



Курс:

Фізичні методи дослідження матеріалів

Тема:

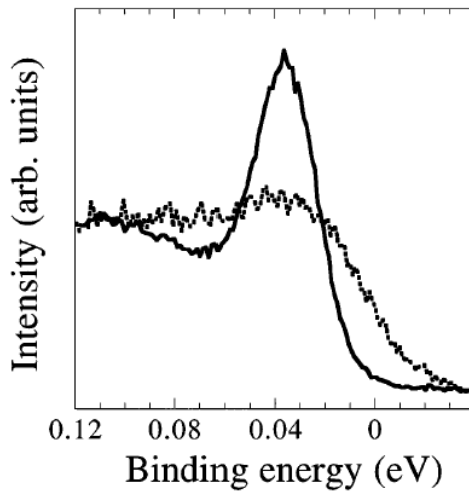
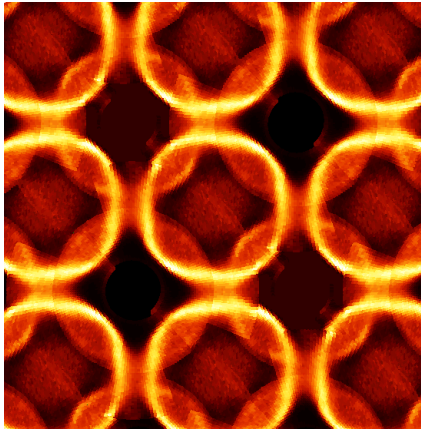
Спектральні функції та перетворення Фур'є

Лектор: О. А. Кордюк

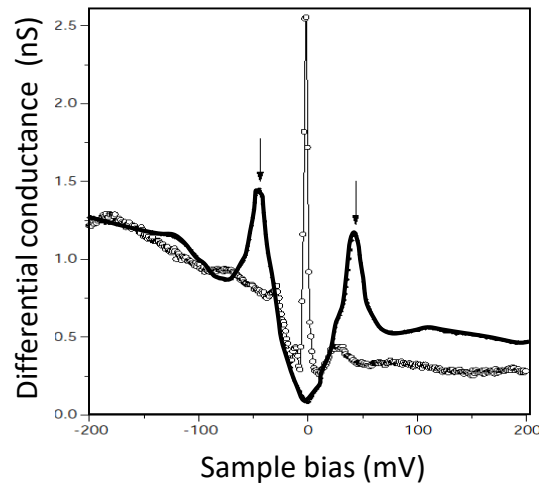
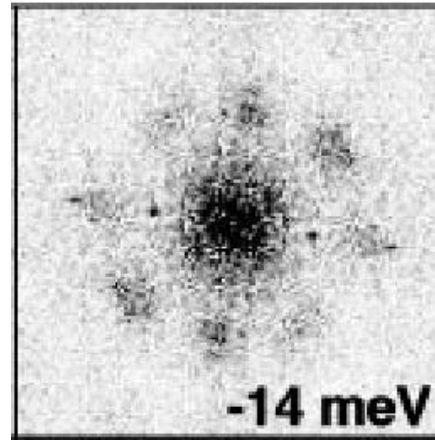
1 particle vs 2 particle spectra

Modern momentum resolving techniques

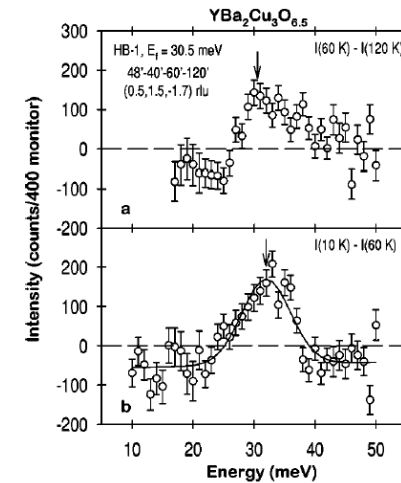
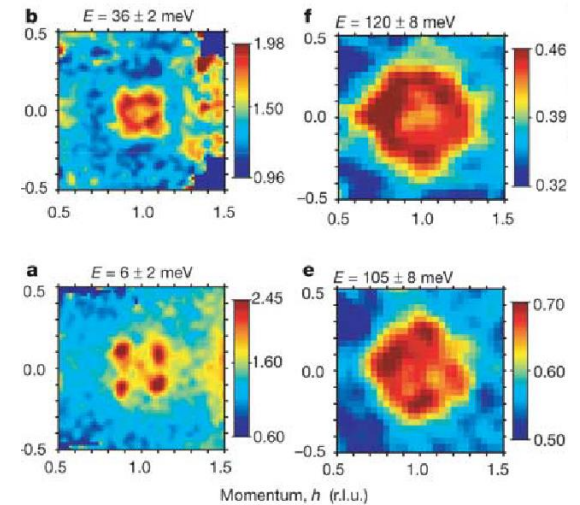
ARPES



STS



INS



Одно-частинкова спектральна функція = Im G

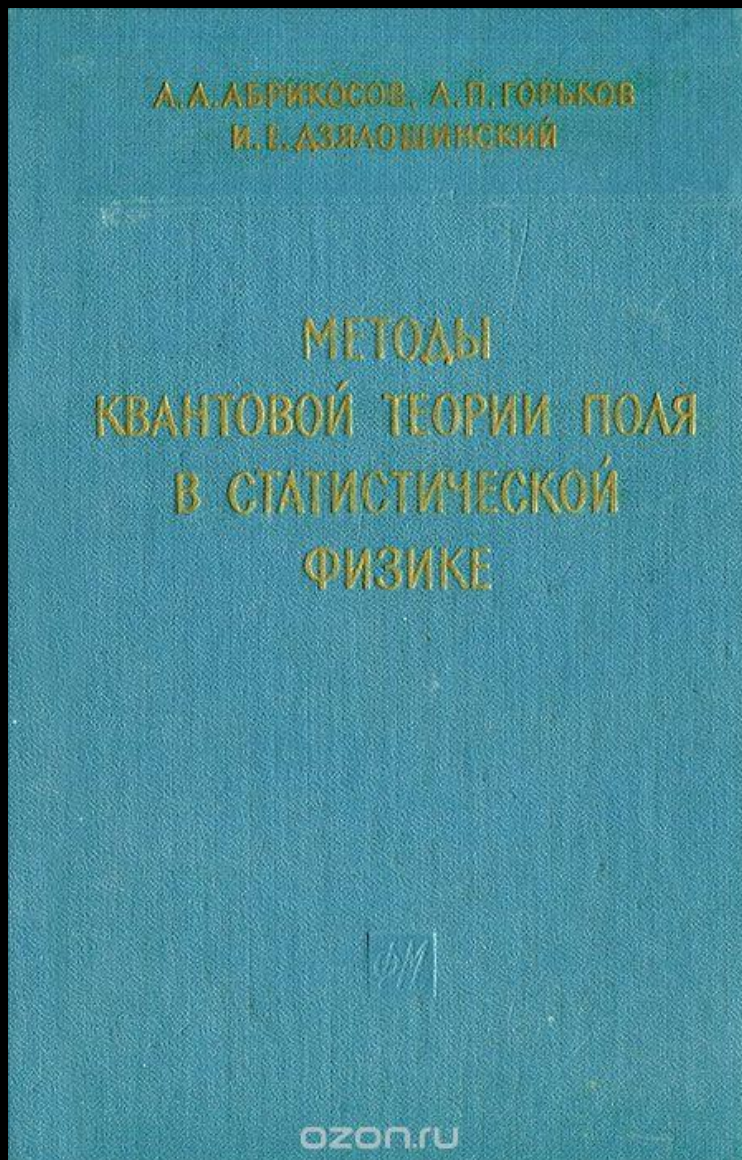
$$G_0(\omega, \mathbf{k}) = \frac{1}{\omega - \varepsilon(\mathbf{k}) + i\delta}$$

$$G(\omega, \mathbf{k}) = \frac{1}{\omega - \varepsilon(\mathbf{k}) + \Sigma(\omega)}$$

$$A(\omega, \mathbf{k}) = \text{Im } G(\omega, \mathbf{k})$$

$$A(\omega, \mathbf{k}) = -\frac{1}{\pi} \frac{\Sigma''(\omega)}{(\omega - \varepsilon(\mathbf{k}) - \Sigma'(\omega))^2 + \Sigma''(\omega)^2}$$

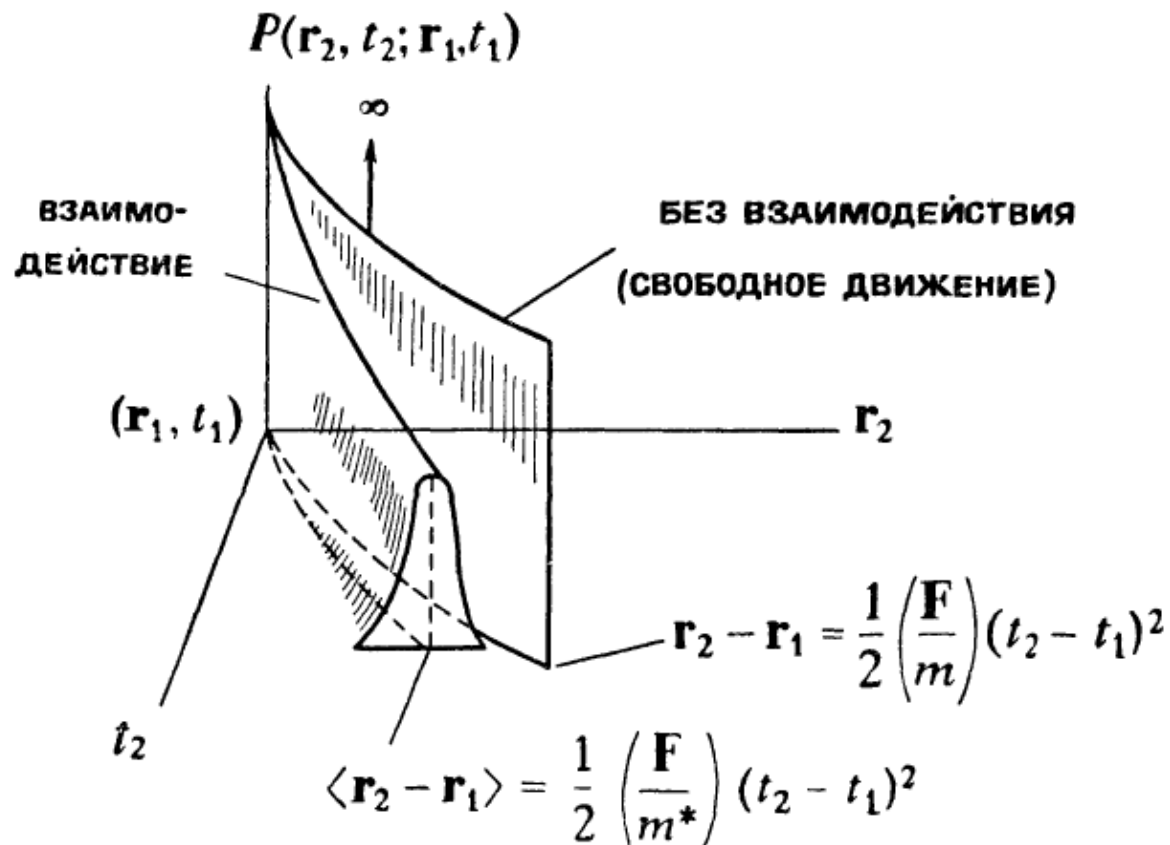
Спектральна функція = Im (функція Гріна)



$$\mathbf{r}_2 - \mathbf{r}_1 = \frac{1}{2} \left(\frac{\mathbf{F}}{m} \right) (t_2 - t_1)^2$$

$$P_0(\mathbf{r}_2, t_2; \mathbf{r}_1, t_1) = \delta \left[(\mathbf{r}_2 - \mathbf{r}_1) - \frac{1}{2} \left(\frac{\mathbf{F}}{m} \right) (t_2 - t_1)^2 \right]$$

$P(\mathbf{r}_2, t_2; \mathbf{r}_1, t_1)$ = Плотность вероятности (вероятность на единицу объема) того, что если частица в момент t_1 помещена в данную систему в точку \mathbf{r}_1 то она будет найдена в точке \mathbf{r}_2 в более поздний момент времени t_2 .



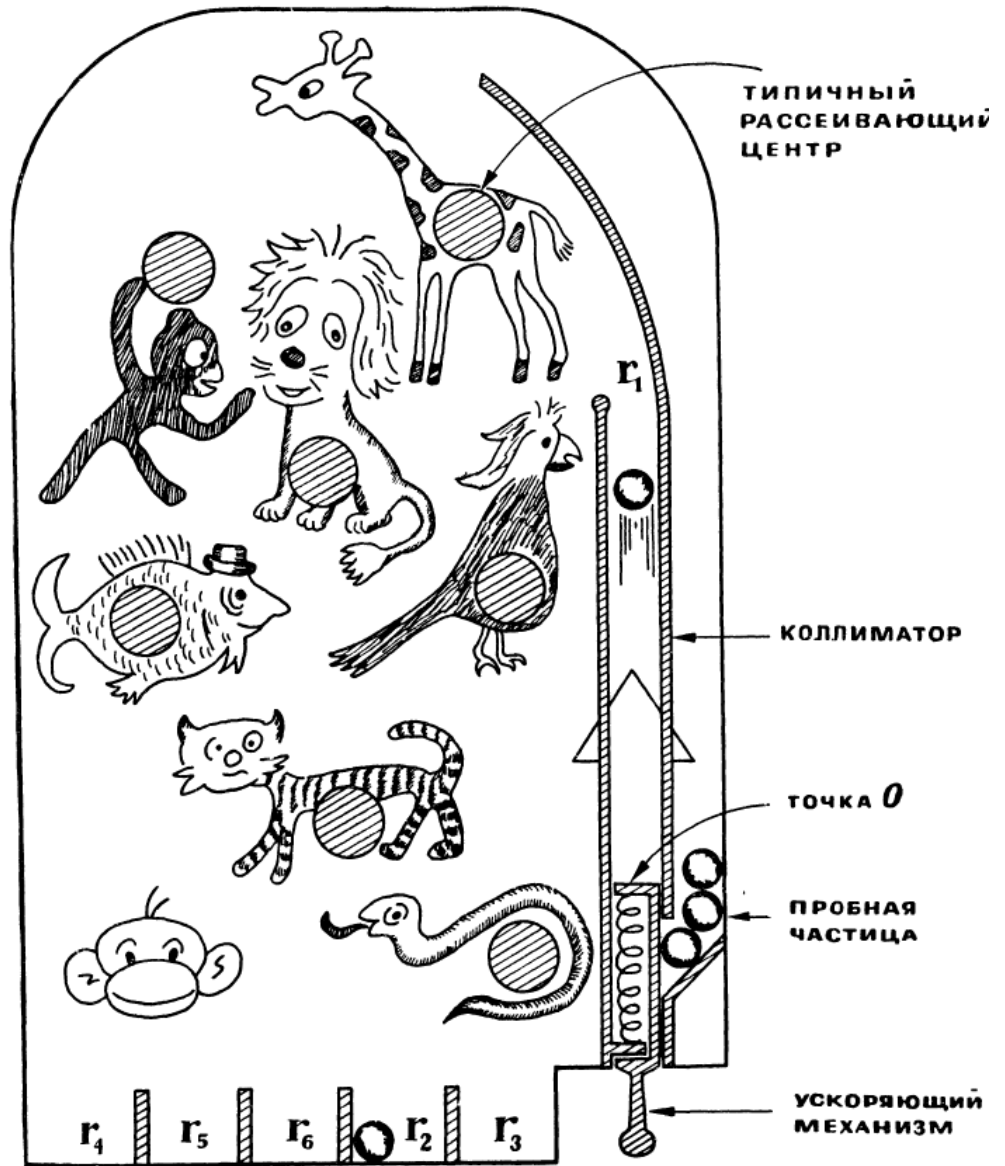
$$\langle \mathbf{r}_2 - \mathbf{r}_1 \rangle = \frac{1}{2} \left(\frac{\mathbf{F}}{m^*} \right) (t_2 - t_1)^2 \quad \text{для максимального значения } P$$

$$P_{\text{макс}}(\mathbf{r}_2, t_2; \mathbf{r}_1, t_1) \sim e^{-(t_2 - t_1)/\tau}$$

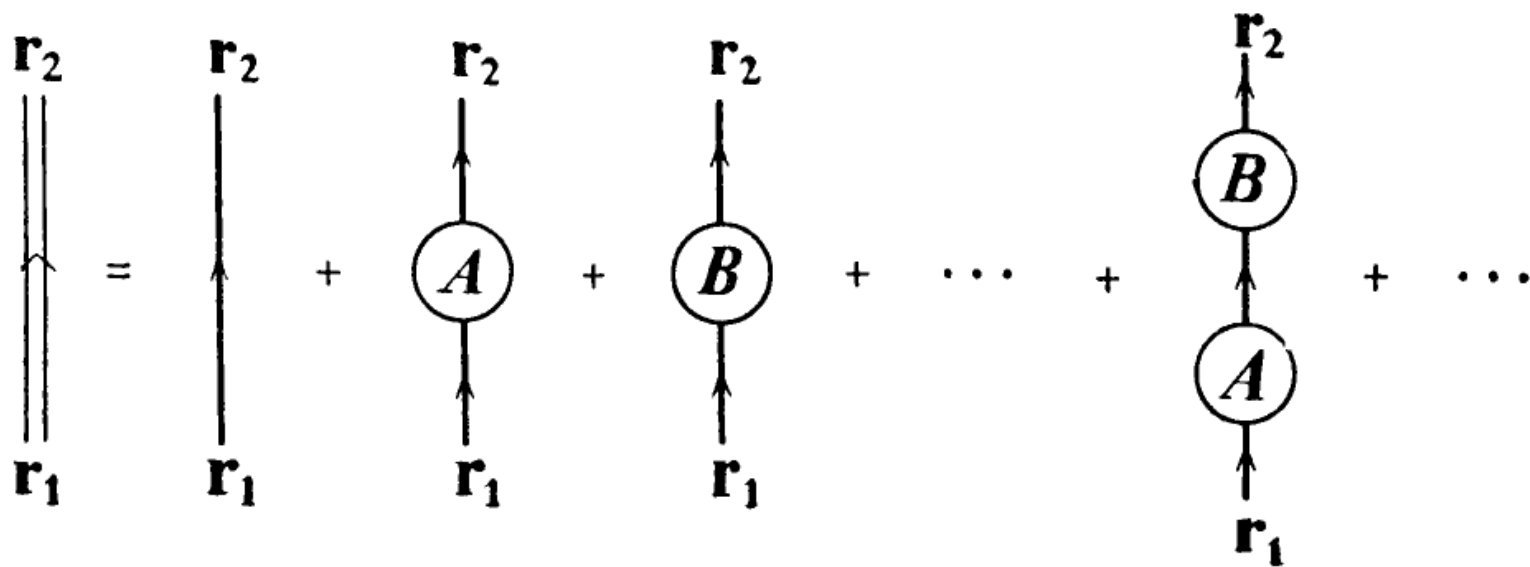
Р. Маттук



СЗОНТУ



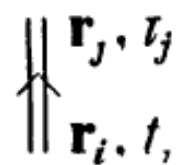
$$\begin{aligned}
 P(\mathbf{r}_2, \mathbf{r}_1) &= P_0(\mathbf{r}_2, \mathbf{r}_1) + P_0(\mathbf{r}_A, \mathbf{r}_1) P(A) P_0(\mathbf{r}_2, \mathbf{r}_A) + \\
 &\quad + P_0(\mathbf{r}_B, \mathbf{r}_1) P(B) P_0(\mathbf{r}_2, \mathbf{r}_B) + \dots \\
 &\quad + P_0(\mathbf{r}_A, \mathbf{r}_1) P(A) P_0(\mathbf{r}_B, \mathbf{r}_A) P(B) P_0(\mathbf{r}_2, \mathbf{r}_B) + \dots
 \end{aligned}$$



