

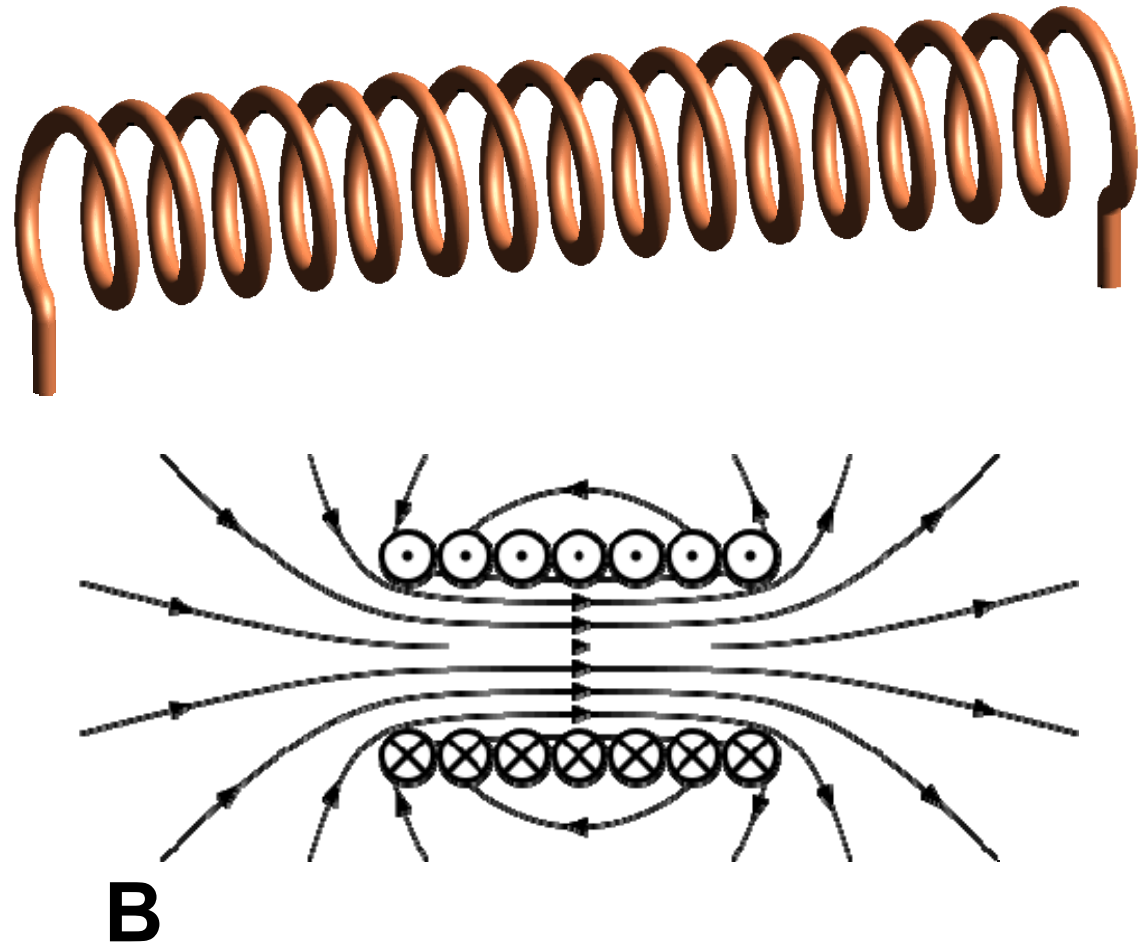
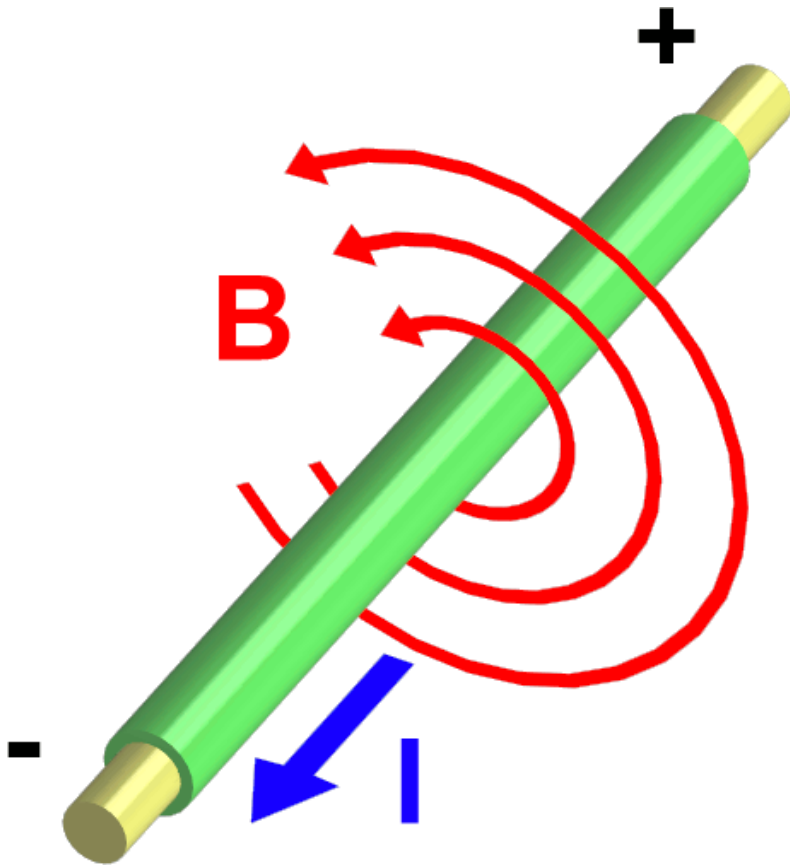
Physical phenomena in high magnetic fields

