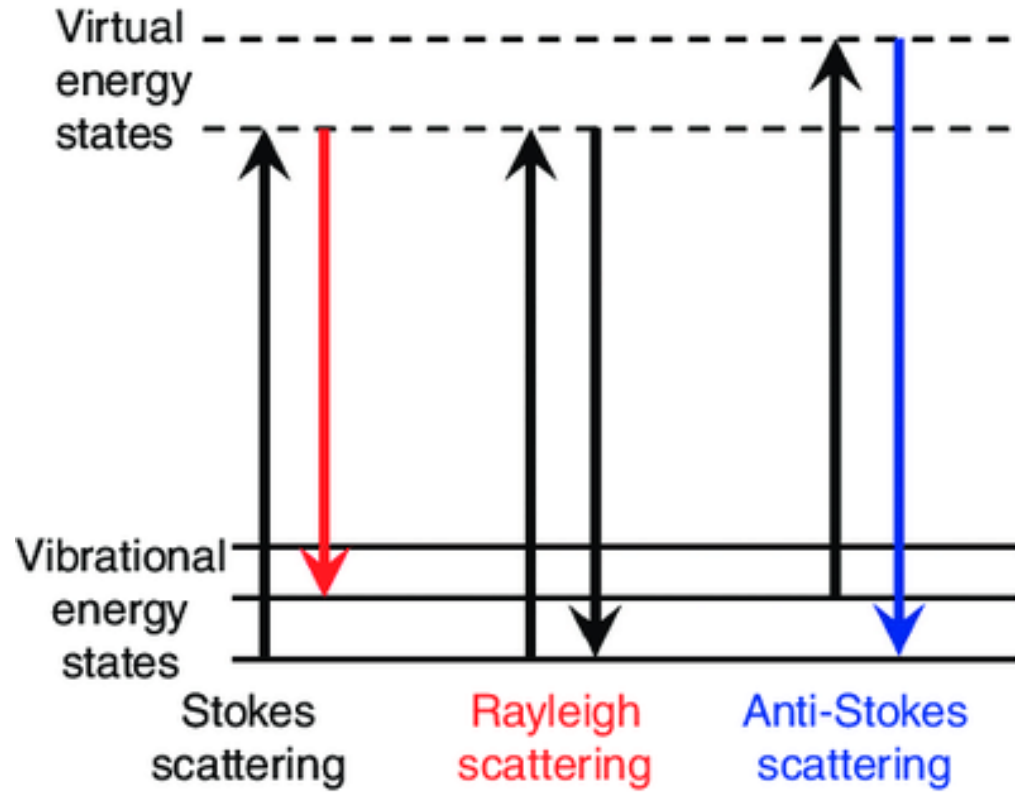
1. Modern Raman spectroscopy : a practical approach / Ewen Smith, Geoff Dent  
[http://www.chemistry.uoc.gr/lapkin/Modern\\_Raman\\_Spectroscopy\\_A\\_Practical\\_Approach.pdf](http://www.chemistry.uoc.gr/lapkin/Modern_Raman_Spectroscopy_A_Practical_Approach.pdf)
2. M. Cardona. Raman scattering in high-Tc superconductors. *Physica C* **185-189**, 65 (1991) [pdf](#)
3. B. Moritz et al. An investigation of particle-hole asymmetry in the cuprates via electronic raman scattering. *Physical Review B* **84**, 235114 (2011)  
<https://arxiv.org/pdf/1106.5798.pdf>