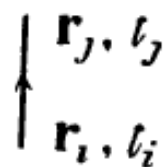
Слово

Диаграмма

$$P(\mathbf{r}_j, \mathbf{r}_i, t_j - t_i)$$

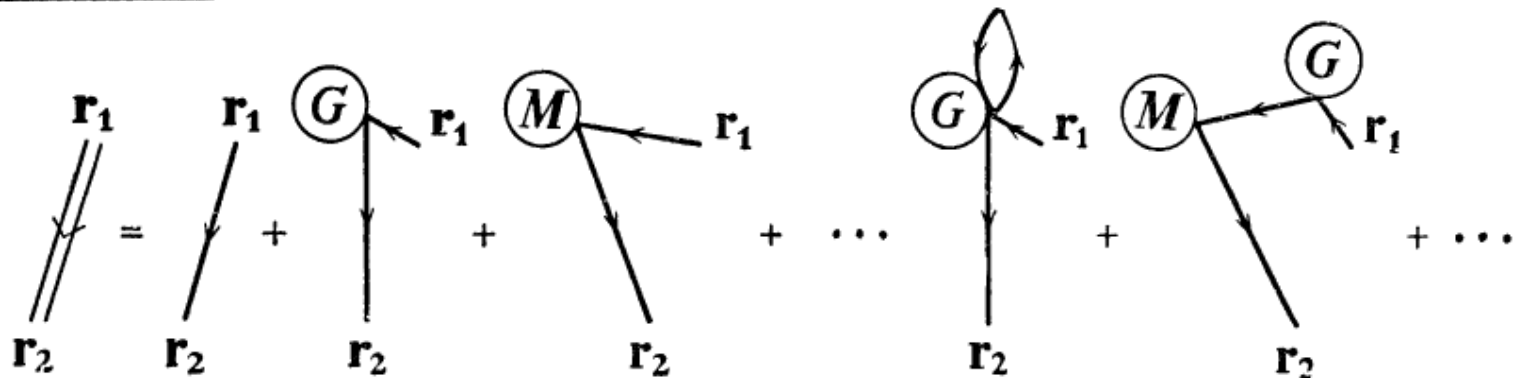
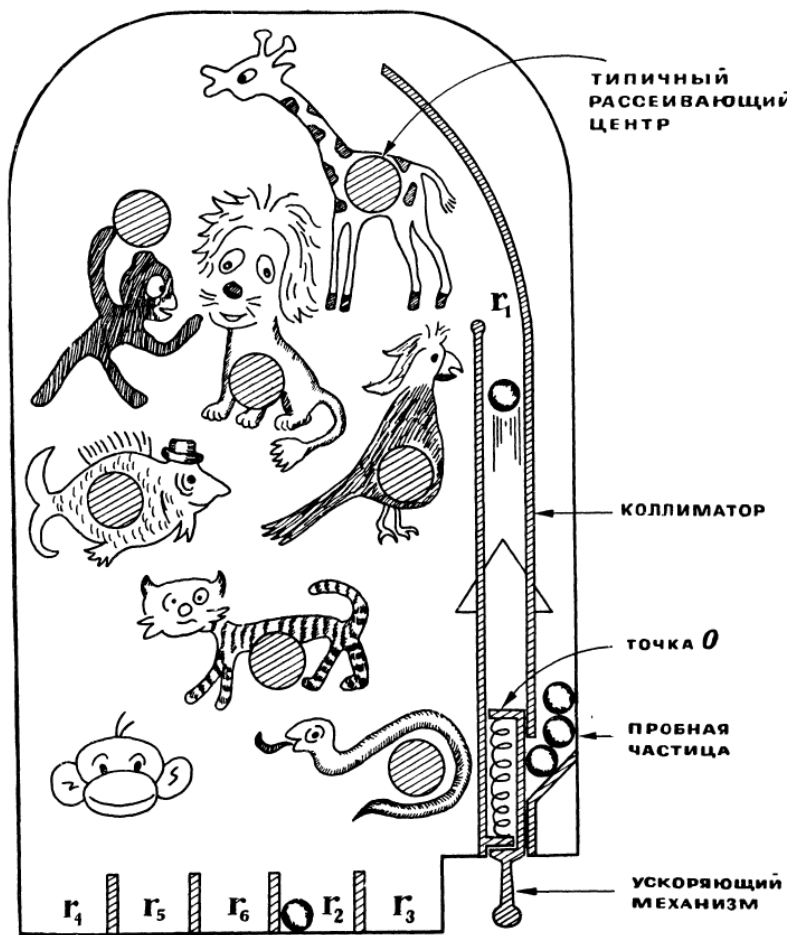


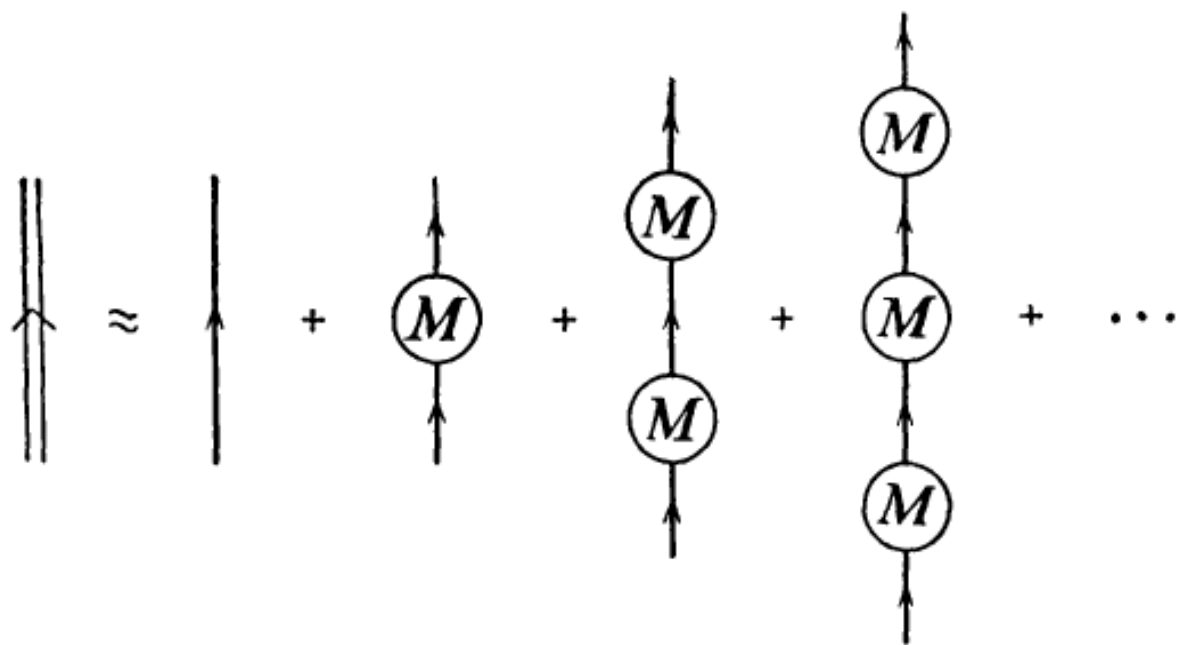
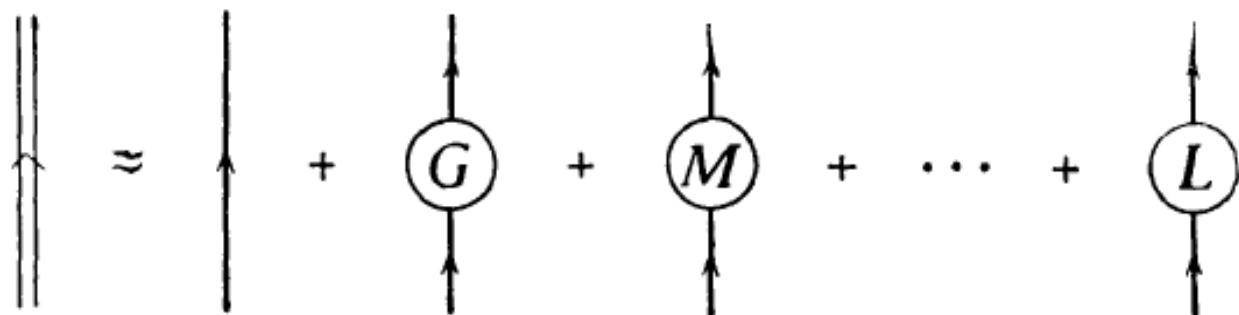
$$P_0(\mathbf{r}_j, \mathbf{r}_i, t_j - t_i)$$

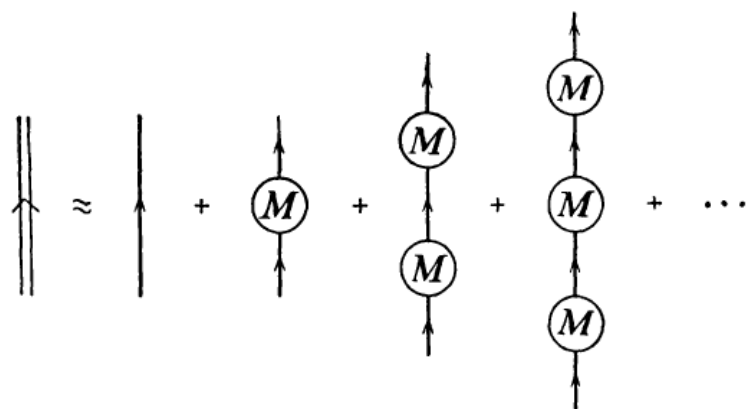


$$P(A)$$



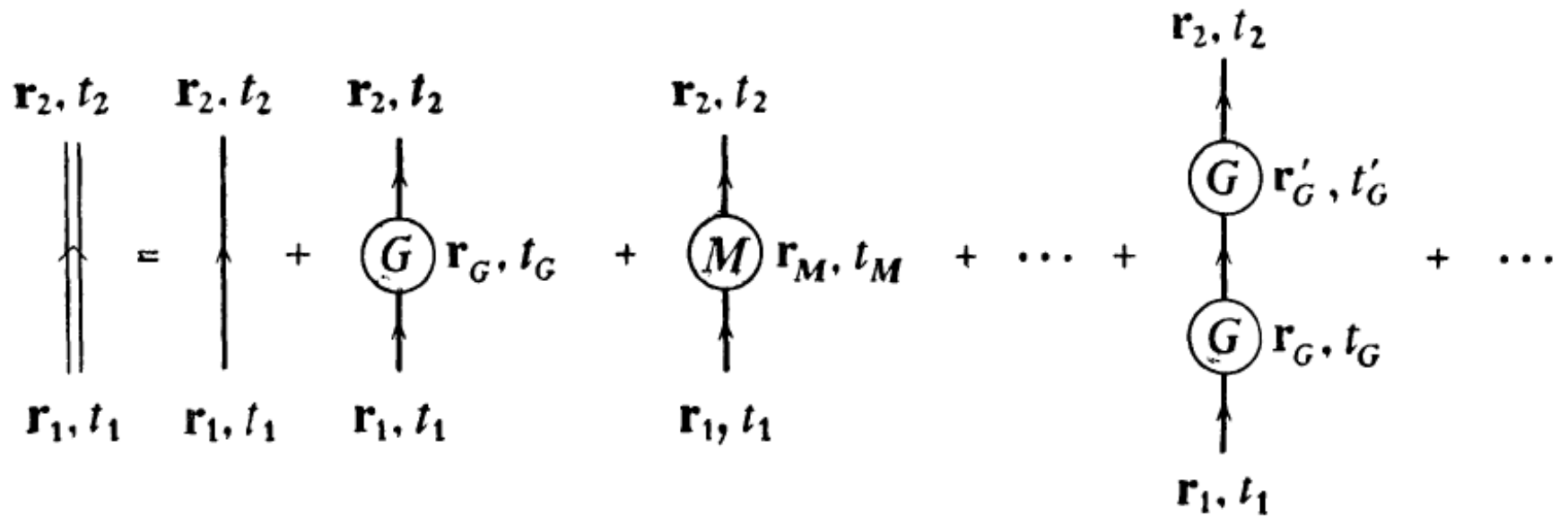






$$P(\mathbf{r}_2, \mathbf{r}_1) \approx P_0(\mathbf{r}_2, \mathbf{r}_1) + P_0(\mathbf{r}_M, \mathbf{r}_1) P(M) P_0(\mathbf{r}_2, \mathbf{r}_M) + \\ + P_0(\mathbf{r}_M, \mathbf{r}_1) P(M) P_0(\mathbf{r}_M, \mathbf{r}_M) P(M) P_0(\mathbf{r}_2, \mathbf{r}_M) + \dots$$

$$P(\mathbf{r}_2, \mathbf{r}_1) \approx P_0(\mathbf{r}_2, \mathbf{r}_1) + P_0(\mathbf{r}_M, \mathbf{r}_1) P(M) P_0(\mathbf{r}_2, \mathbf{r}_M) \times \\ \times [1 + P(M) P_0(\mathbf{r}_M, \mathbf{r}_M) + P(M)^2 P_0(\mathbf{r}_M, \mathbf{r}_M)^2 + \dots] = \\ = P_0(\mathbf{r}_2, \mathbf{r}_1) + \frac{P_0(\mathbf{r}_M, \mathbf{r}_1) P(M) P_0(\mathbf{r}_2, \mathbf{r}_M)}{1 - P(M) P_0(\mathbf{r}_M, \mathbf{r}_M)}$$



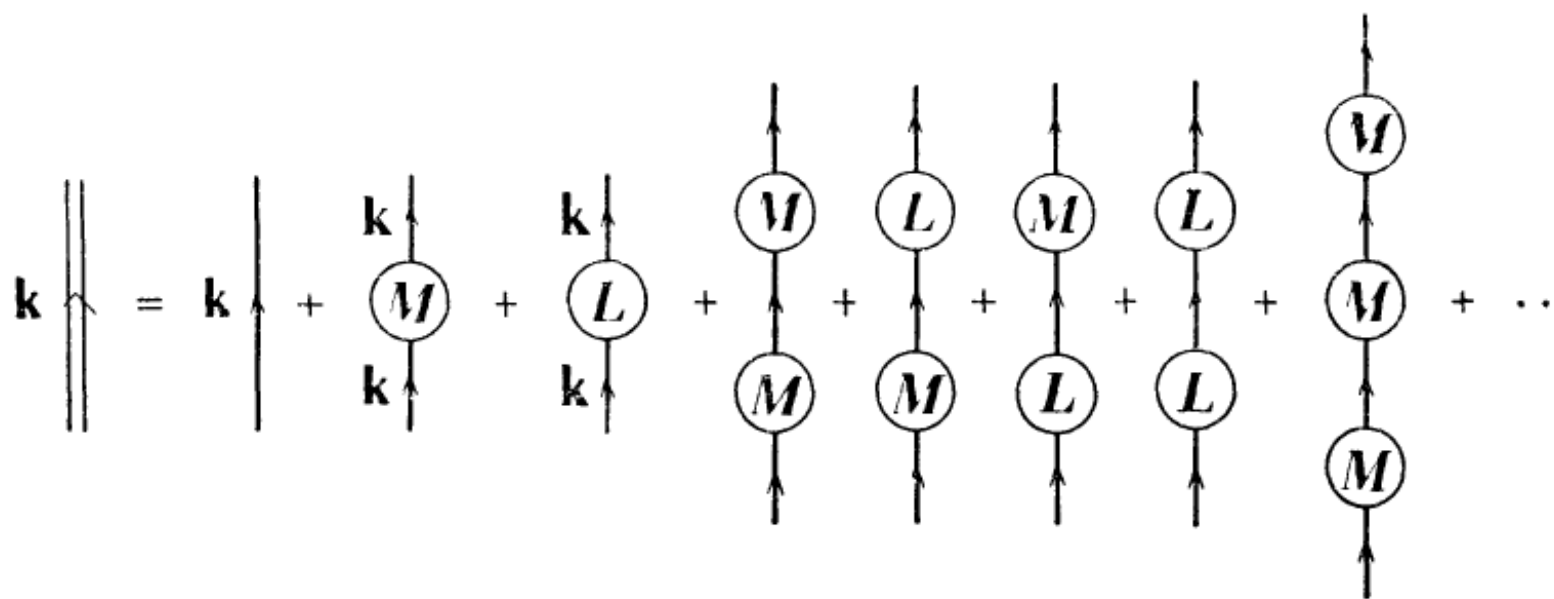
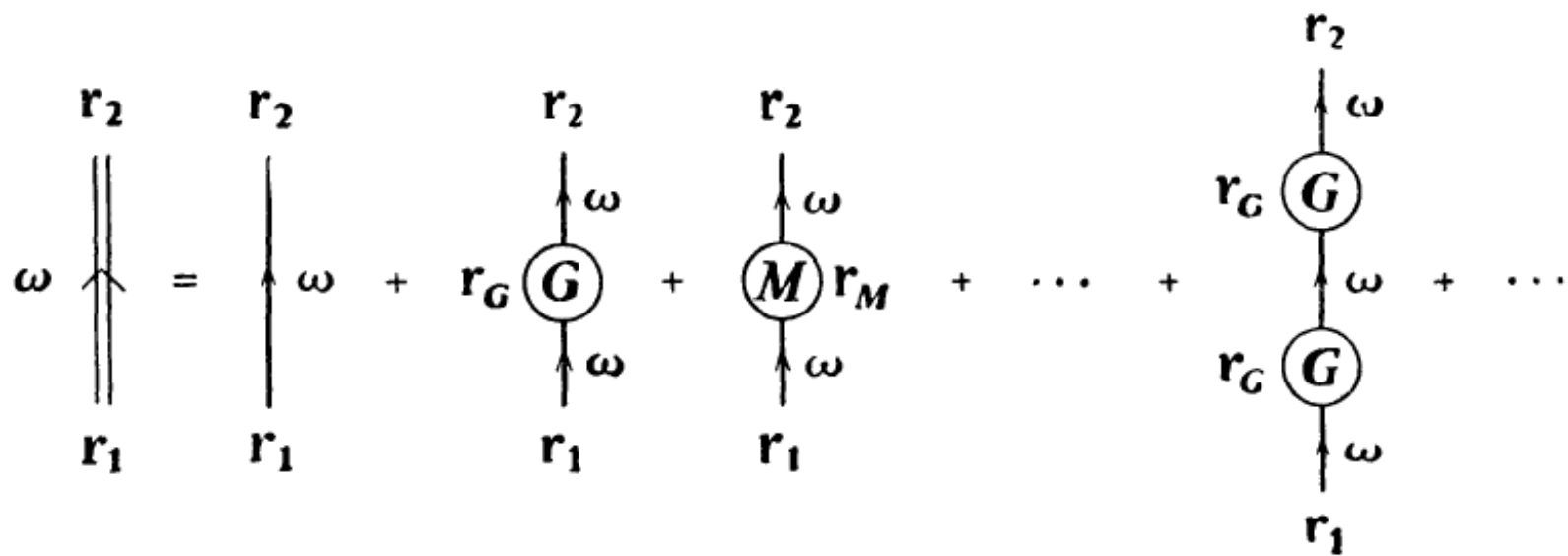
$$\begin{aligned}
 P(\mathbf{r}_2, \mathbf{r}_1, t_2 - t_1) &= P_0(\mathbf{r}_2, \mathbf{r}_1, t_2 - t_1) + \\
 &+ \int_{t_1}^{t_2} dt_G P_0(\mathbf{r}_G, \mathbf{r}_1, t_G - t_1) P(G) P_0(\mathbf{r}_2, \mathbf{r}_G, t_2 - t_G) + \\
 &+ \int dt_M \dots + \int \int + \dots + \int \int \int + \dots + \dots
 \end{aligned}$$

$$P_0(\mathbf{r}_j, \mathbf{r}_i, t_j - t_i) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} d\omega e^{-i\omega(t_j - t_i)} P_0(\mathbf{r}_j, \mathbf{r}_i, \omega)$$

$$\begin{aligned}
 P_0(\mathbf{r}_2, \mathbf{r}_1, t_2 - t_1) &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} d\omega e^{-i\omega(t_2 - t_1)} P_0(\mathbf{r}_2, \mathbf{r}_1, \omega), \\
 \int_{-\infty}^{+\infty} dt_G P_0(\mathbf{r}_G, \mathbf{r}_1, t_G - t_1) P(G) P_0(\mathbf{r}_2, \mathbf{r}_G, t_2 - t_G) &= \\
 &= \int_{-\infty}^{+\infty} dt_G \left[\frac{1}{2\pi} \int_{-\infty}^{+\infty} d\omega' e^{-i\omega'(t_G - t_1)} P_0(\mathbf{r}_G, \mathbf{r}_1, \omega') \right] \times \\
 &\times P(G) \left[\frac{1}{2\pi} \int_{-\infty}^{+\infty} d\omega e^{-i\omega(t_2 - t_G)} P_0(\mathbf{r}_2, \mathbf{r}_G, \omega) \right] = \\
 &= \frac{1}{(2\pi)^2} \int_{-\infty}^{+\infty} d\omega \int_{-\infty}^{+\infty} d\omega' P_0(\mathbf{r}_G, \mathbf{r}_1, \omega') \times \\
 &\times P(G) P_0(\mathbf{r}_2, \mathbf{r}_G, \omega) e^{+i(\omega' t_1 - \omega t_2)} \underbrace{\int_{-\infty}^{+\infty} dt_G e^{-it_G(\omega' - \omega)}}_{2\pi\delta(\omega' - \omega)} = \\
 &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} d\omega e^{-i\omega(t_2 - t_1)} P_0(\mathbf{r}_G, \mathbf{r}_1, \omega) P(G) P_0(\mathbf{r}_2, \mathbf{r}_G, \omega).
 \end{aligned}$$