Dmytro Kamenskyi (d.kamenskyi@science.ru.nl)

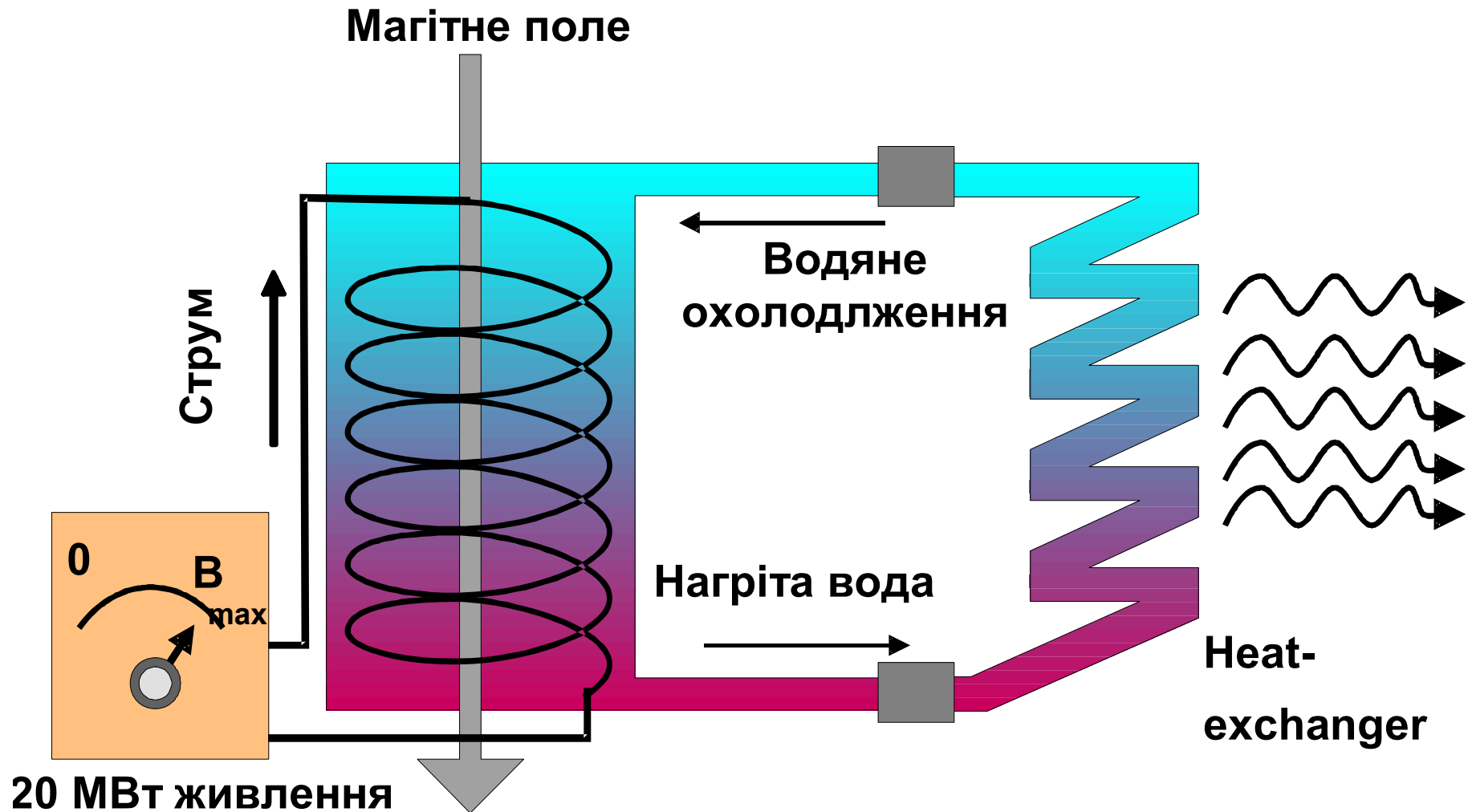
FELIX and HFML



Що таке магнітне поле?



Як отримати сильне магнітне поле?

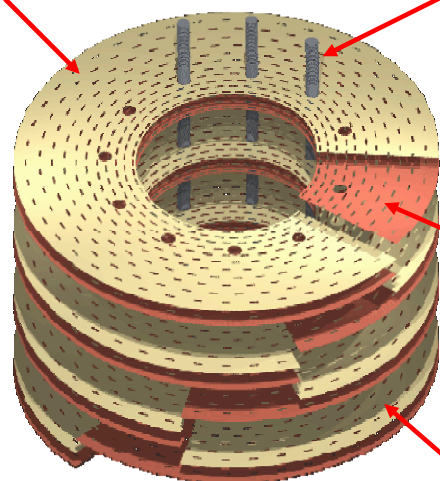