# Raman spectroscopy

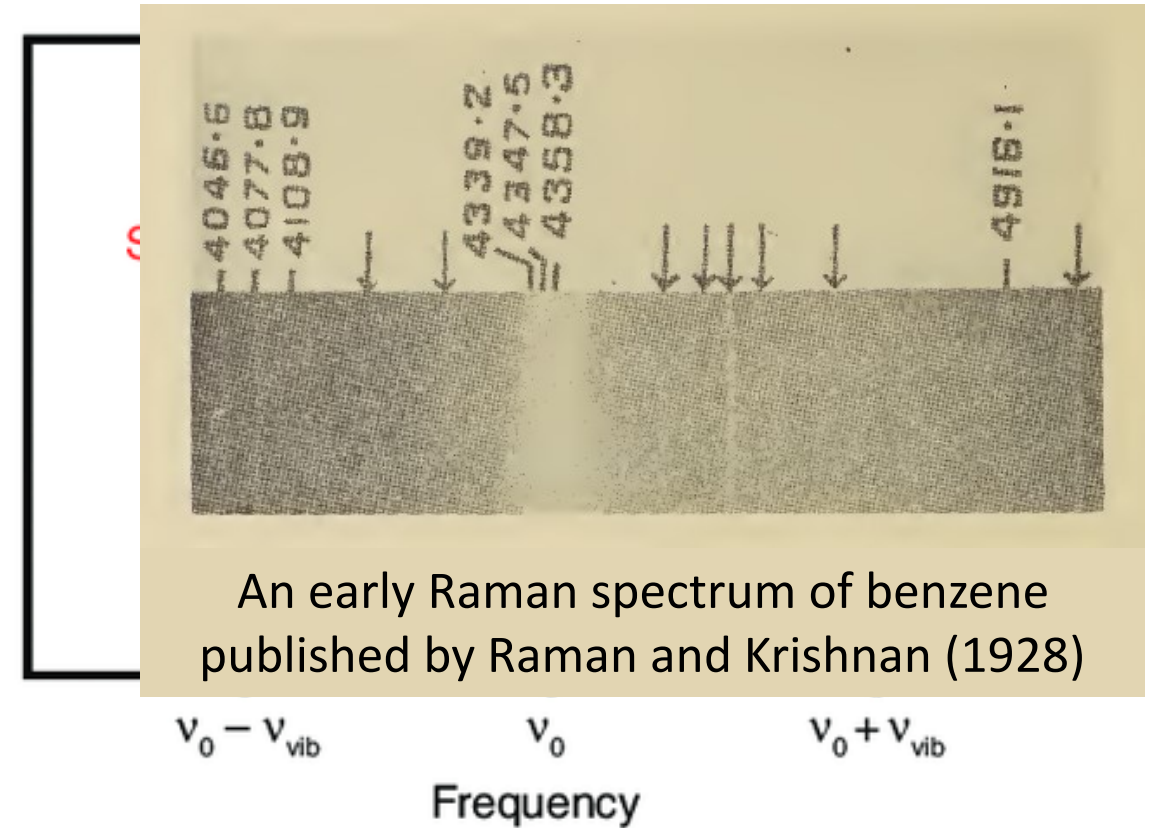
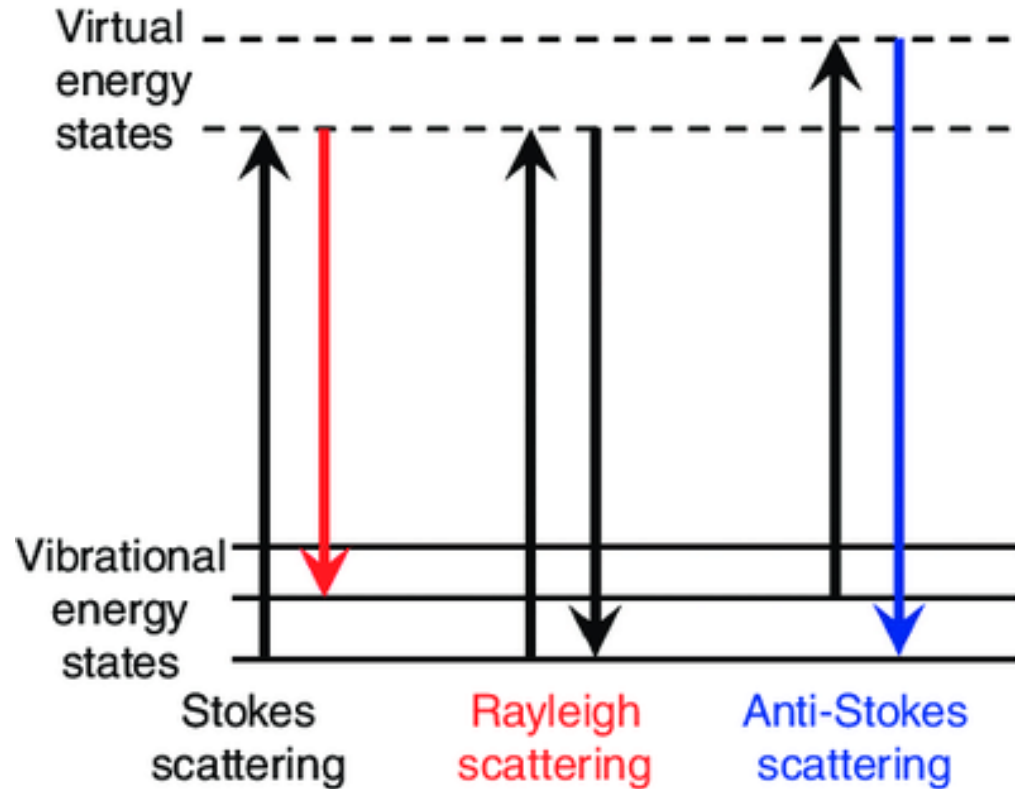