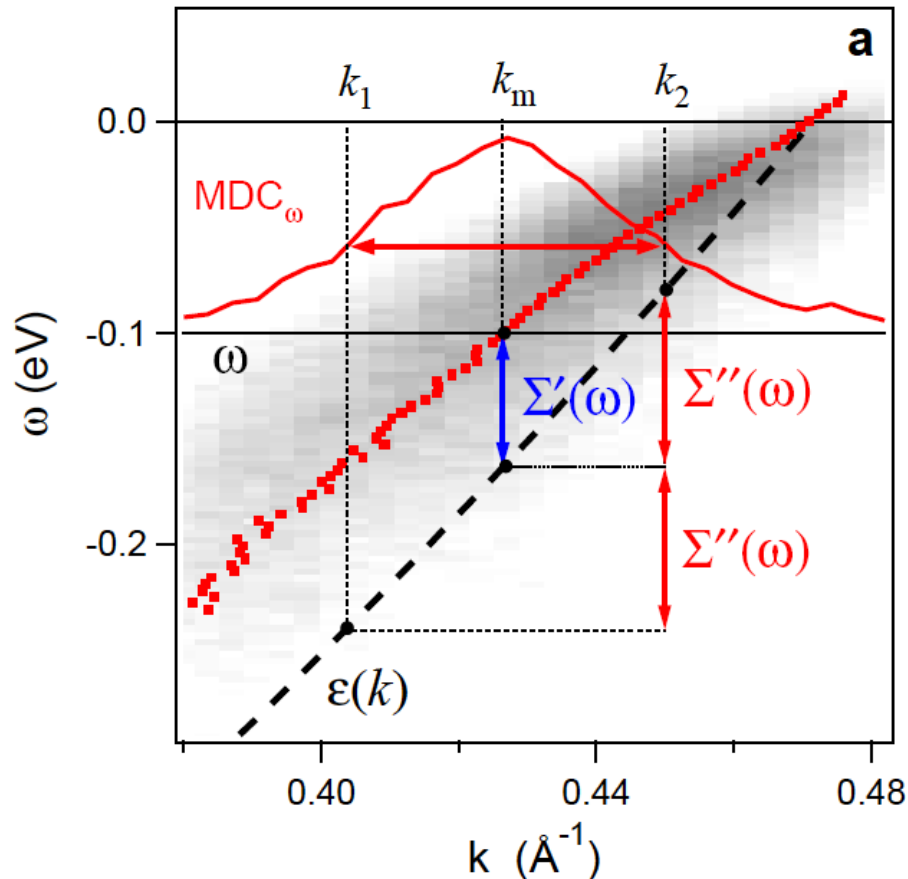
$$P(\mathbf{r}_2, \mathbf{r}_1, \omega) =$$

$$= P_0(\mathbf{r}_2, \mathbf{r}_1, \omega) + P_0(\mathbf{r}_G, \mathbf{r}_1, \omega) P(G) P_0(\mathbf{r}_2, \mathbf{r}_G, \omega) + \dots$$



Спектральна функція

$$A(\omega, \mathbf{k}) = -\frac{1}{\pi} \frac{\Sigma''(\omega)}{(\omega - \varepsilon(\mathbf{k}) - \Sigma'(\omega))^2 + \Sigma''(\omega)^2}$$



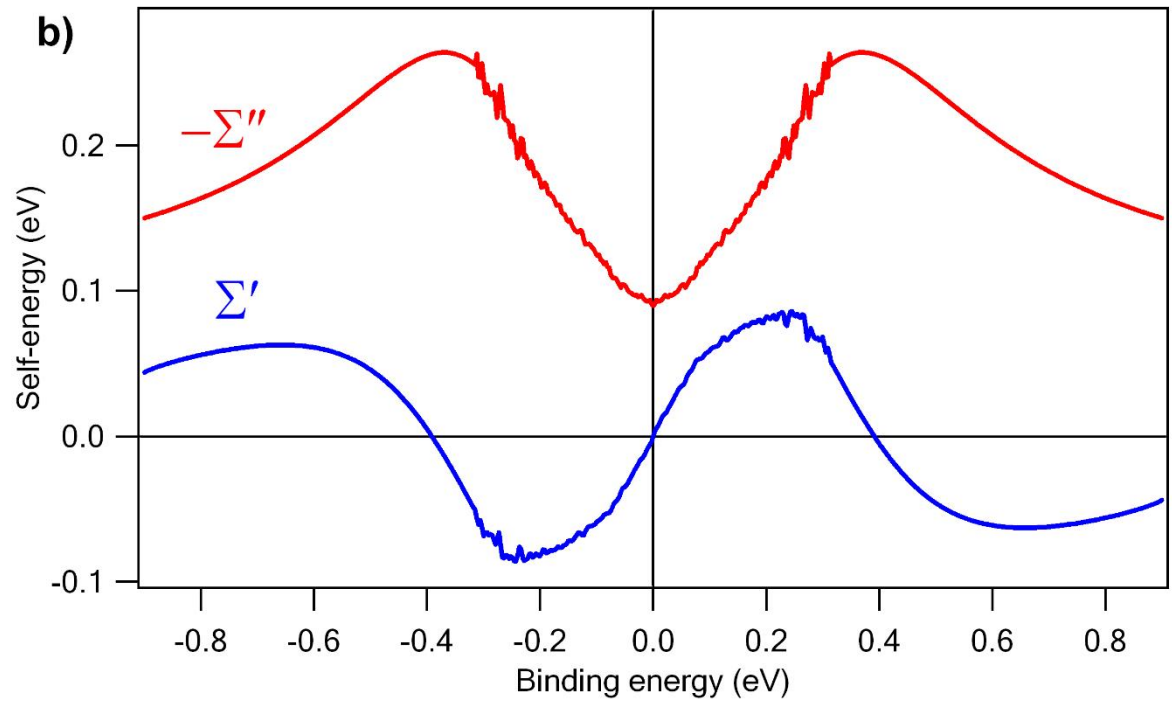
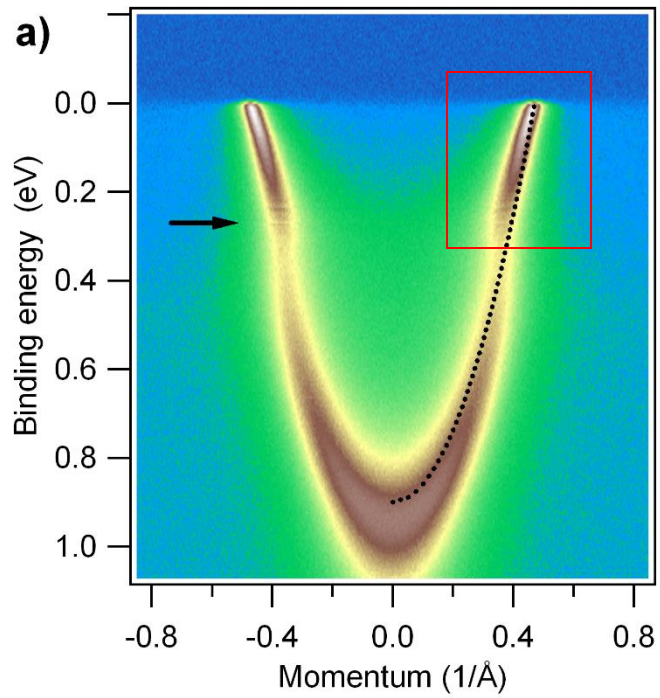
$$\text{EDC}(\omega) = A(\omega)_k$$

$$\text{MDC}(k) = A(k)_\omega$$

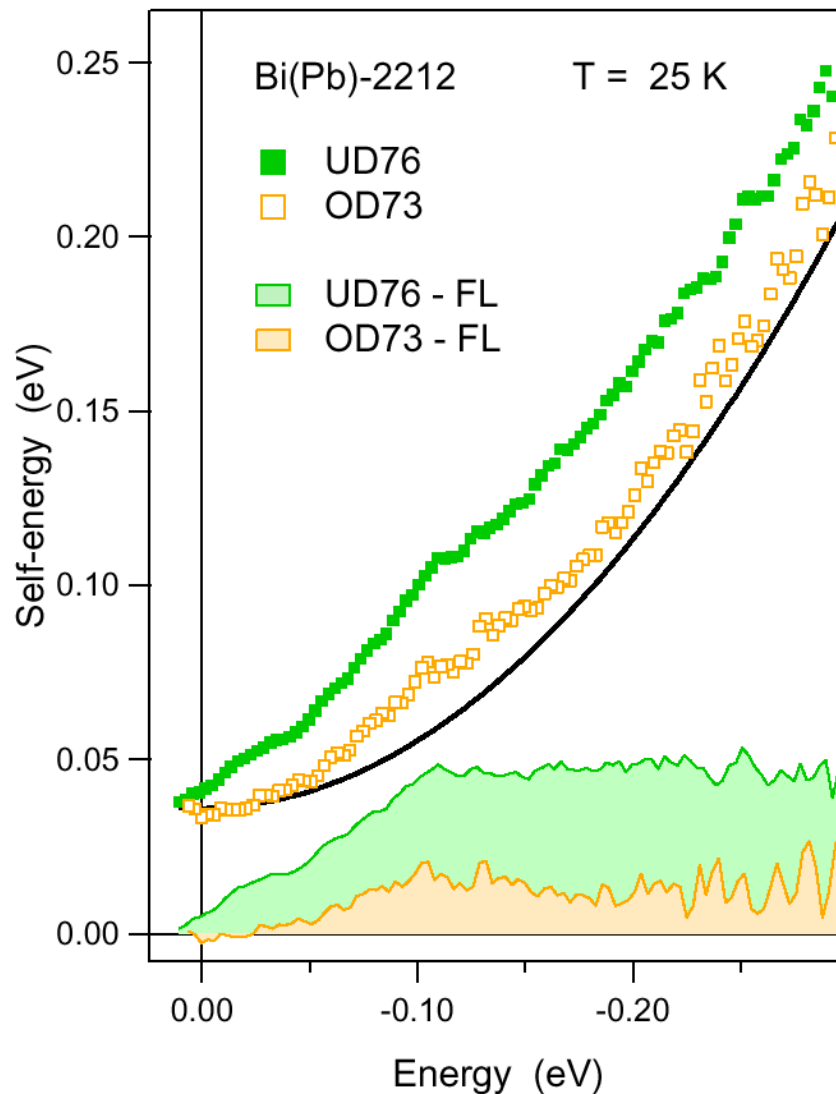
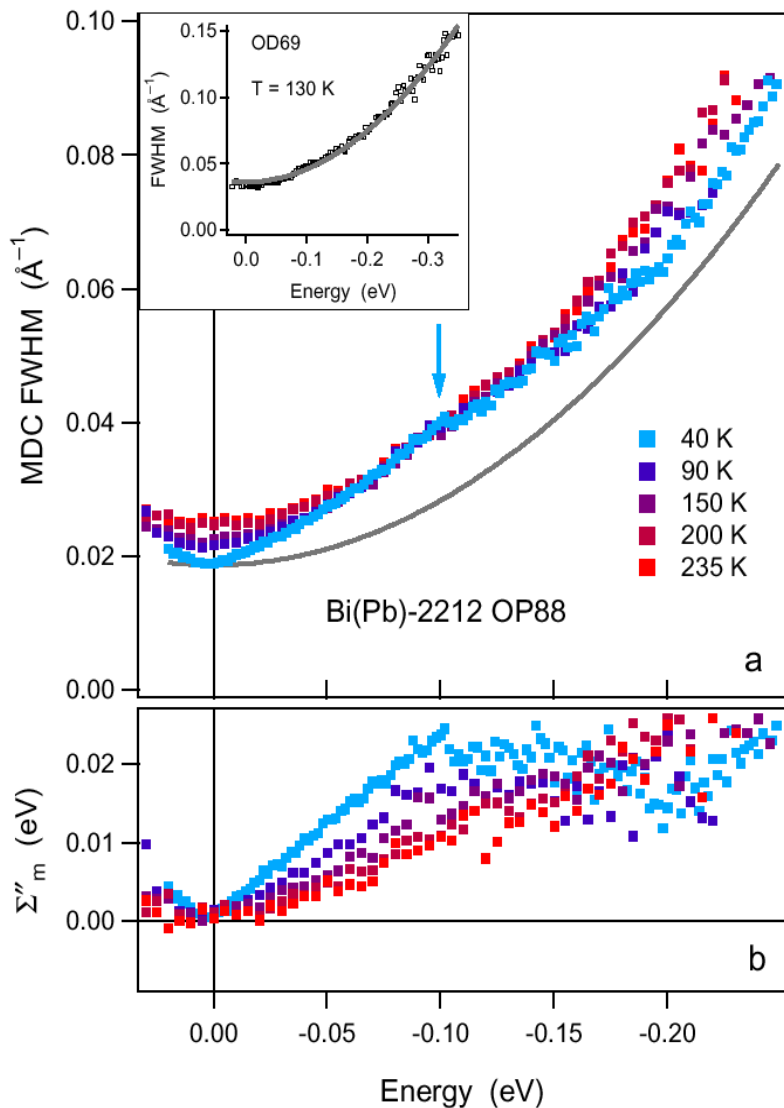
$$\Sigma'(\omega) = \omega - \varepsilon(k_m)$$

$$\Sigma''(\omega) = -v_F W(\omega)$$

KK

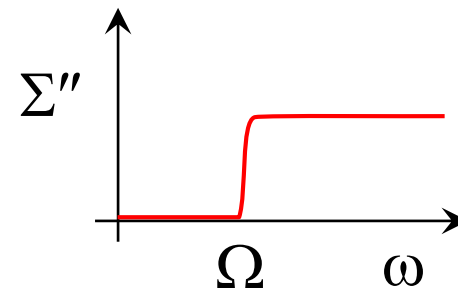
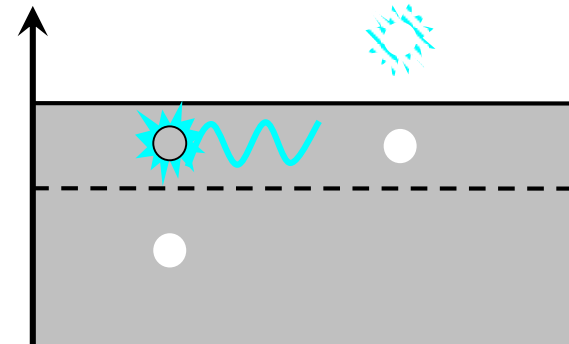
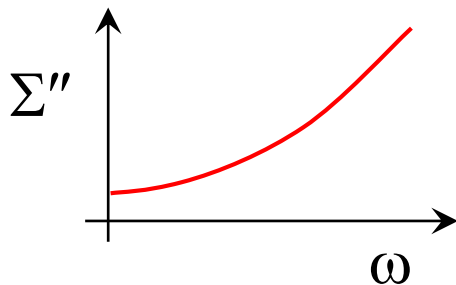
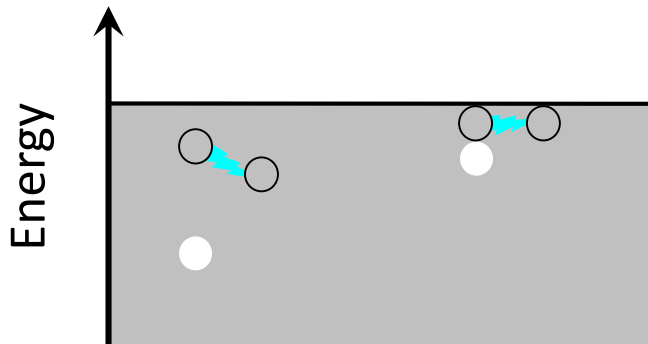
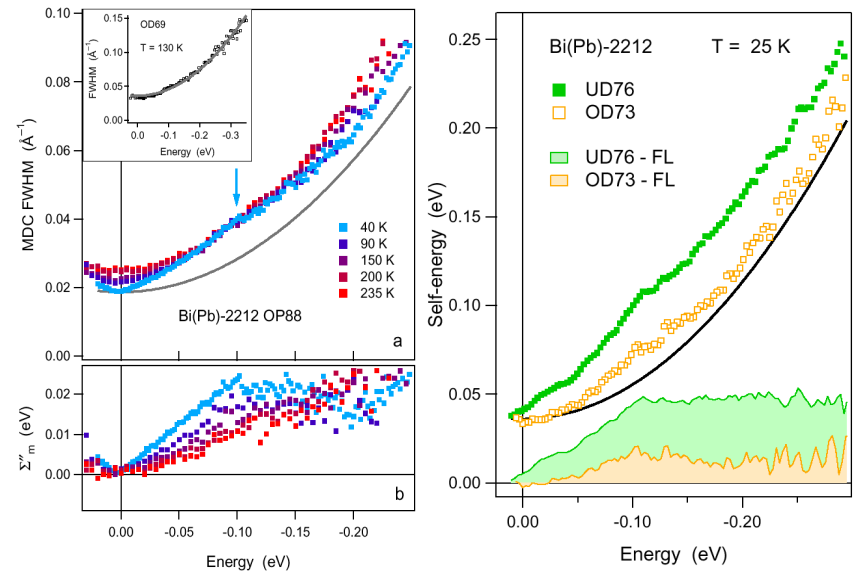