33 Тесла резистивний магніт

дизайн Біттера

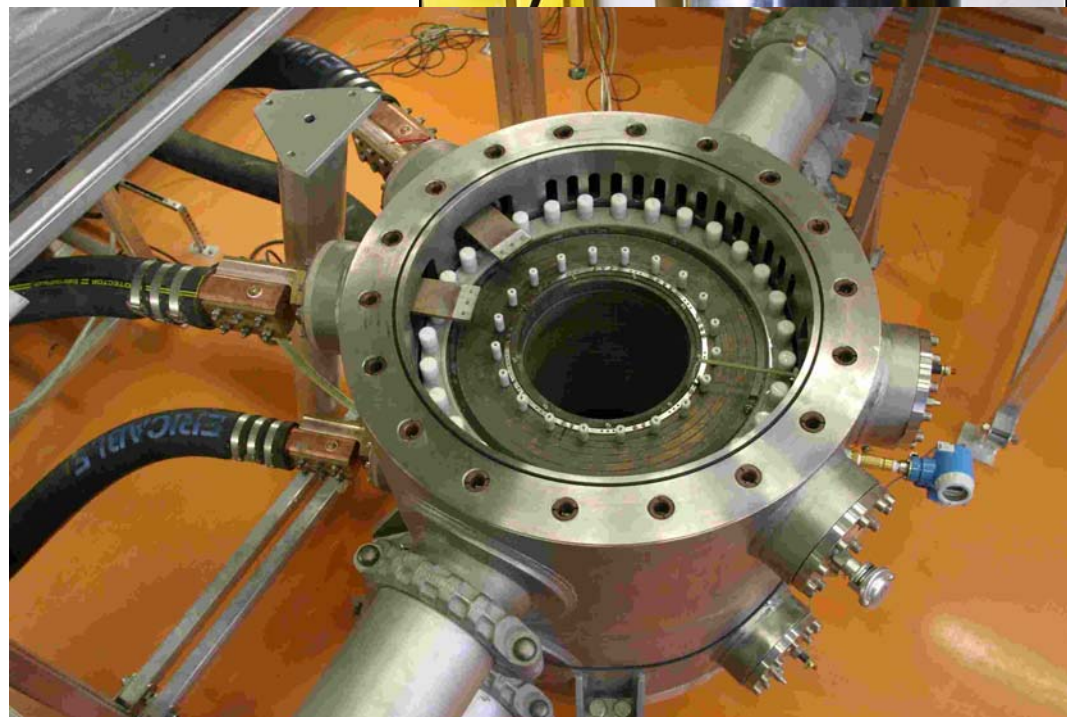
Отвори для охолодження

Стрижні для фіксації дисків



Си диск

Ізоляція



33 Т магніт:

Складається з 3-х катушок

- 36 кА, 18 МВт
- 145 літрів/сек вода

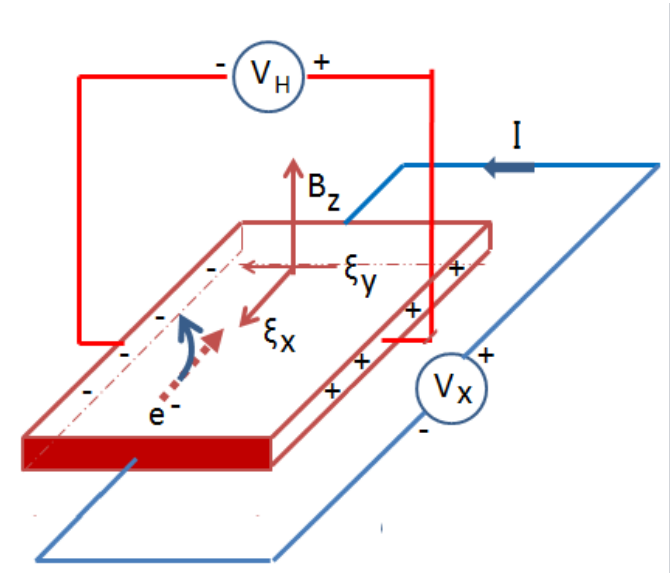
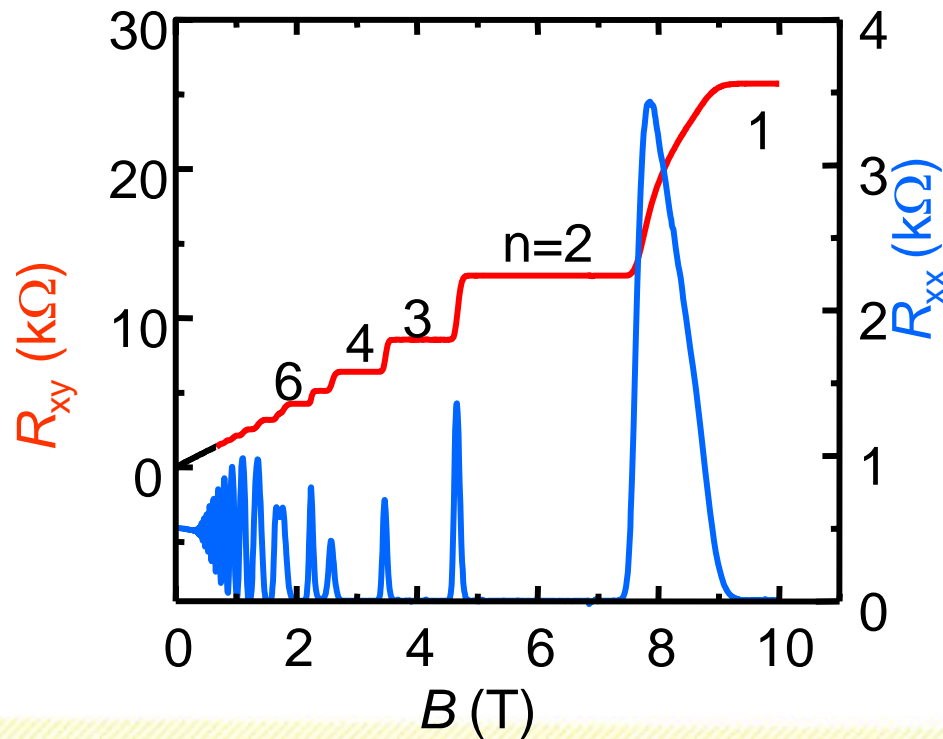
Why we want to have high magnetic fields?