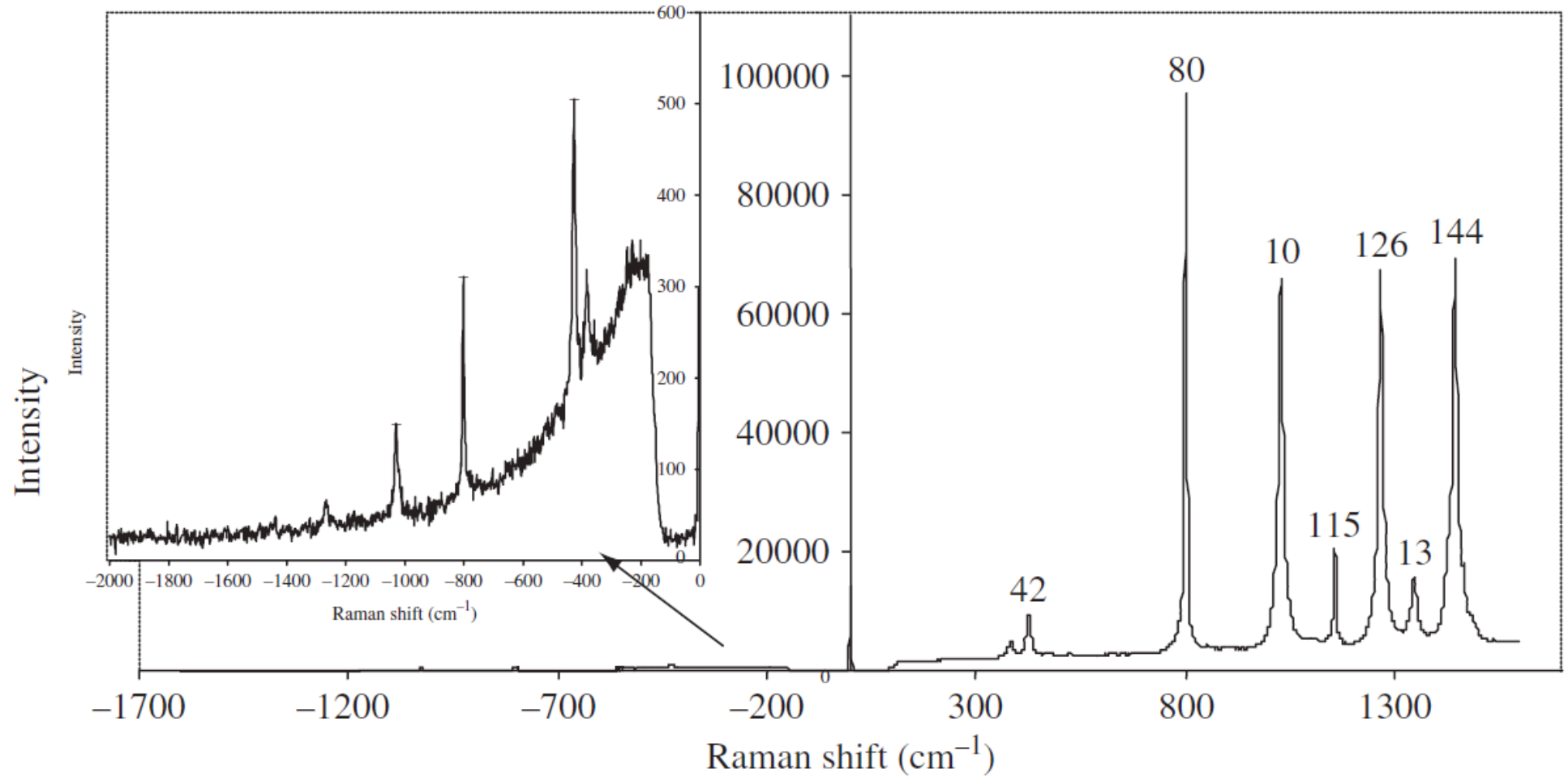
# Raman spectroscopy



# Raman spectroscopy

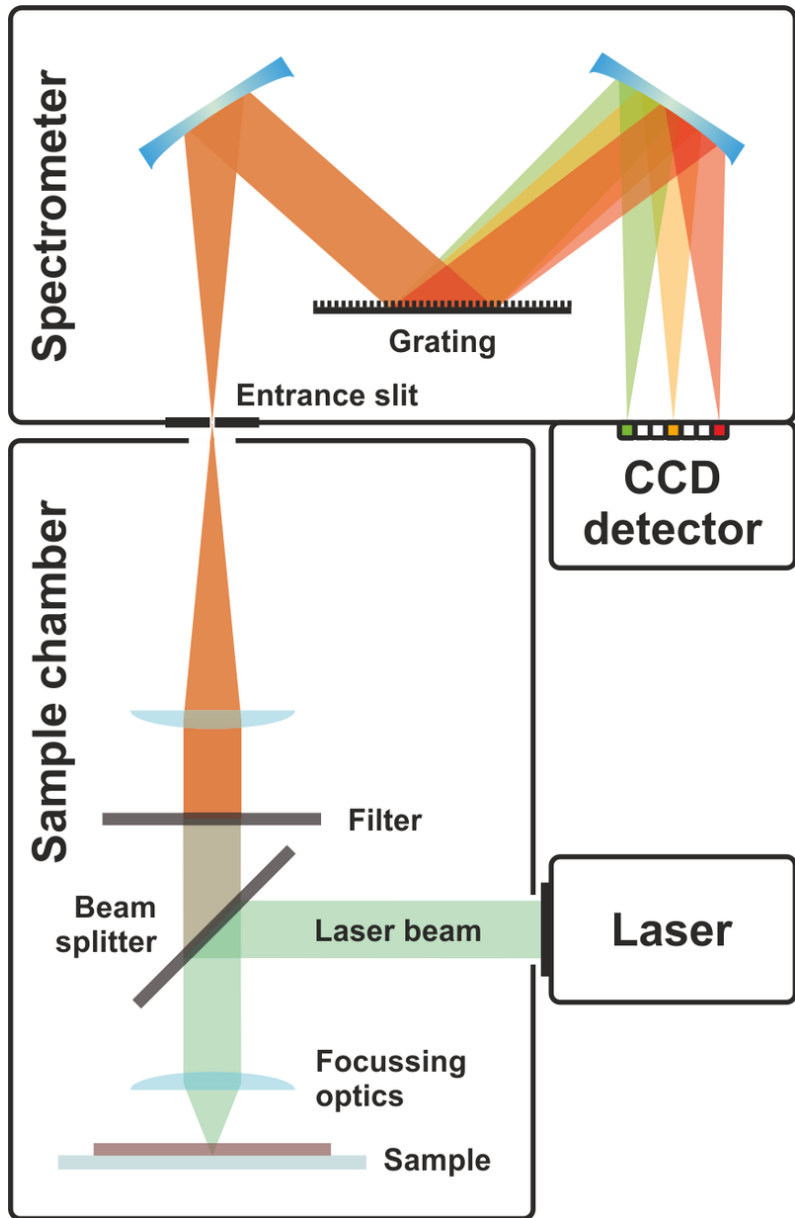


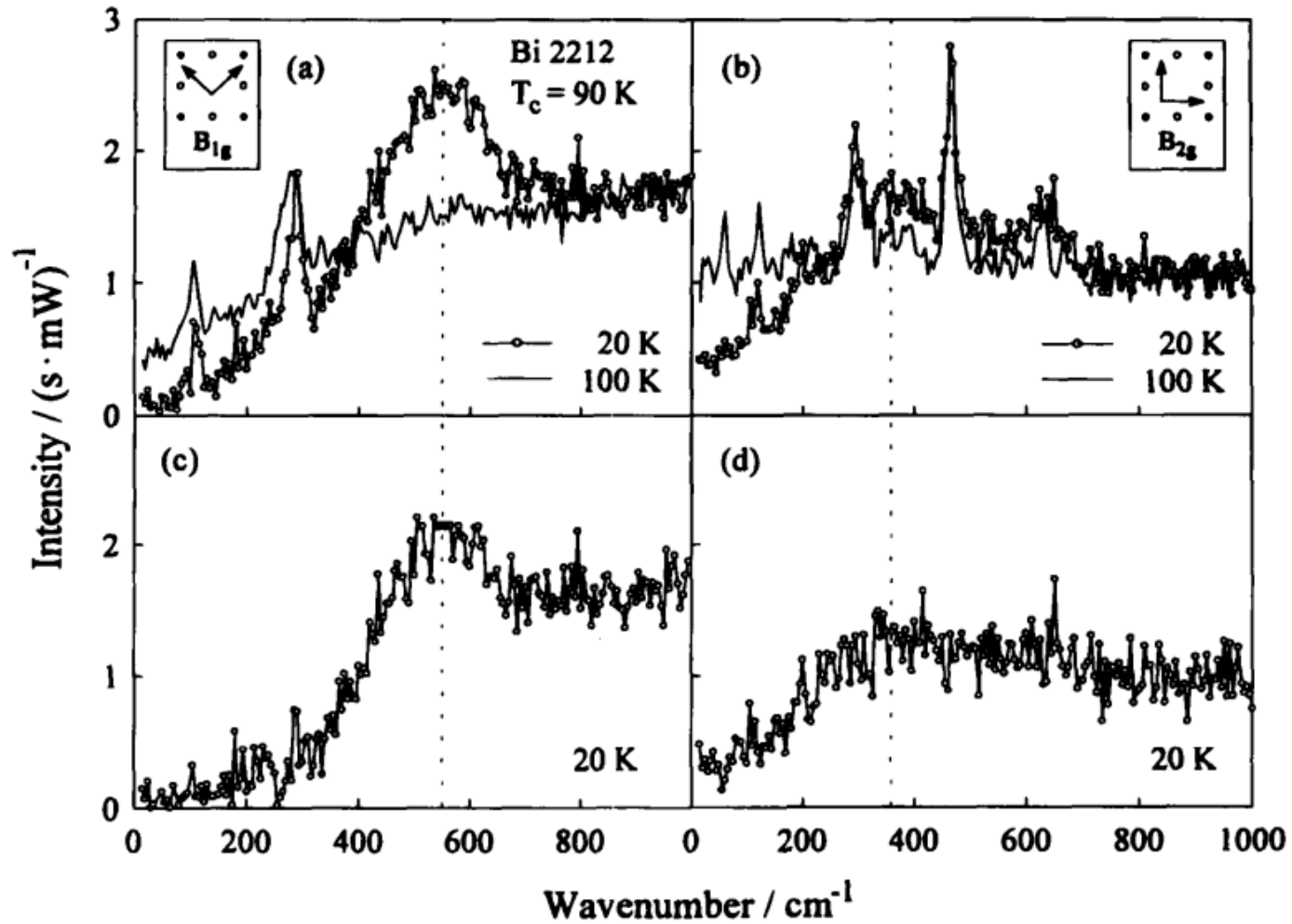
An early Raman spectrum of benzene published by Raman and Krishnan (1928)



Stokes and anti-Stokes scattering for cyclohexane. To show the weak anti-Stokes spectrum, the y-axis has been extended in the inset.

# Raman spectrometer





(a, b) Normal and superconducting spectra (raw data) of Bi<sub>2</sub>2212 for  $B_{1g}$  and  $B_{2g}$  symmetry. (c, d) Superconducting spectra after subtraction of the phonons.