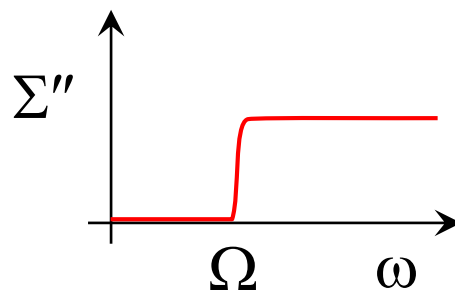
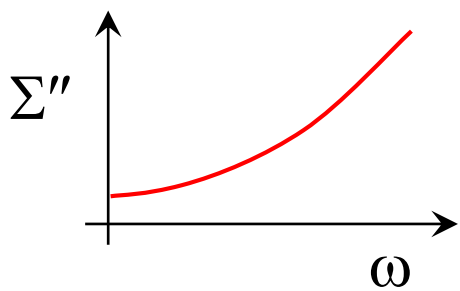
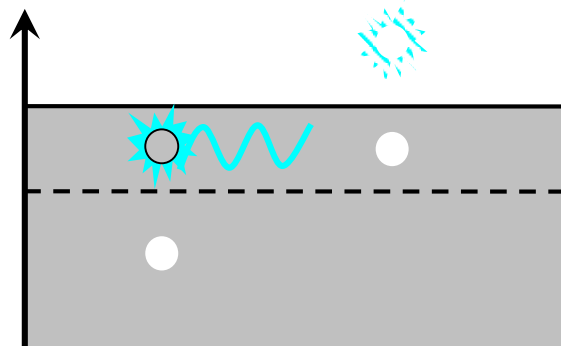
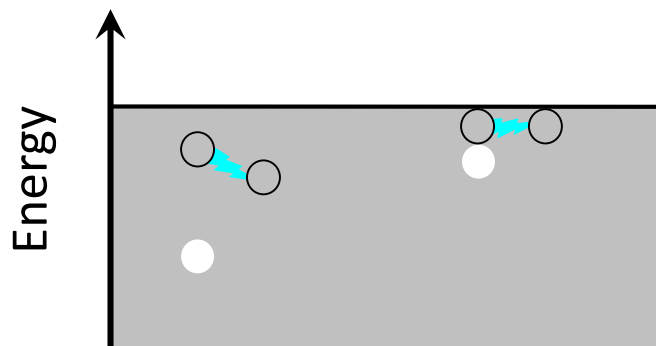
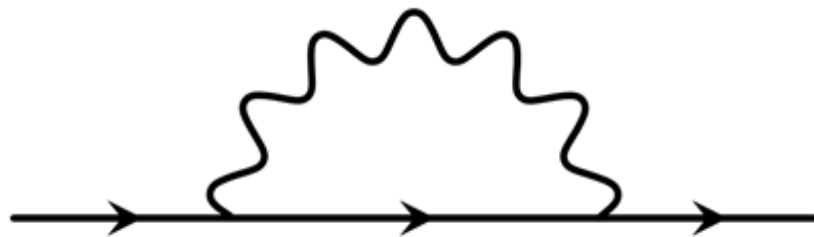
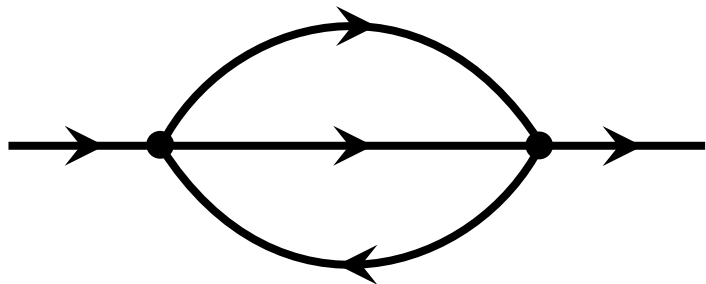
Scattering rate: T - and x -dependence



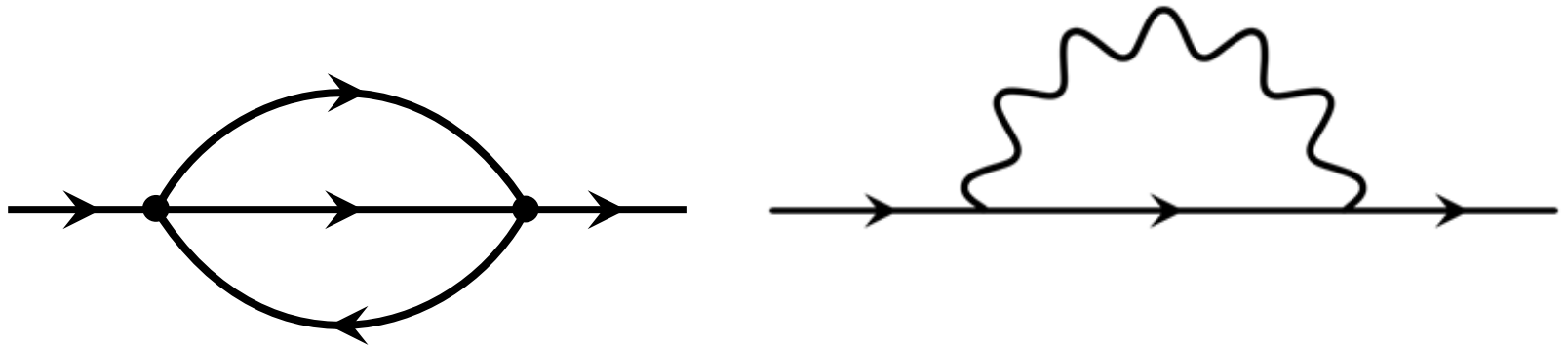
Scattering rate: Two channels

There are two channels:
1st electron-electron scattering and
2nd electron-boson scattering

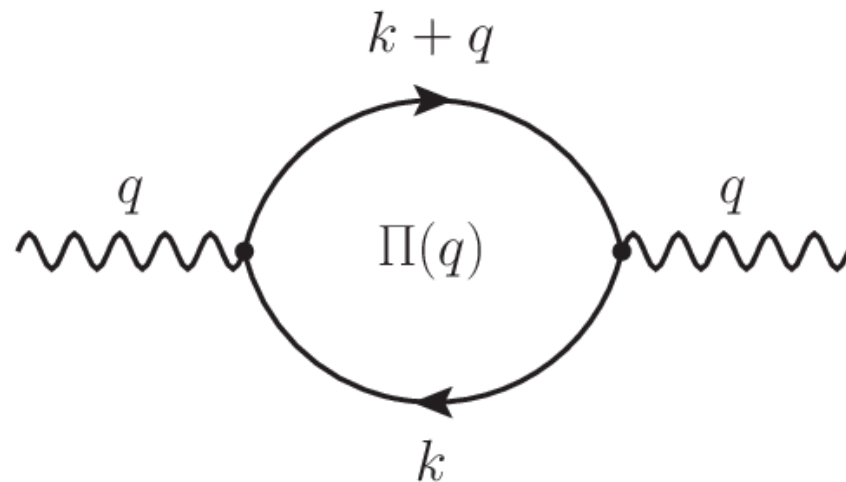




1 particle & 2 particle spectral functions

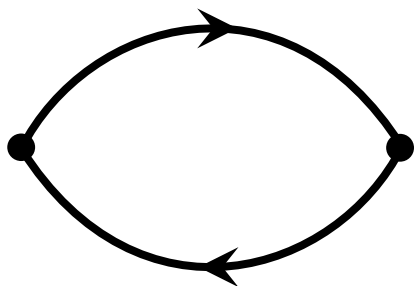


$$\Sigma \sim (G \star \chi)_{\mathbf{k}, \omega}$$



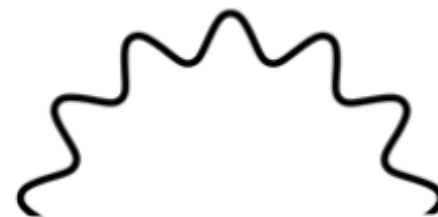
polarization bubble

$$\Sigma \sim G \star \chi$$



$$\chi \sim G \star G$$

$$\Sigma \sim G \star F$$

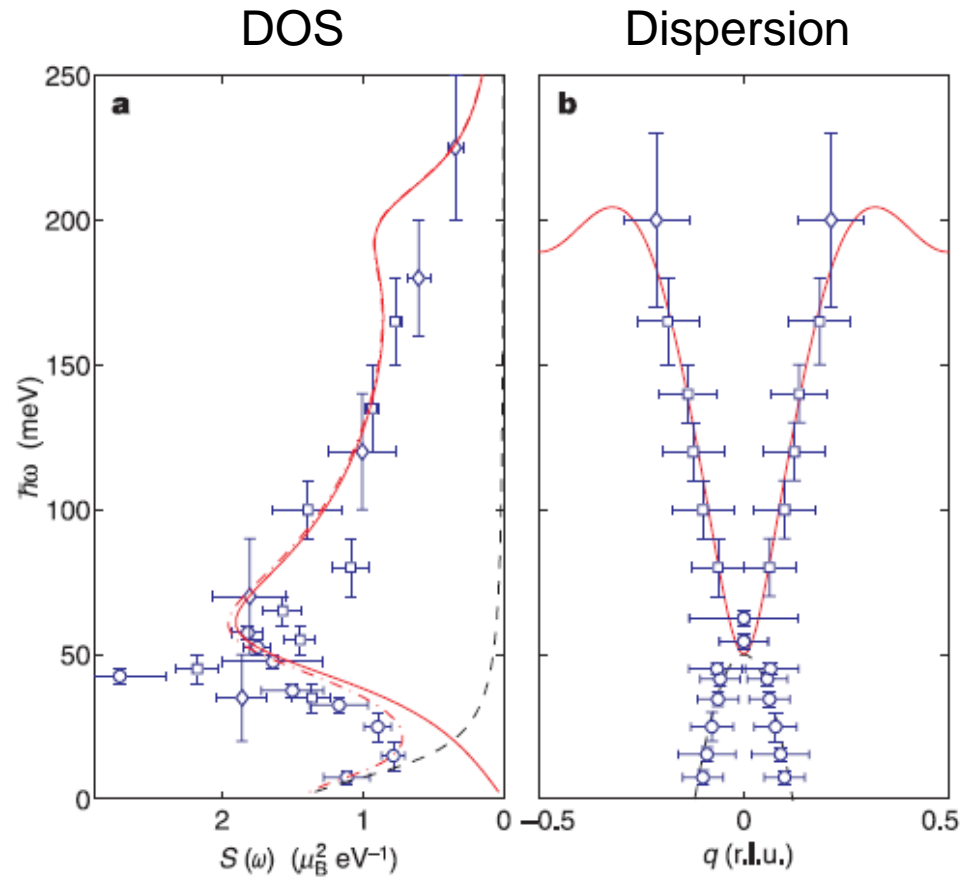
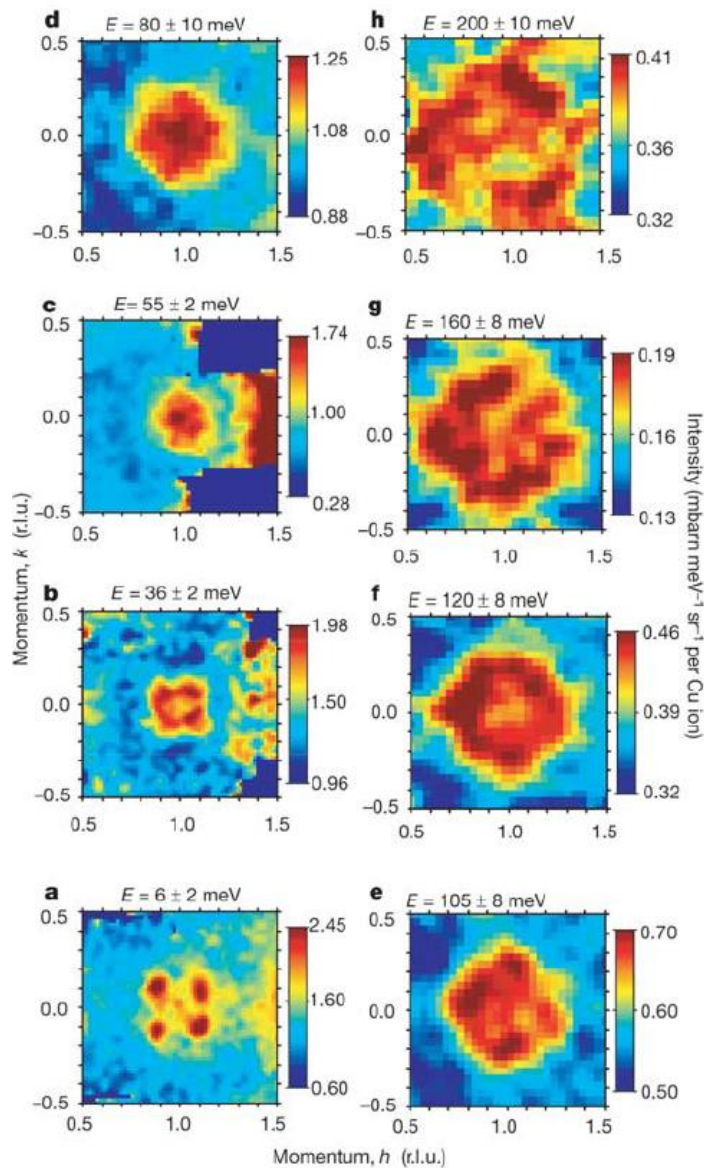


F

**Двохчастинкова функція –
функція Лінхардта**

$$\chi(\Omega, q)$$

Spin susceptibility structure

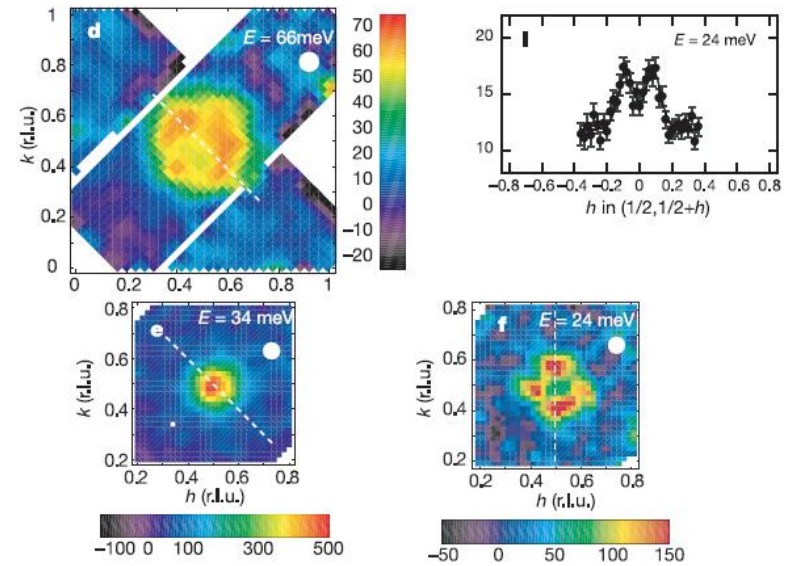
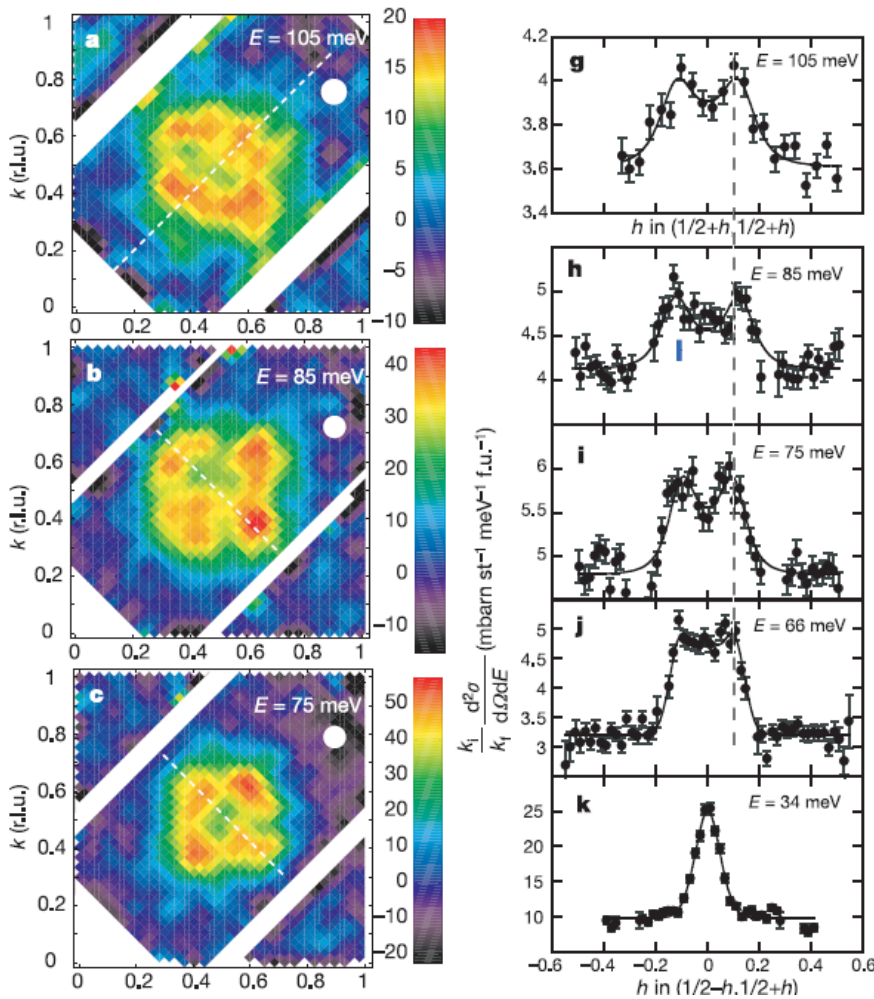


$\text{La}_{2-x}\text{Ba}_x\text{CuO}_4$ ('Zurich' oxide)

Tranquada *Nature* 2004

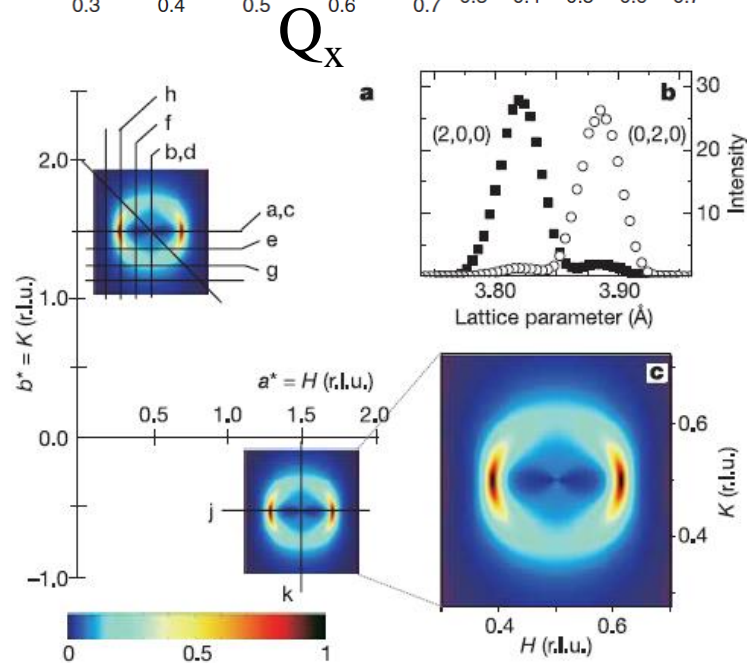
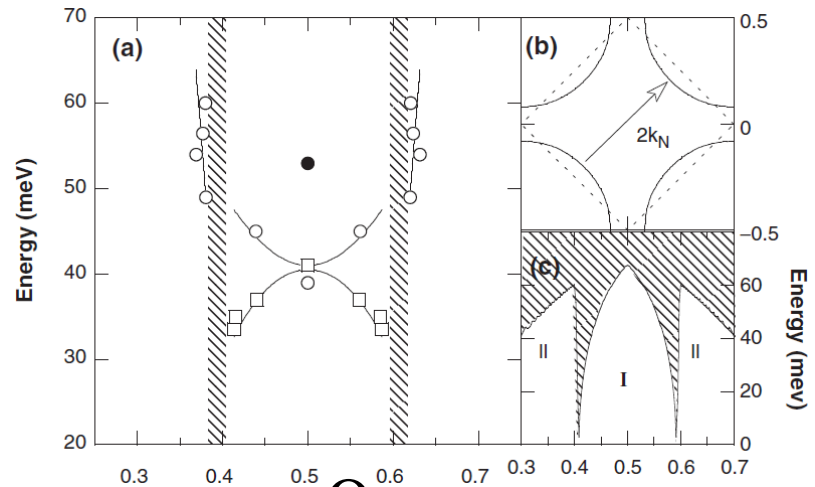
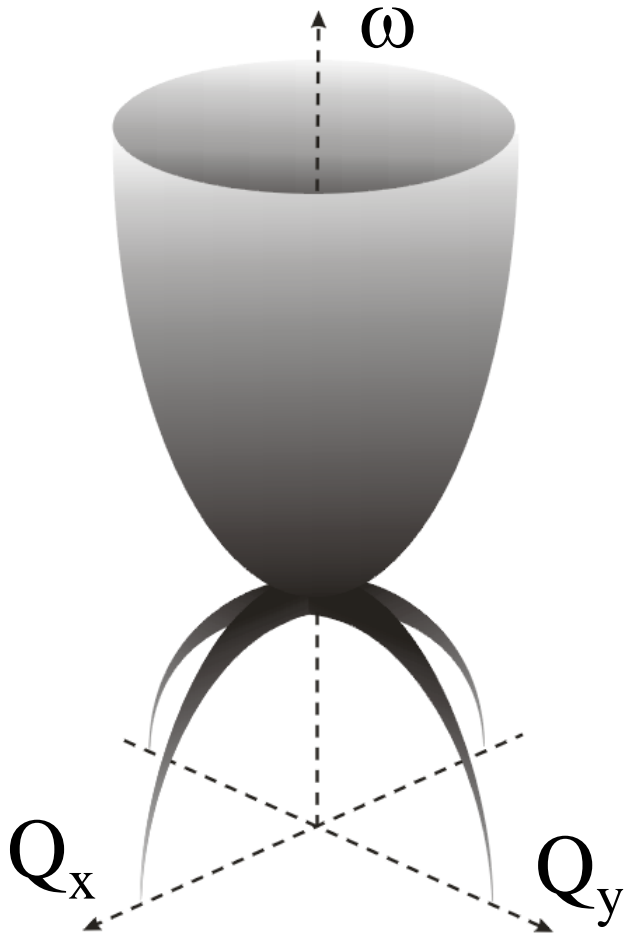
Spin susceptibility structure