- To put materials in extreme conditions

Hall effect

- Free electrons in a magnetic field.

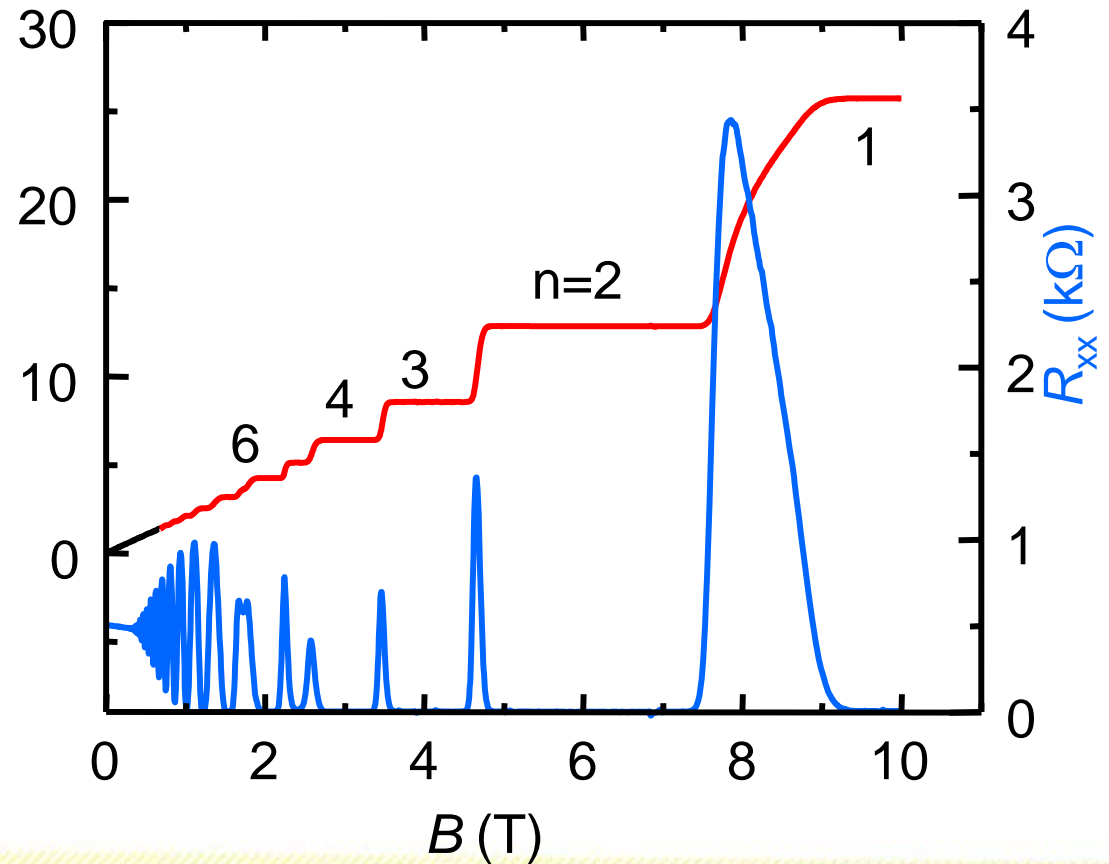
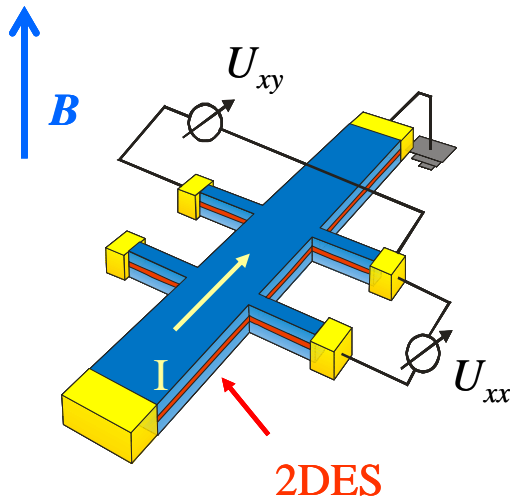
$$F_L = q\mathbf{v} \times \mathbf{B} \quad V_H = \frac{1}{ne} jB$$