Resonant inelastic X-ray scattering

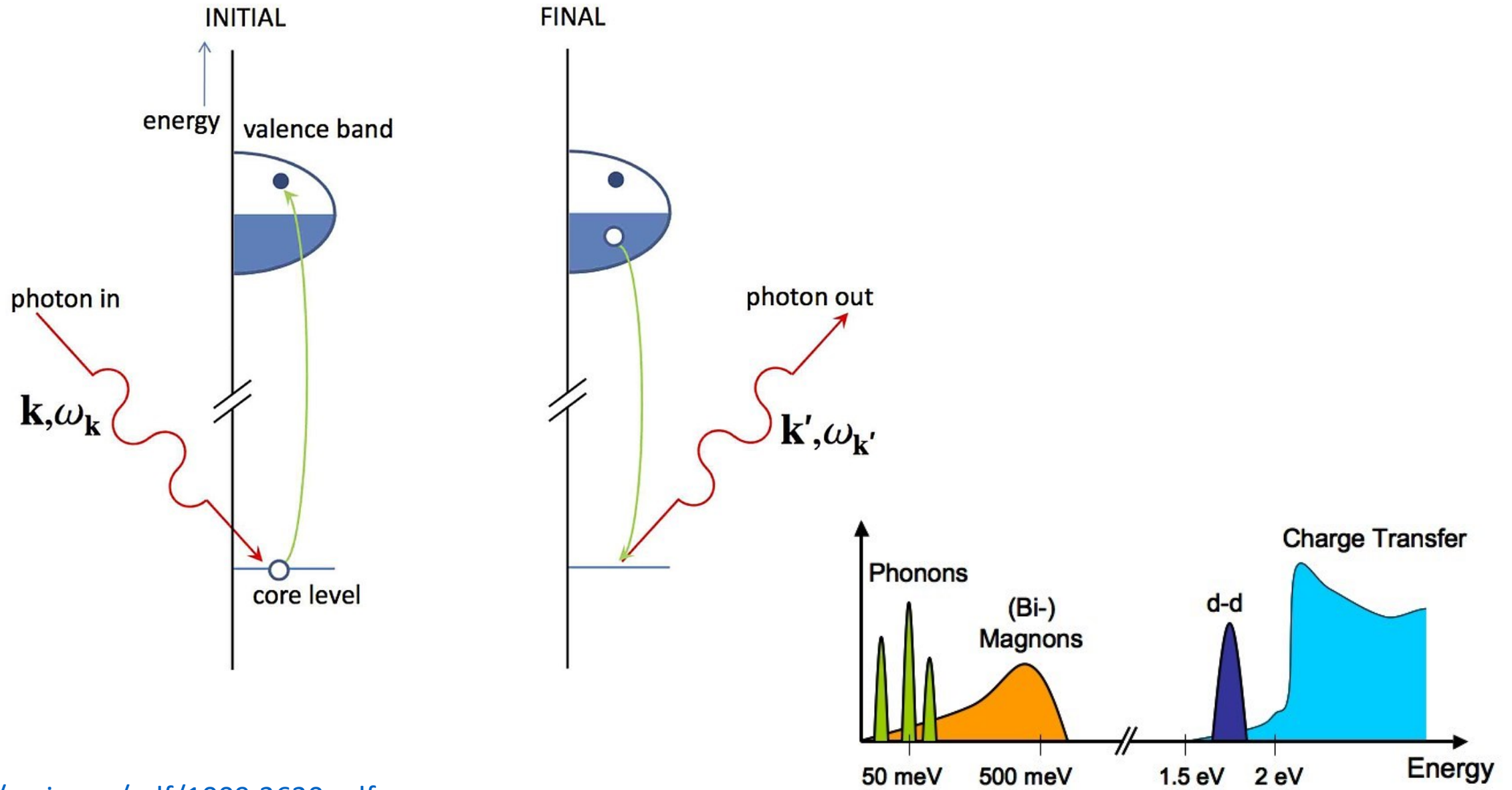
**RIXS**

= Resonant X-ray Raman

## *To read*

1. L. J. P. Ament et al. Resonant Inelastic X-ray Scattering Studies of Elementary Excitations. *Rev. Mod. Phys.* **83**, 705 (2011) <https://arxiv.org/pdf/1009.3630.pdf>