$$\chi(\omega, \mathbf{Q})$$



$\text{YBa}_2\text{Cu}_3\text{O}_{6.6}$

Spin susceptibility structure



$\text{YBa}_2\text{Cu}_3\text{O}_{6.6}$

Looking for "fingerprints"

if 2nd order perturbation theory works

① $(\Delta, \Sigma) = \text{EE}(\Delta, \Sigma, \varepsilon, \chi)$ **SC**

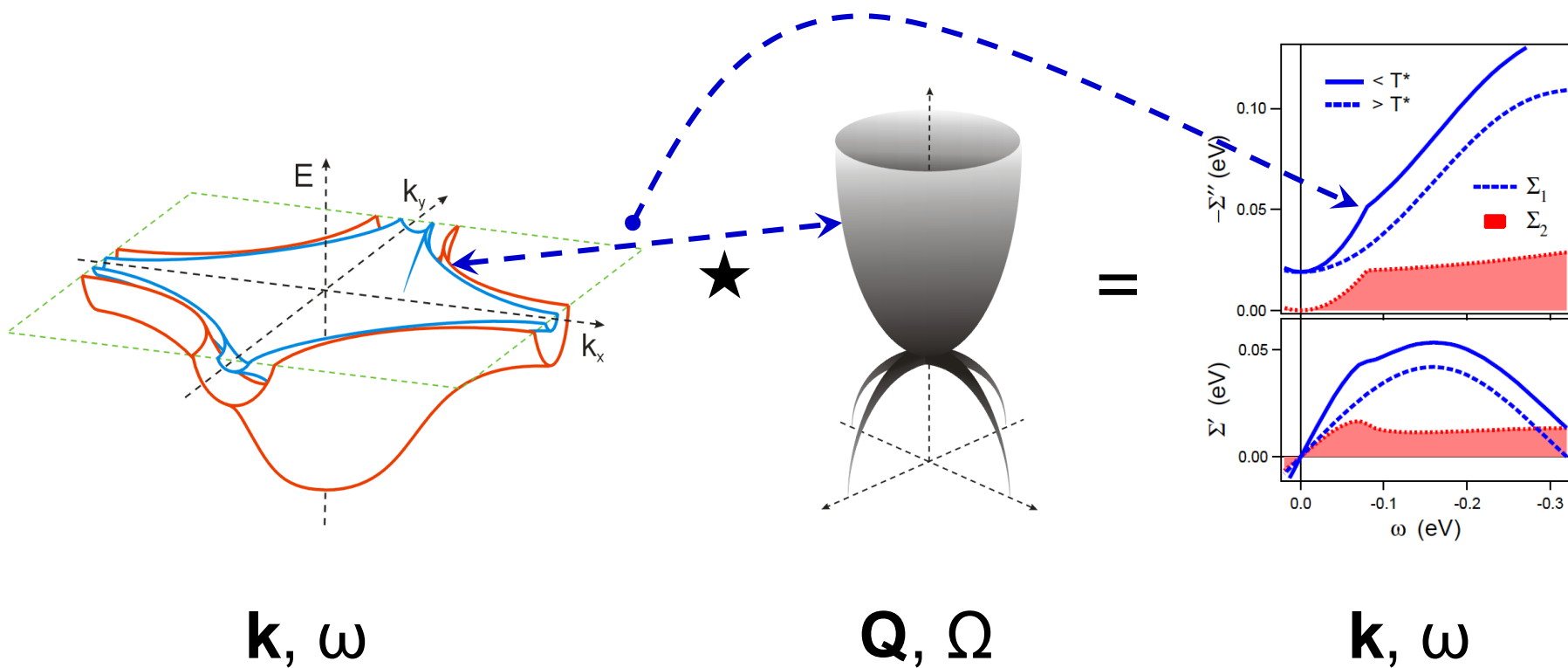
$$\Sigma \sim (G \star \chi)_{\mathbf{k}, \omega} \quad \mathbf{N}$$

② $X_{\text{it}} \sim (G \star G)_{\mathbf{k}, \omega}$ "itinerant"
magnetism

①

$$G \star \chi \sim \Sigma$$

$$\Sigma(\mathbf{k}, \omega) \sim \int G(\mathbf{k} + \mathbf{Q}, \omega + \Omega) \chi(\mathbf{Q}, \Omega) d\mathbf{Q} d\Omega$$



①

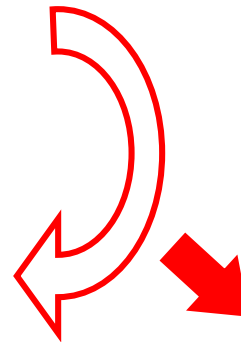
LDA or
ARPES

$$G_0 \star X_{\text{exp}} \sim \Sigma_i$$

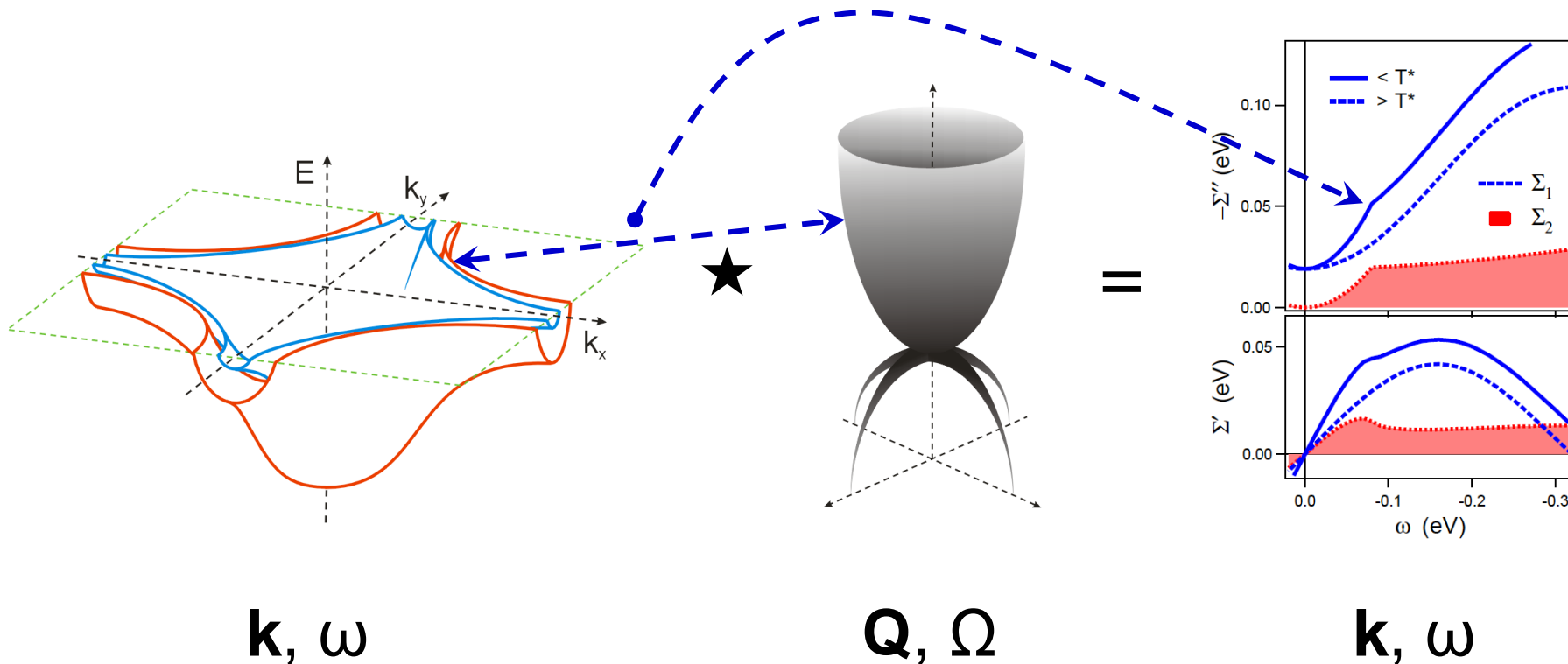


INS

$$G_i^{-1} = G_0^{-1} + \Sigma_i$$



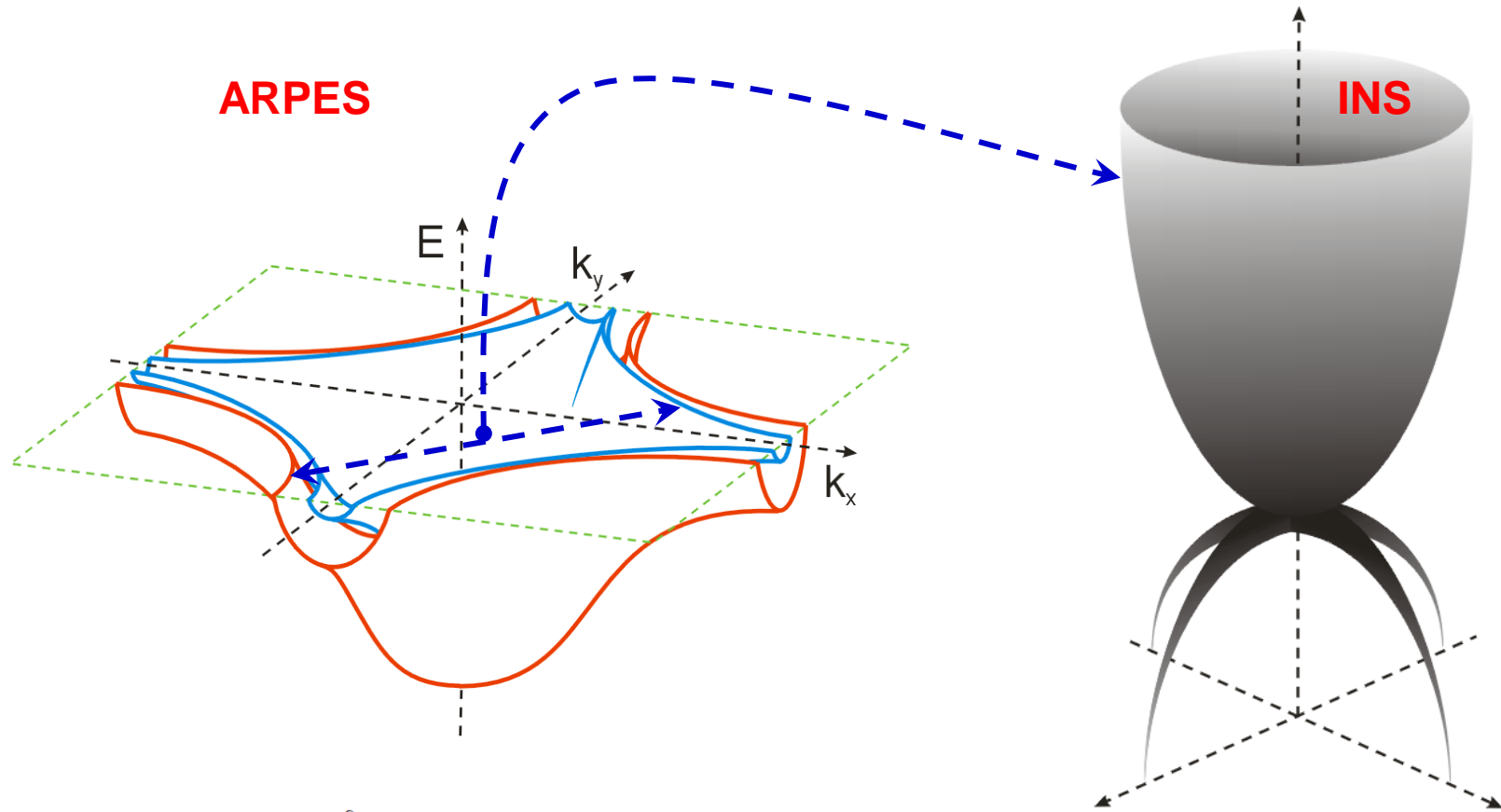
ARPES



②

itinerant

$$\chi \sim G \star G$$



$$\chi_0(\mathbf{Q}, \Omega) \propto -2i \int G(\mathbf{k}, \omega) G(\mathbf{k} + \mathbf{Q}, \omega + \Omega) d^2k d\omega$$

$$\chi(\mathbf{Q}, \Omega) = \chi_0(\mathbf{Q}, \Omega) / [1 + J_{\mathbf{Q}} \chi_0(\mathbf{Q}, \Omega)] \quad \text{RPA}$$

② $\chi \sim G \star G$

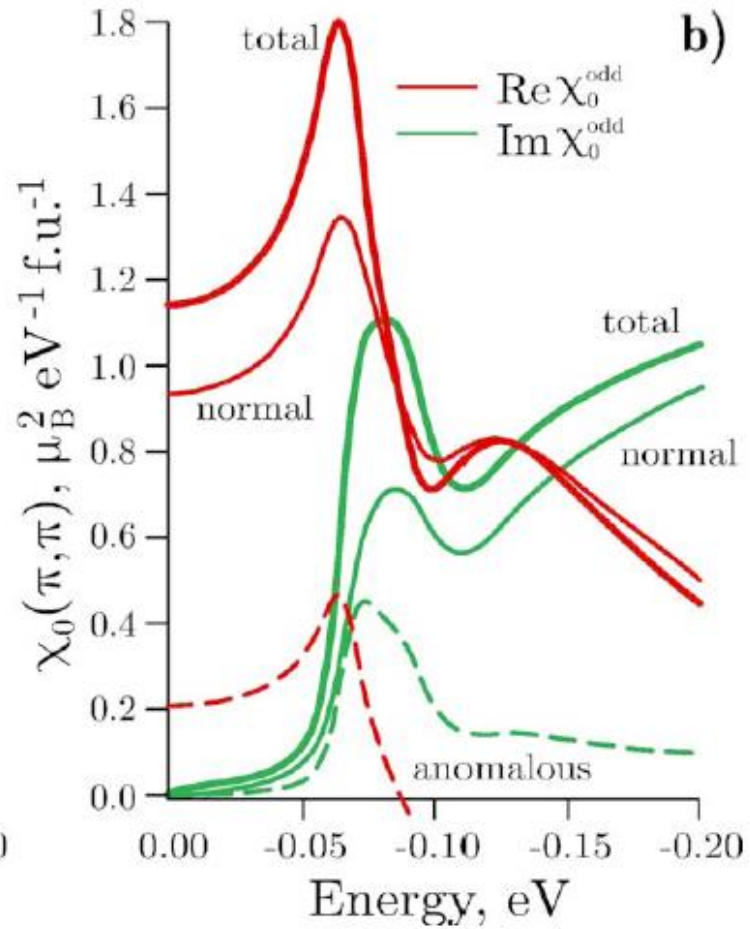
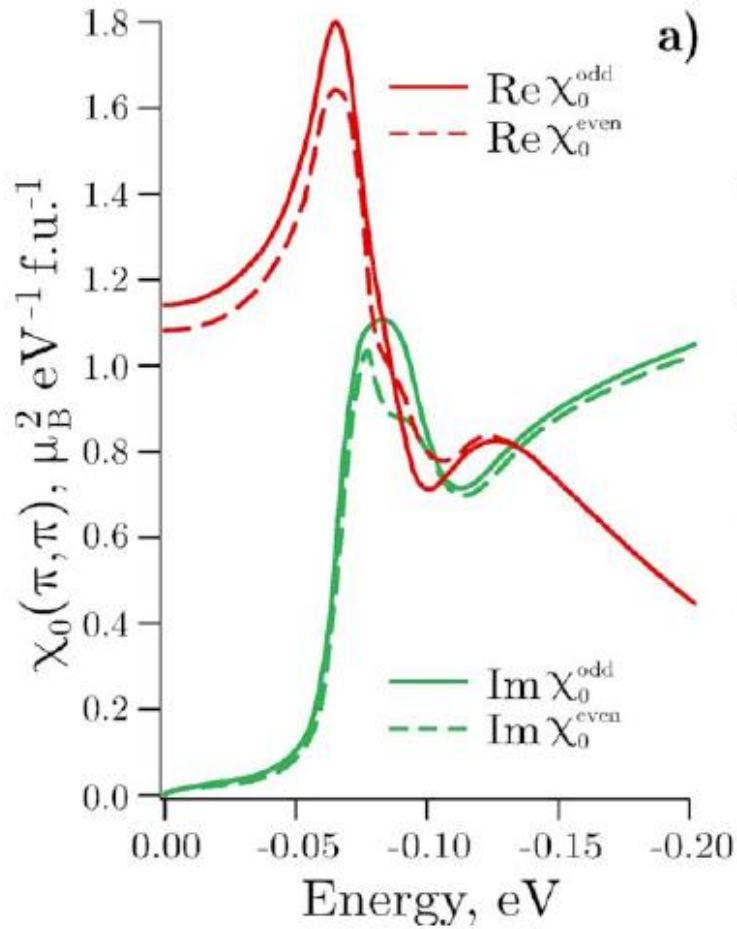
$$\chi_0(\mathbf{Q}, i\Omega_n) = \frac{1}{\pi^2} \int \sum_m G(\mathbf{k}, i\omega_m) G(\mathbf{k} + \mathbf{Q}, i\omega_m + i\Omega_n) d\mathbf{k}$$

$$\chi_0^{o,e}(\mathbf{Q}, \Omega) = \sum_{\substack{i=j(o) \\ i \neq j(e)}} \iint_{-\infty}^{+\infty} C_{ij}(\mathbf{k}, \epsilon, \nu) \frac{n_f(\nu) - n_f(\epsilon)}{\Omega + \nu - \epsilon + i\Gamma} d\nu d\epsilon$$

where

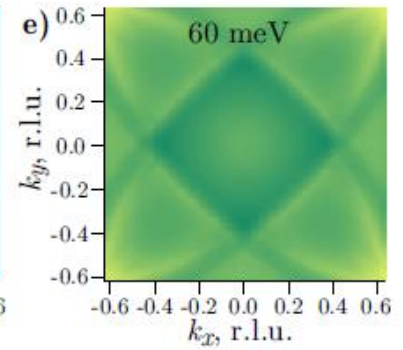
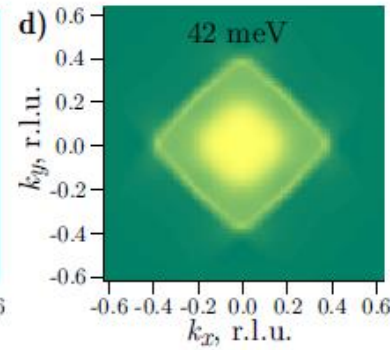
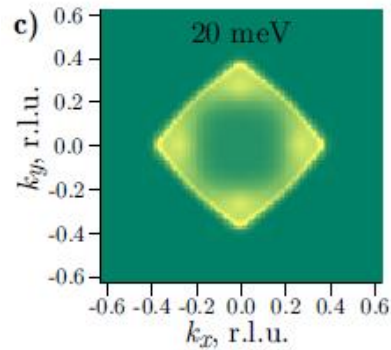
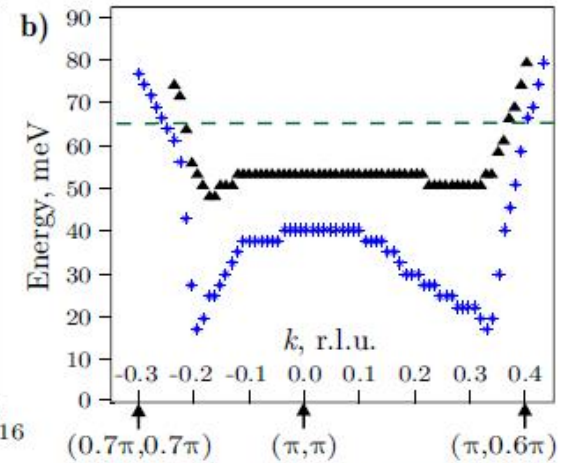
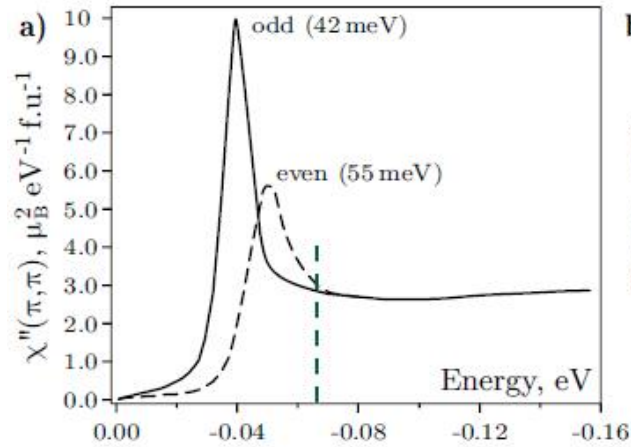
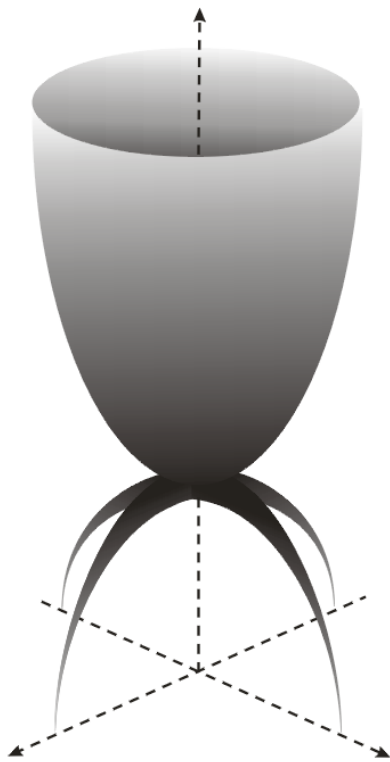
$$C_{ij}(\mathbf{k}, \epsilon, \nu) = \frac{1}{\pi^2} \int [\text{Im } G_i(\mathbf{k}, \epsilon) \text{Im } G_j(\mathbf{k} + \mathbf{Q}, \nu) + \text{Im } F_i(\mathbf{k}, \epsilon) \text{Im } F_j(\mathbf{k} + \mathbf{Q}, \nu)] d\mathbf{k}.$$

②



②

$$\chi^{o,e}(\mathbf{Q}, \Omega) = \chi_0^{o,e}(\mathbf{Q}, \Omega) / [1 - J_{\mathbf{Q}}^{o,e} \chi_0^{o,e}(\mathbf{Q}, \Omega)]$$



Fourier transform (FT) & Correlation

Fourier transform (FT)

$$F(\omega) = \int_{-\infty}^{\infty} f(t)e^{-i\omega t} dt$$

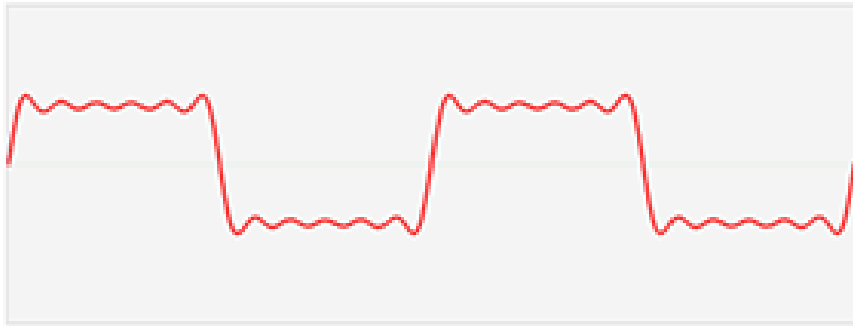
$$f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega)e^{i\omega t} d\omega$$

Fourier transform (FT)

$$\hat{f}(\xi) = \int_{-\infty}^{\infty} f(x) e^{-2\pi i x \xi} dx,$$

$$f(x) = \int_{-\infty}^{\infty} \hat{f}(\xi) e^{2\pi i x \xi} d\xi,$$

Fourier series

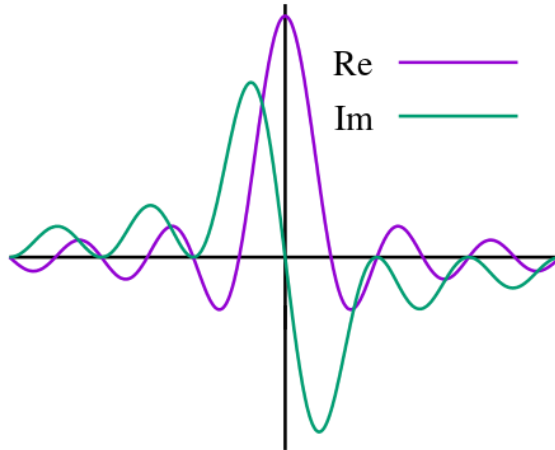
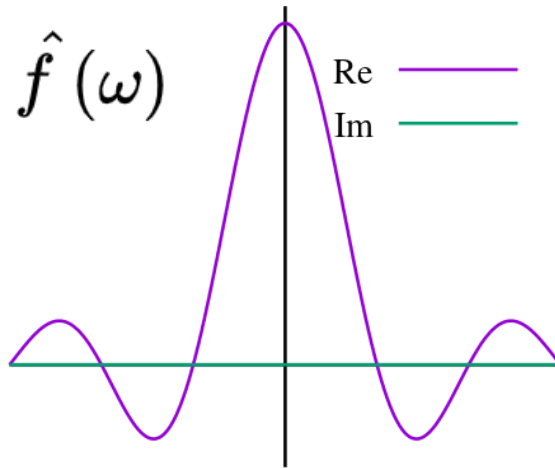
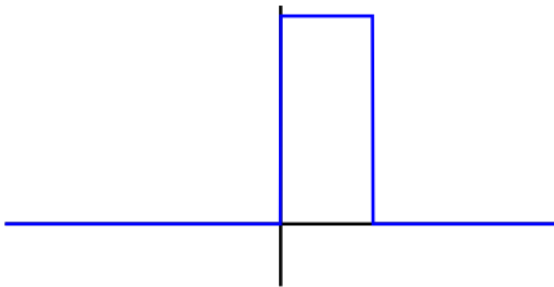
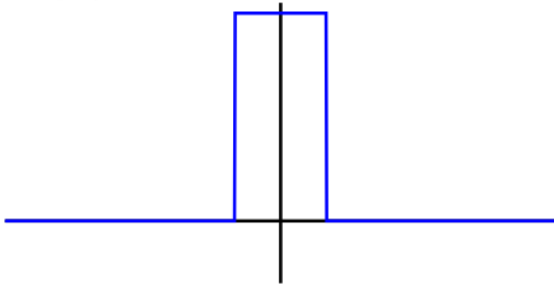


$$f(x) = \sum_{n=-\infty}^{\infty} c_n e^{2\pi i \left(\frac{n}{T}\right) x}$$

$$c_n = \frac{1}{T} \int_{-\frac{T}{2}}^{\frac{T}{2}} f(x) e^{-2\pi i \left(\frac{n}{T}\right) x} dx.$$

Fourier transform (FT)

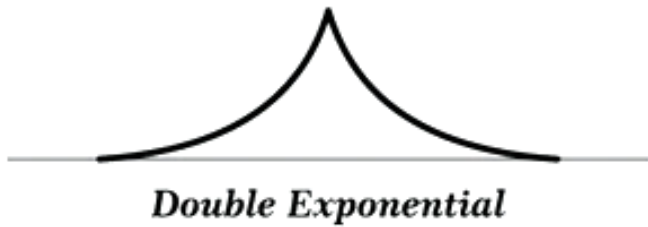
$f(t)$



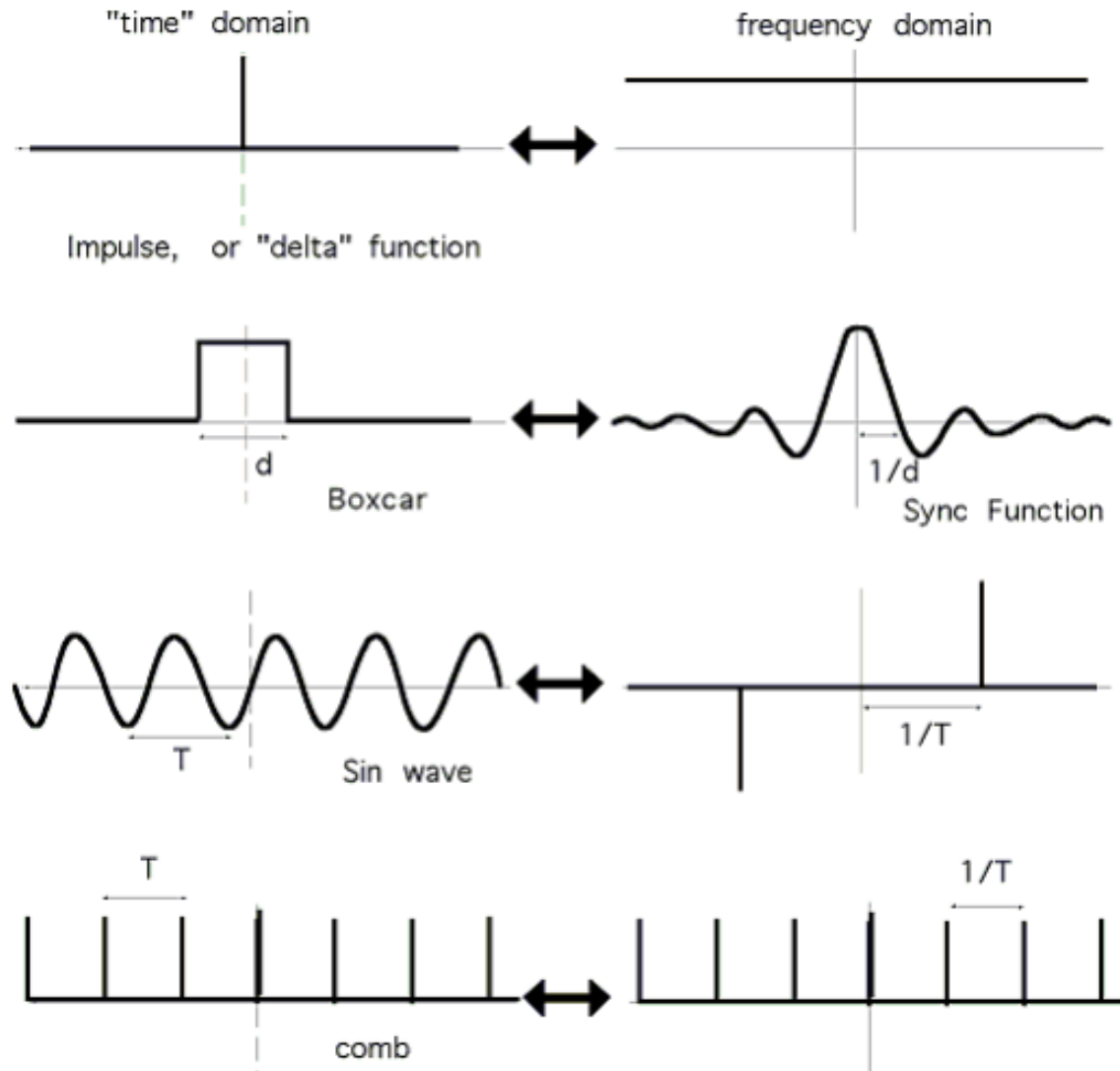
Fourier transform (FT)

$f(t)$

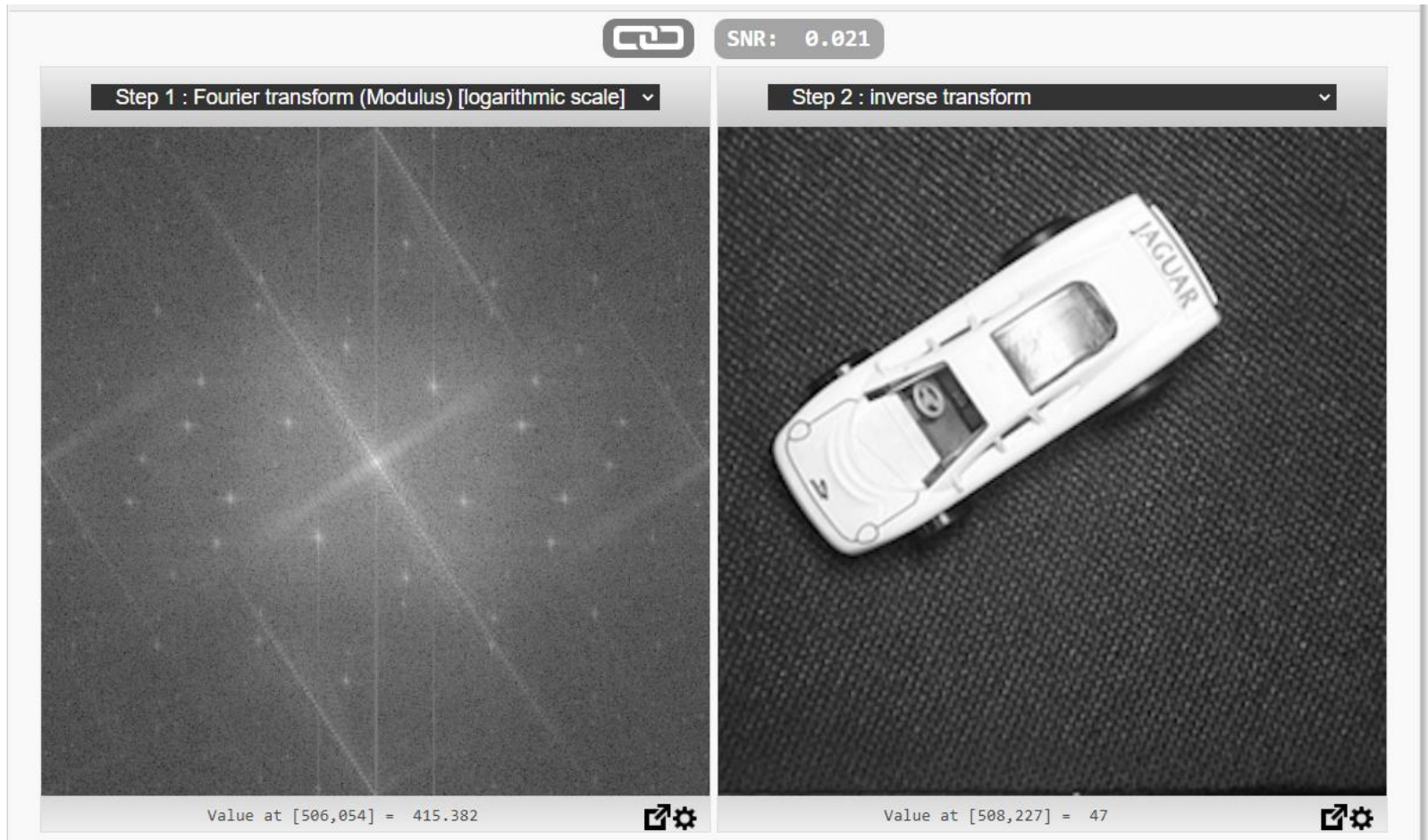
$\hat{f}(\omega)$



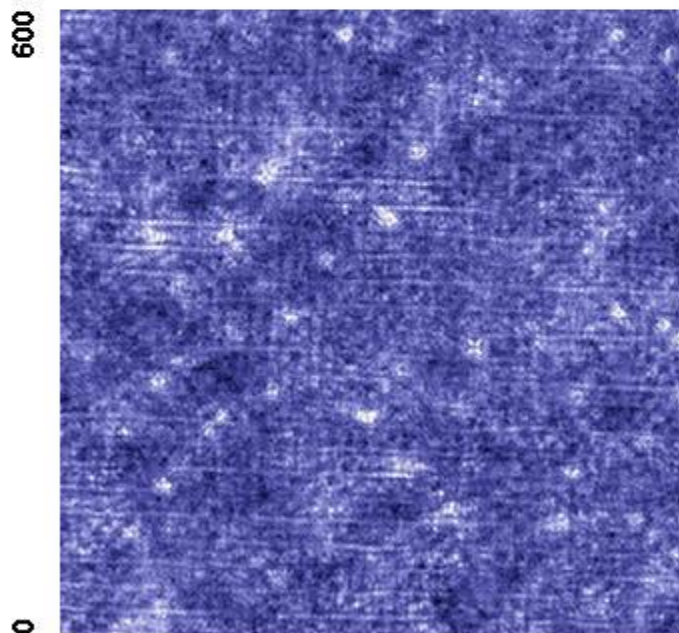
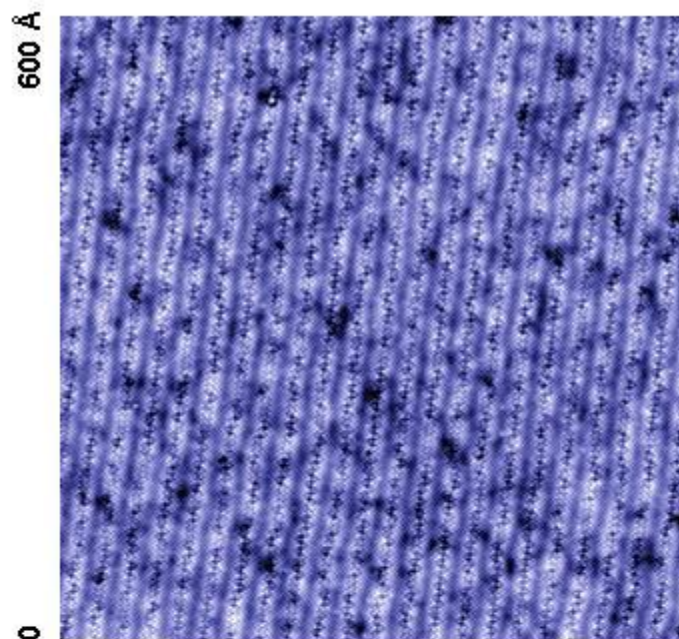
Fourier transform (FT)



Fourier transform (FT)