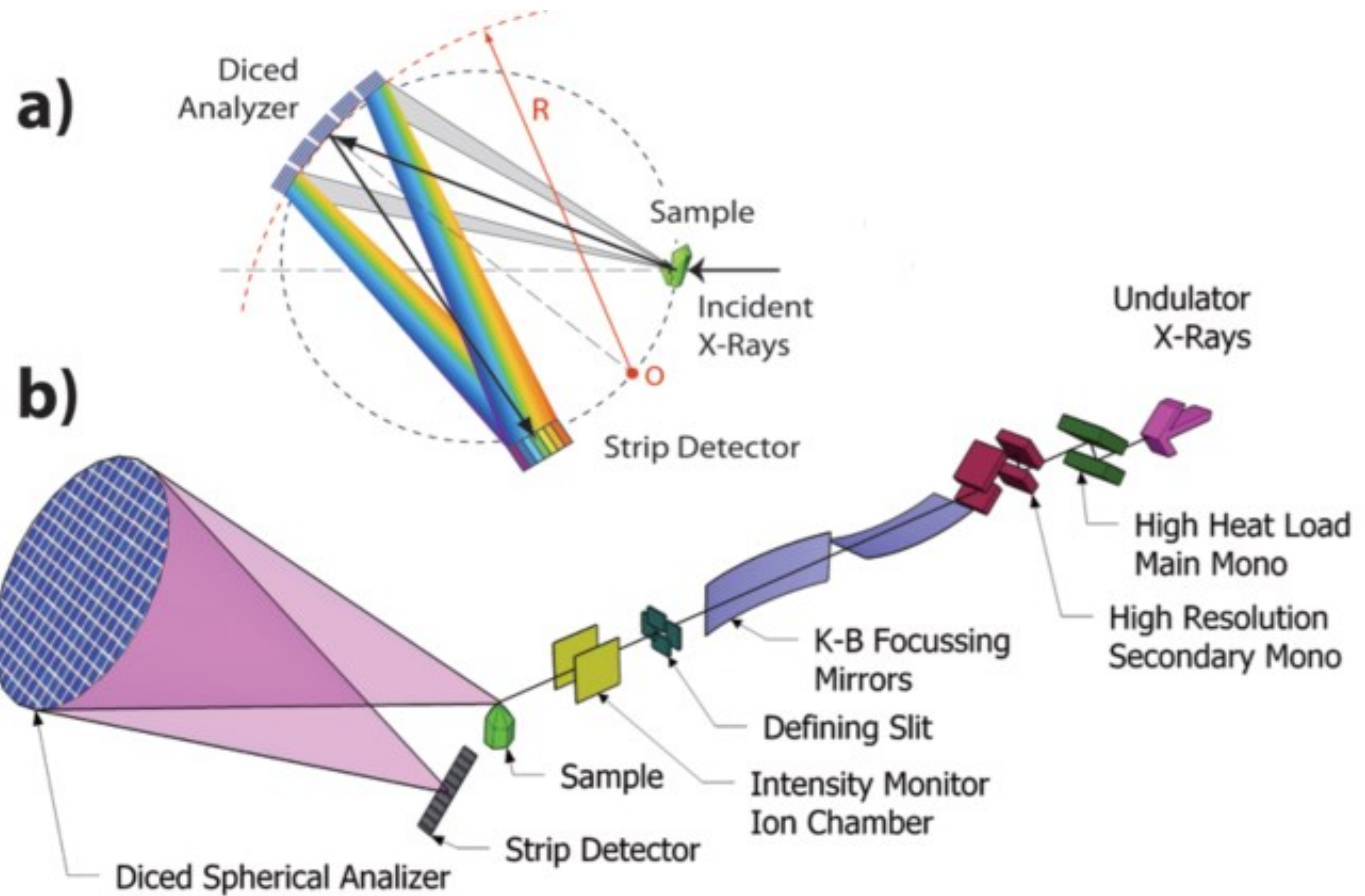
# RIXS

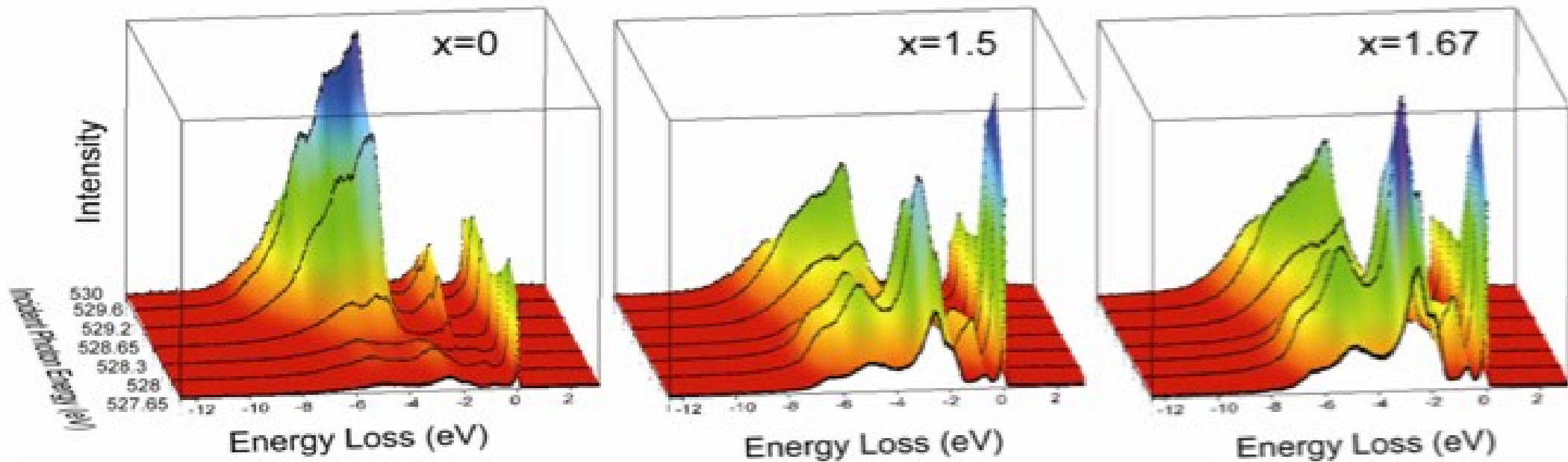






## Resonant Inelastic X-ray Scattering

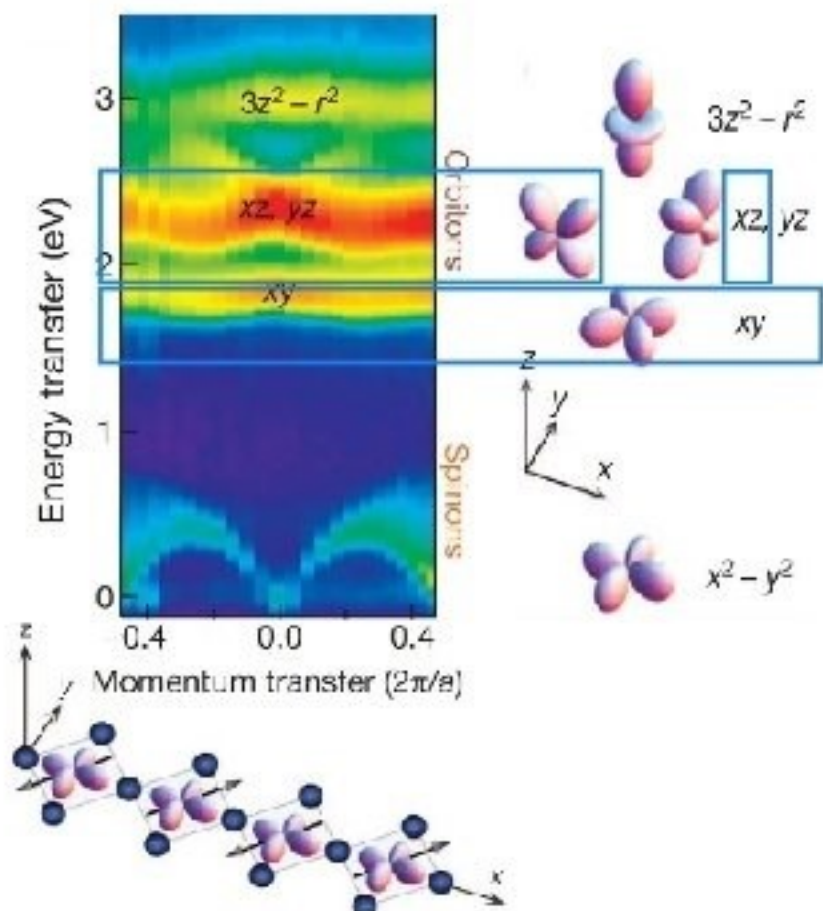




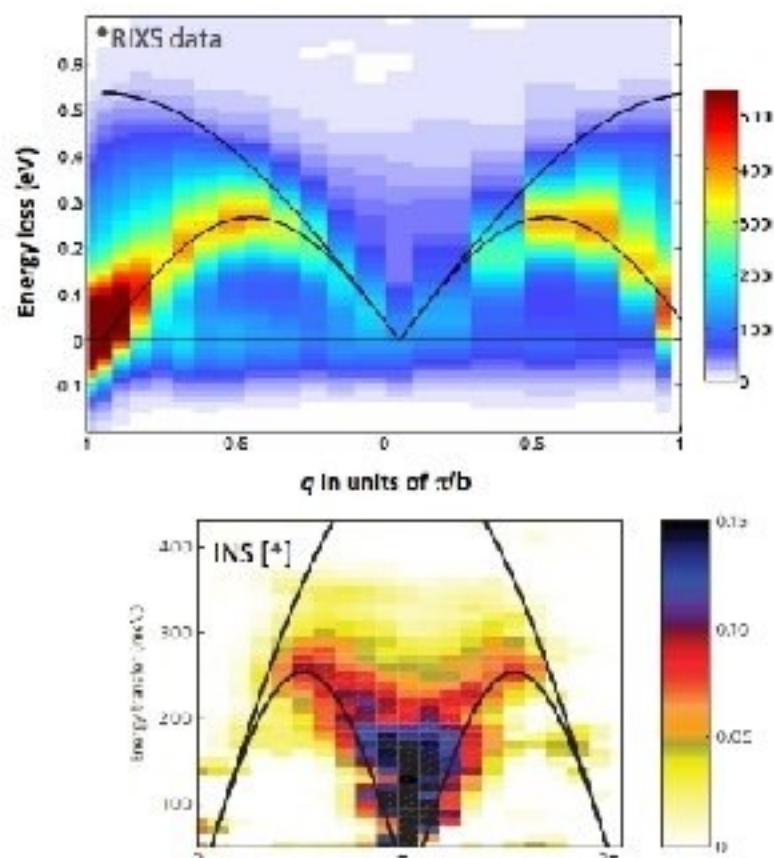
<https://arpes.stanford.edu/research/tool-development/resonant-x-ray-scattering>

# 2p3d RIXS in cuprates

$\text{Sr}_2\text{CuO}_3$



$\text{CaCu}_2\text{O}_3$



2p3d RIXS

Schlappa et al., Nature (2012)

Bisogni et al., [arXiv:1310.8346](https://arxiv.org/abs/1310.8346)