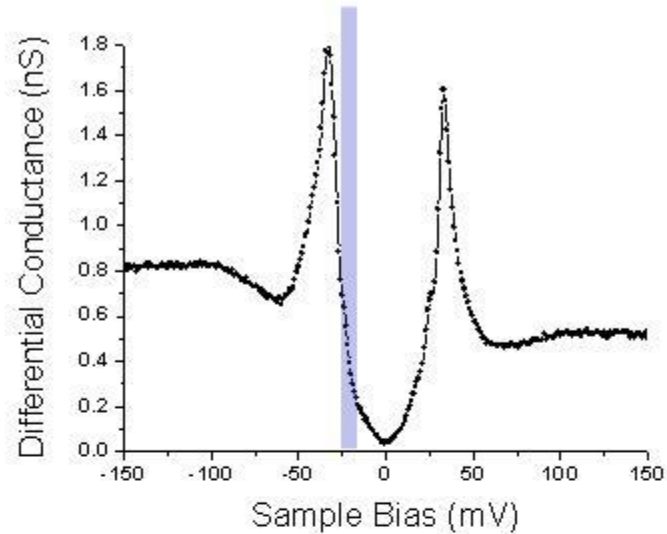


<http://bigwww.epfl.ch/demo/ip/demos/FFT/>

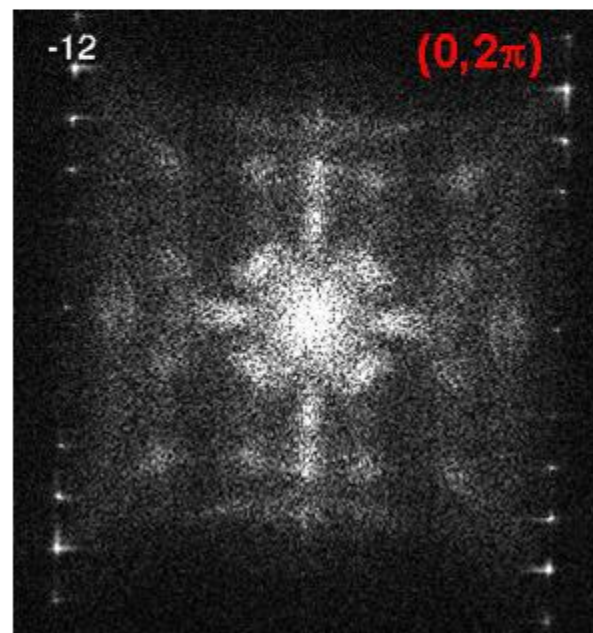


$g(\vec{r}, E = -12 \text{ meV})$

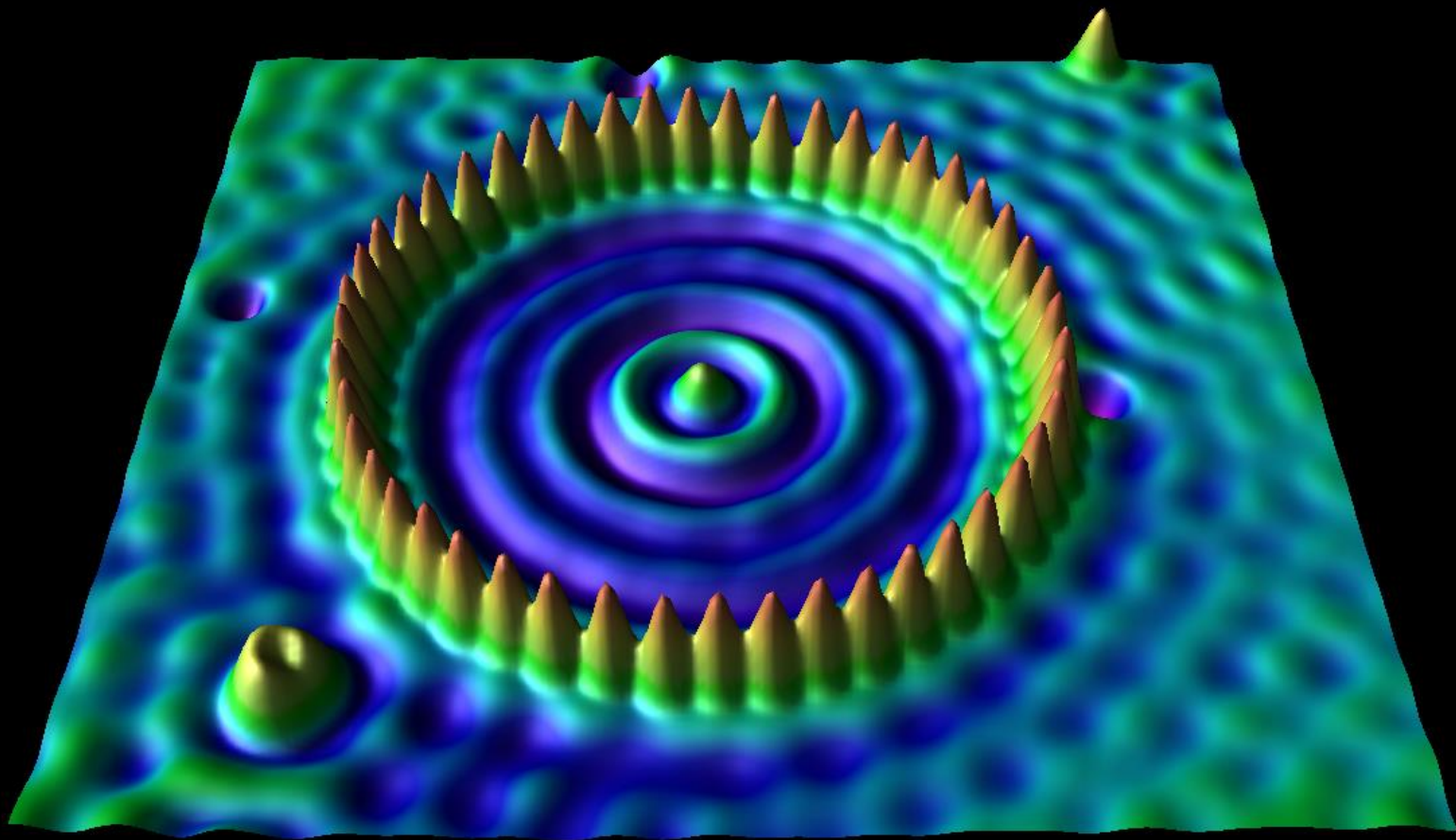
FFT shows
q-vector of
LDOS
modulations



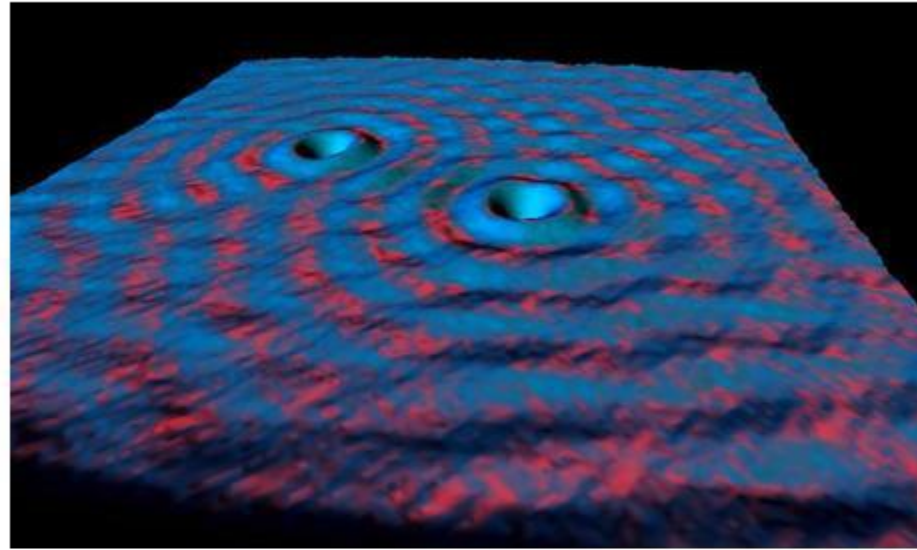
Point $dI/dV \equiv g(\vec{r}, E)$ Spectrum



$g(\vec{q}, E = -12 \text{ meV})$



Quasiparticle Interference at Impurity Atoms

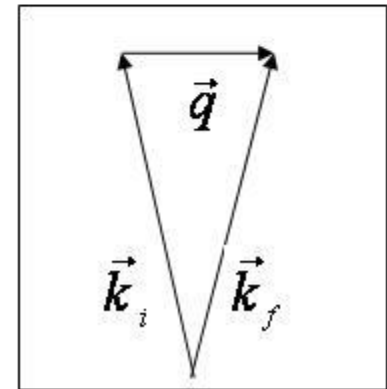


Crommie, Lutz & Eigler, *Nature* **363**, 524 (1993)

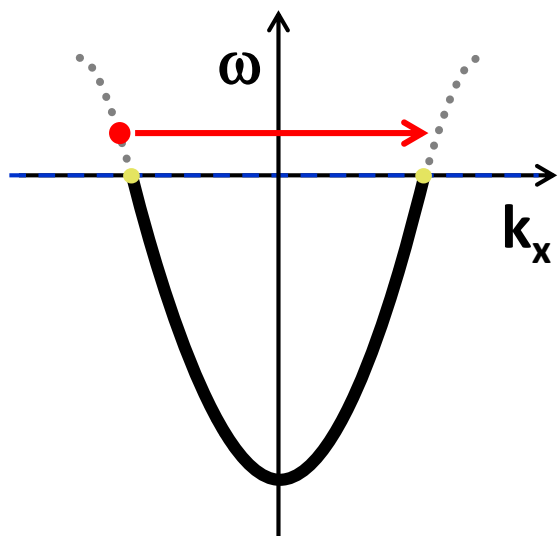
- Mixing of $|\vec{k}_1\rangle$ and $|\vec{k}_2\rangle$ by scattering creates interference term, with wavevector

$$\vec{q} = \vec{k}_2 - \vec{k}_1$$

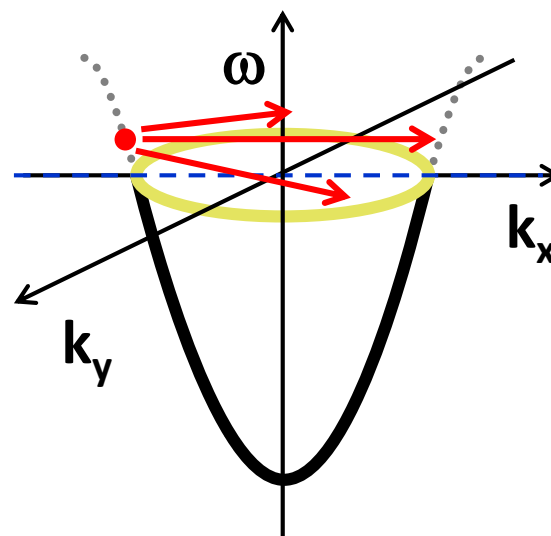
- Interference results in modulations in LDOS with wavelength $\lambda = \frac{2\pi}{|\vec{q}|}$



Impurity scattering

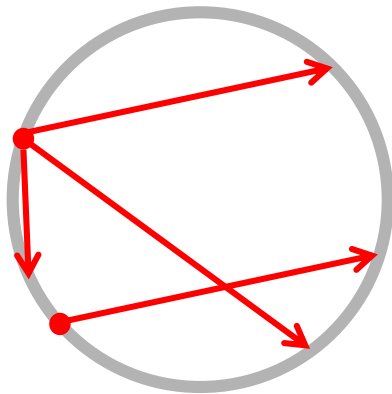


1D

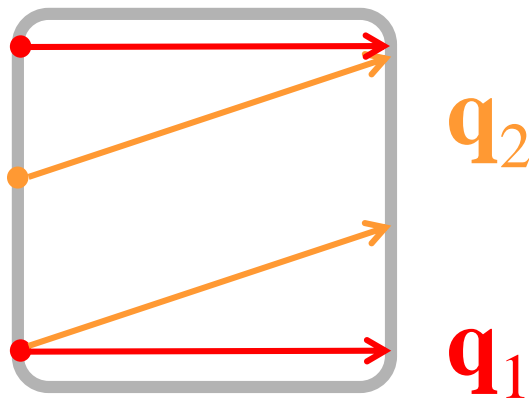


2D

$$\chi \sim G \star G \sim \int A(\mathbf{k})A(\mathbf{k} + \mathbf{q})d\mathbf{k}$$

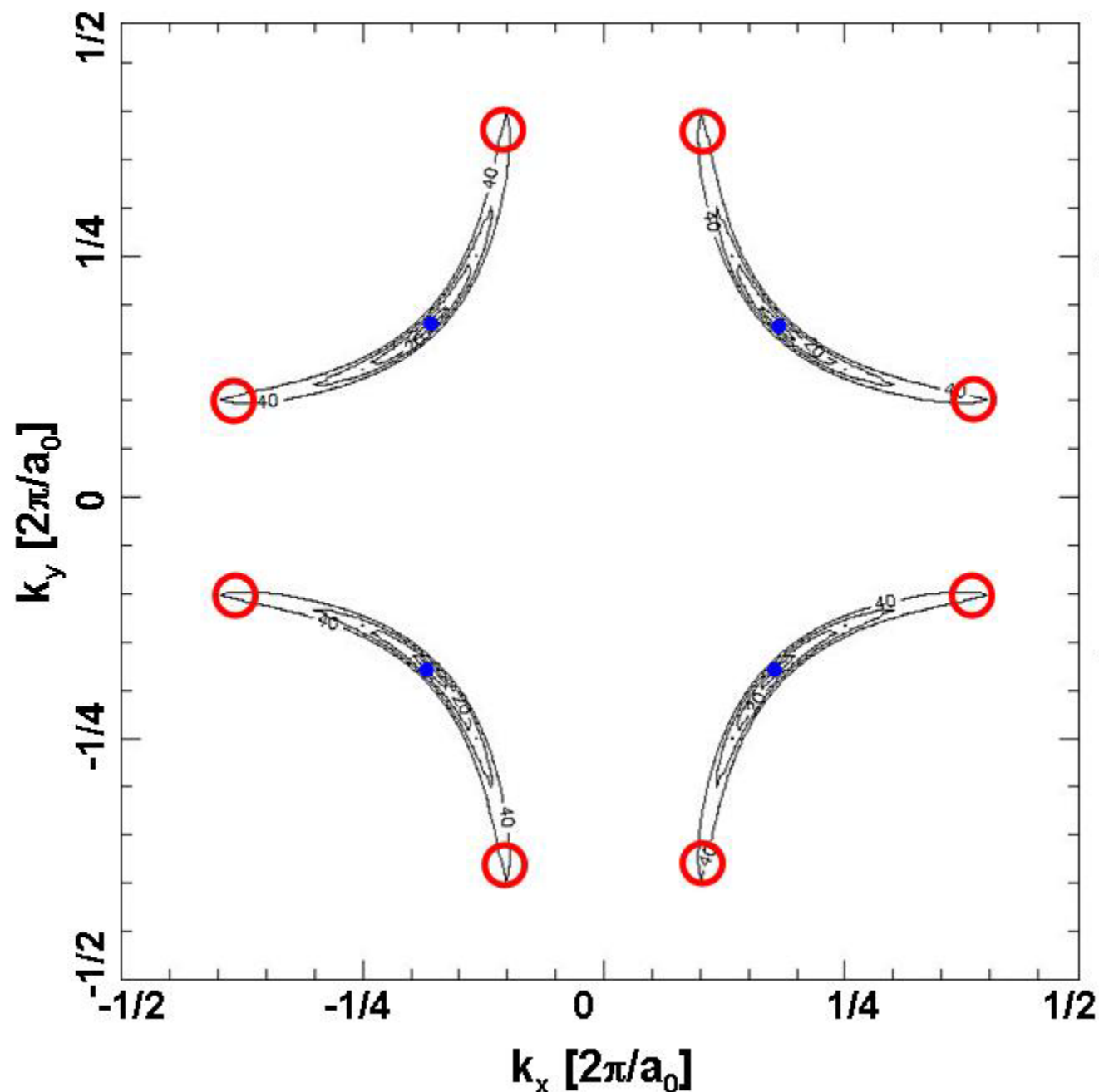


$$\mathbf{q} = \mathbf{k}_1 - \mathbf{k}_0$$



$$\chi(\mathbf{q}_1) > \chi(\mathbf{q}_2)$$

Octet of regions at ends of 'bananas' have smallest $dE/|dk|$



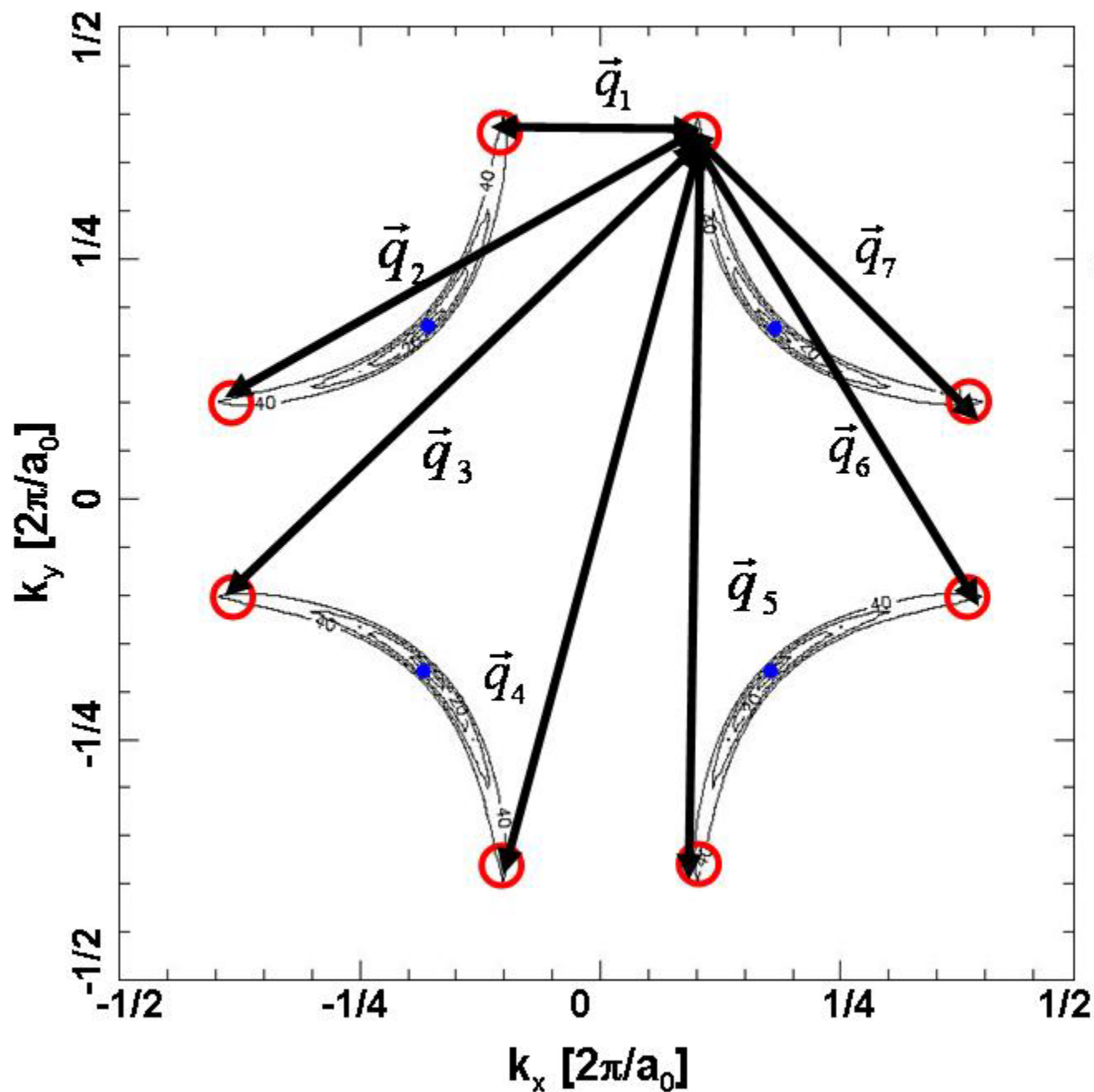
Density of States

$$n(E) = \oint_{E(\vec{k})=E} \frac{1}{|\nabla_{\vec{k}} E(\vec{k})|} d\vec{k}$$



This octet of locations at the tips of the 'bananas' provide maximum contribution to $n_f(E)$ and thus dominate elastic scattering processes.

Octet of regions at ends of 'bananas' have smallest $dE/|dk|$



Density of States

$$n(E) = \oint_{E(\vec{k})=E} \frac{1}{|\nabla_{\vec{k}} E(\vec{k})|} d\vec{k}$$



Characteristic set of quasiparticle interference wavevectors which is different at each energy.

Way to ARPES

Impurity scattering hypothesis: $|FS(\mathbf{r})| = ACA(\mathbf{k})$

Definition of autocorrelation: $ACA(\mathbf{k}) = \int A(\mathbf{k})A(\mathbf{k} + \mathbf{q})d\mathbf{k}$

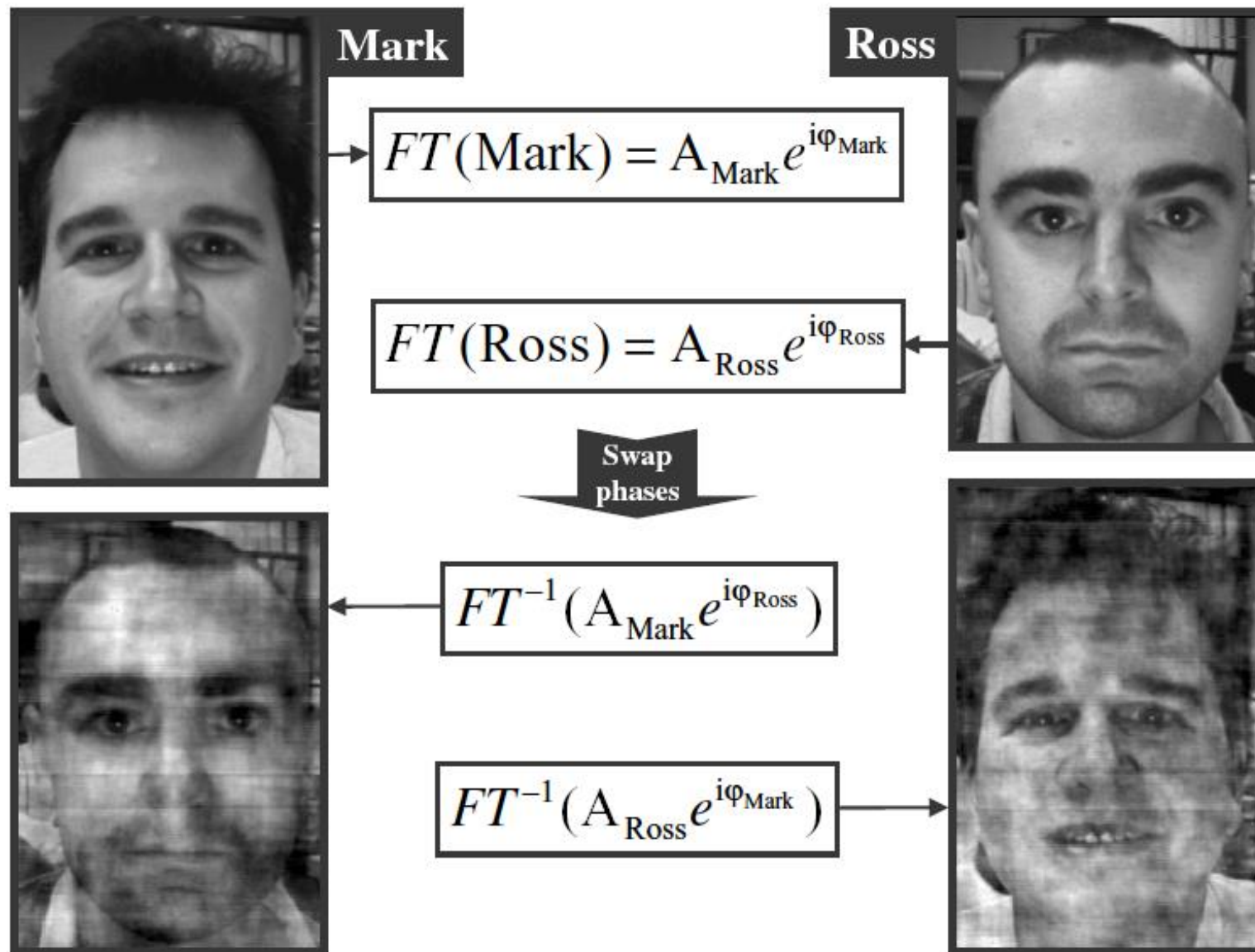
Another definition through
the Fourier transform
(Wiener-Khinchin theorem):

$$ACA(\mathbf{k}) = F|FA(\mathbf{k})|^2$$

$$|FA(\mathbf{k})| = \sqrt{F|FS(\mathbf{r})|} = R(\rho)$$

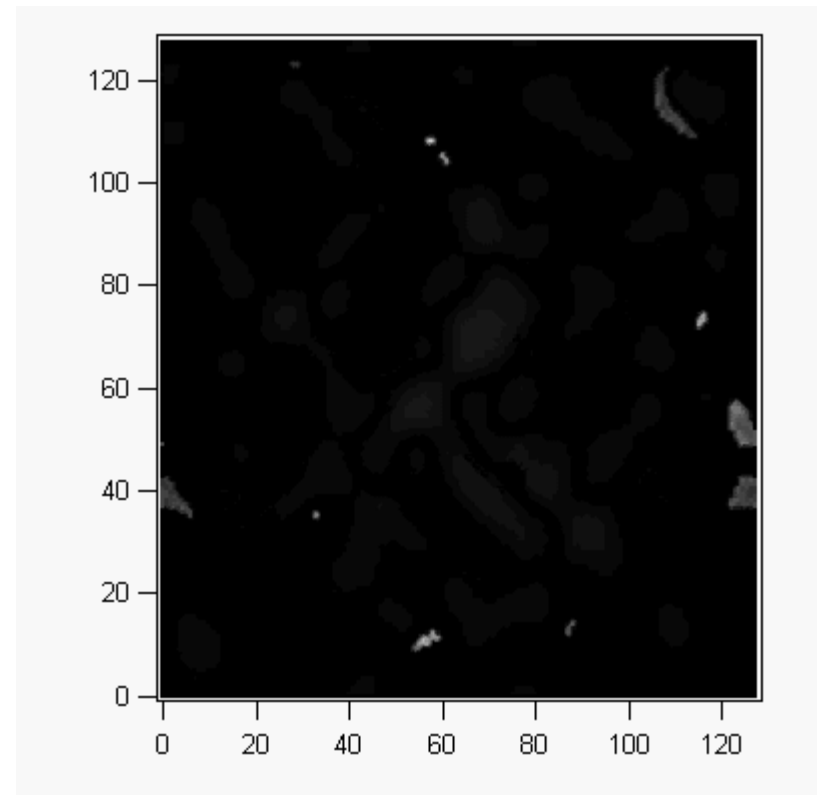
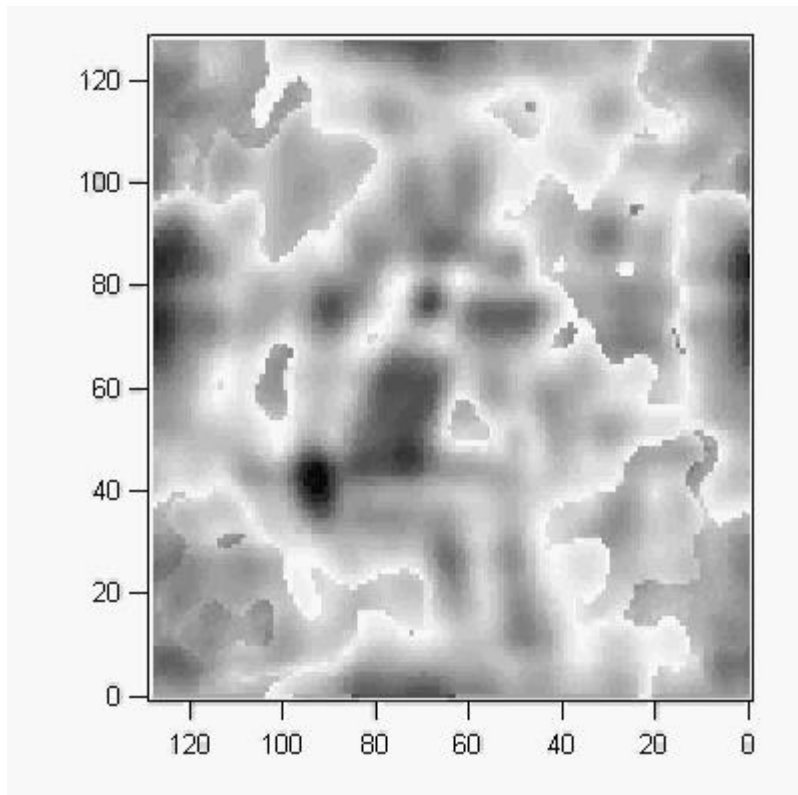
Phase retrieval algorithm: $A(\mathbf{k}) = PRA\sqrt{F|FS(\mathbf{r})|}$

Phase retrieval algorithm

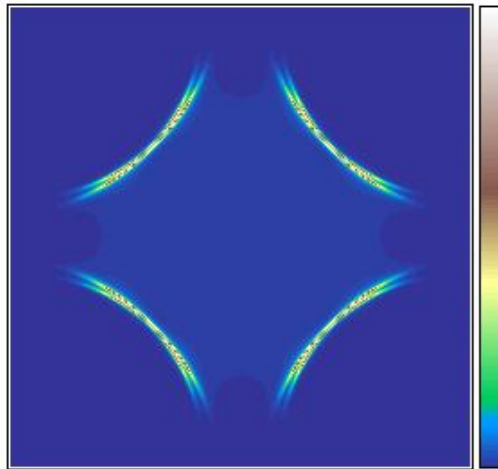


Phase retrieval algorithm

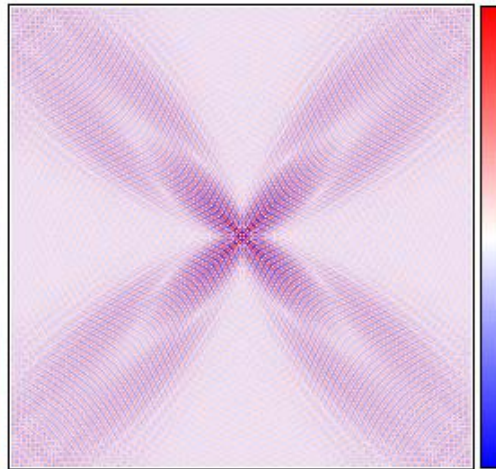
$$A(\mathbf{k}) = \text{PRA} |F A(\mathbf{k})|$$



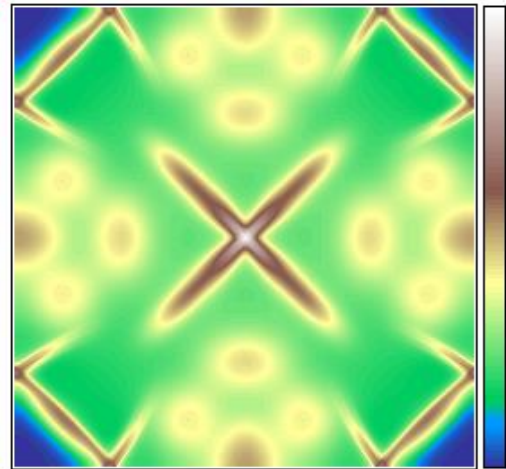
Representations of quasiparticles in different spaces of high- T_c cuprates:



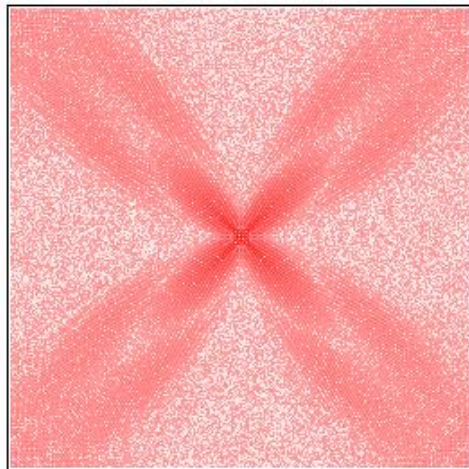
A vs k



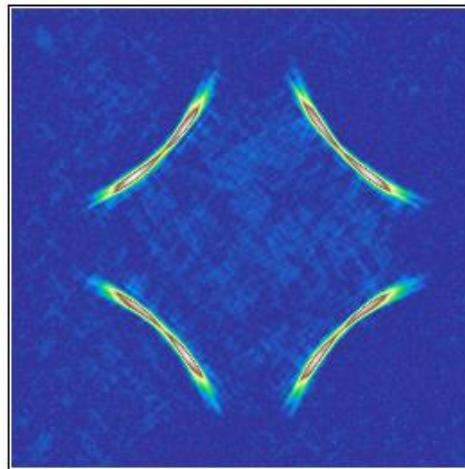
FA vs ρ



ACA vs q



$|FA| + 50\%$ noise vs ρ

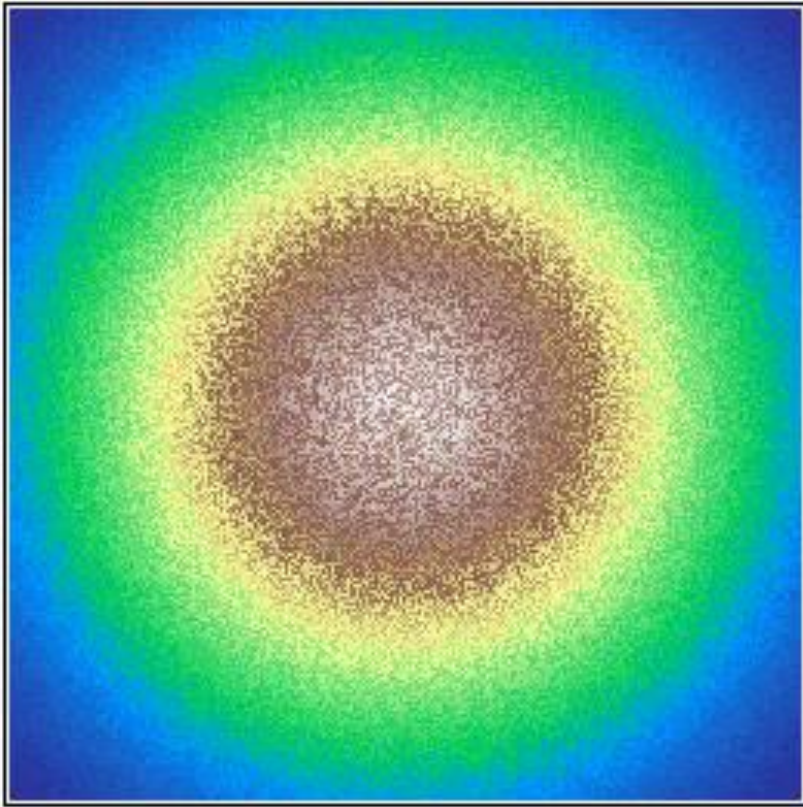


A' vs k



ACA' vs q

Phase retrieval algorithm



$$\mathcal{R}_n = \mathbf{F}A_n,$$

$$\mathcal{R}'_n = R \exp[i\arg(\mathcal{R}_n)],$$

$$\mathcal{A}'_n = \mathbf{F}^{-1}\mathcal{R}'_n,$$

$$A_{n+1} = \begin{cases} \operatorname{Re}(\mathcal{A}'_n) & \text{if } \operatorname{Re}(\mathcal{A}'_n) \geq 0, \\ \operatorname{Re}(A_n - \beta\mathcal{A}'_n) & \text{if } \operatorname{Re}(\mathcal{A}'_n) < 0, \end{cases}$$

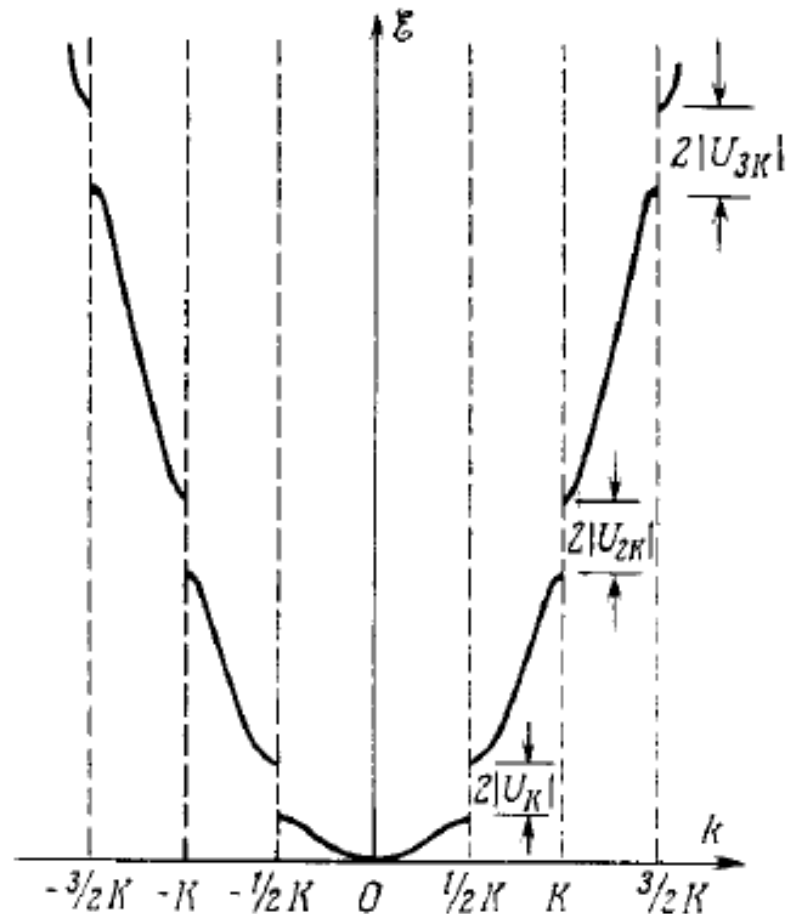
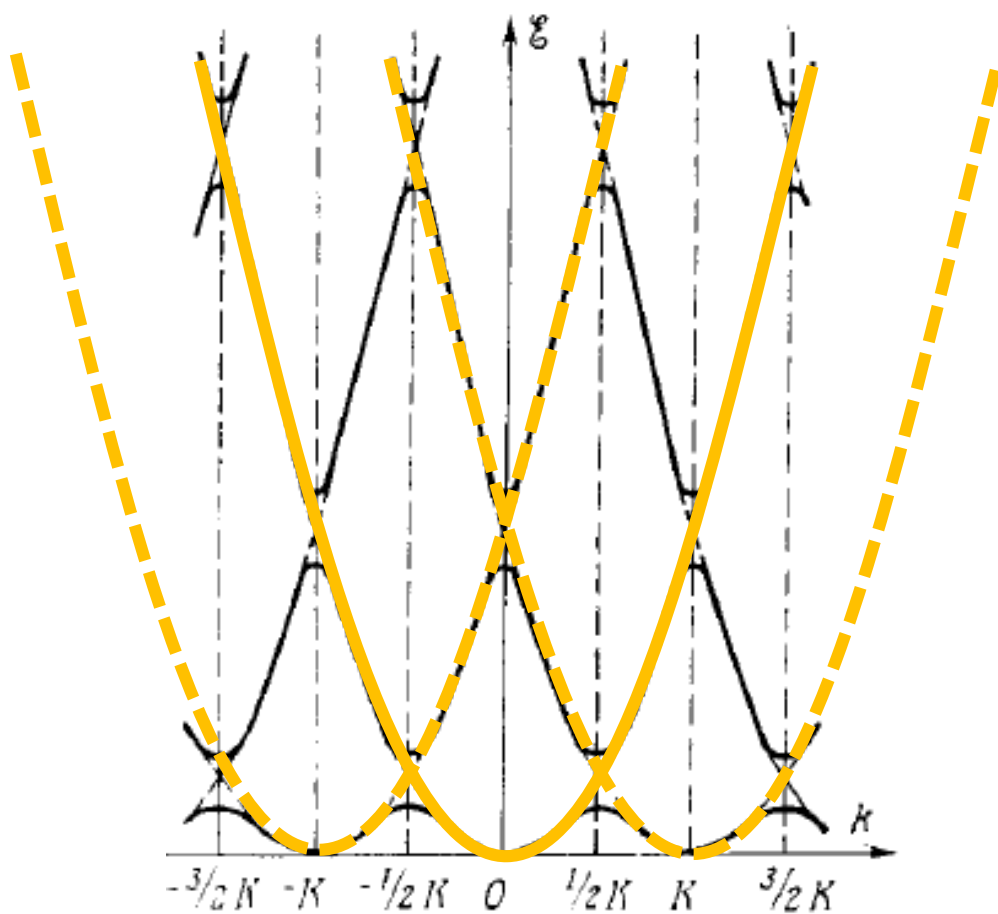
Теорема Блоха

$$\psi(\mathbf{r}) = e^{i\mathbf{k}\cdot\mathbf{r}} u(\mathbf{r})$$

$$\Psi_{n, \mathbf{k}+\mathbf{K}}(\mathbf{r}) = \Psi_{n\mathbf{k}}(\mathbf{r})$$

$$\mathcal{E}_{n, \mathbf{k}+\mathbf{K}} = \mathcal{E}_{n\mathbf{k}}$$

$$\mathcal{E} = \frac{1}{2} (\mathcal{E}_{\mathbf{q}}^0 + \mathcal{E}_{\mathbf{q}-\mathbf{K}}^0) \pm \left[\left(\frac{\mathcal{E}_{\mathbf{q}}^0 - \mathcal{E}_{\mathbf{q}-\mathbf{K}}^0}{2} \right)^2 + |U_{\mathbf{K}}|^2 \right]^{1/2}$$



$2H-Cu_xTaSe_2$