Quantum Hall effect

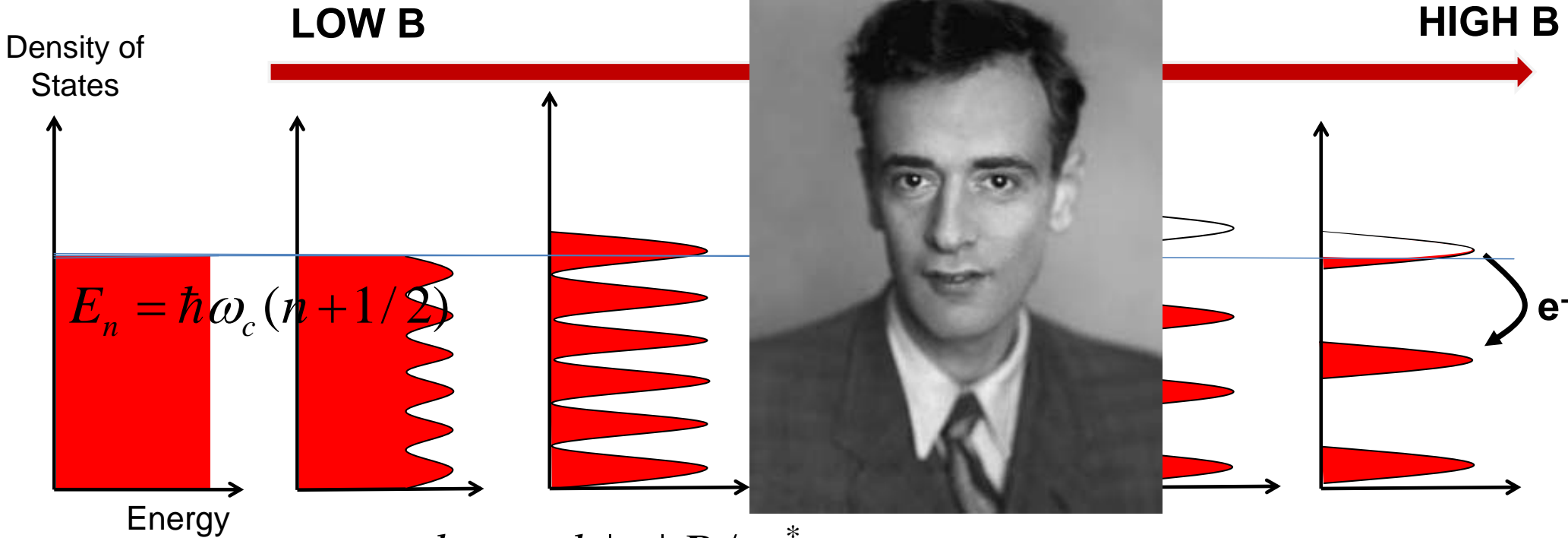


The Nobel Prize in Physics 1985
Klaus von Klitzing



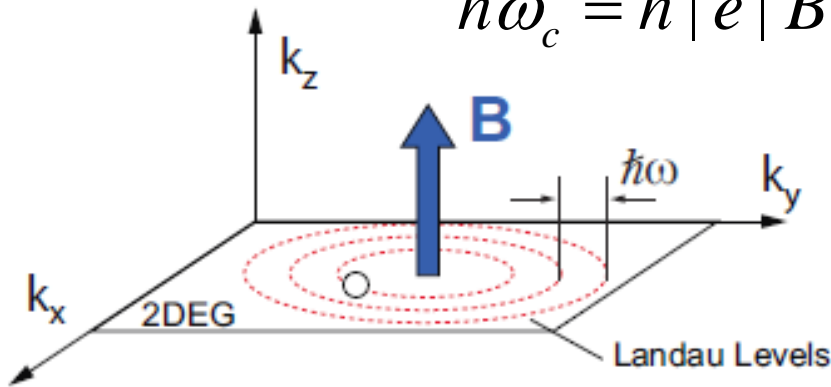
- $\Rightarrow 1/B$ -periodic oscillations in R_{xx}
- \Rightarrow quantized plateaus in $R_{xy} = (h/ne^2)$
- $\Rightarrow R_{xx} \rightarrow 0$ in quantum Hall plateaus

Landau quantization



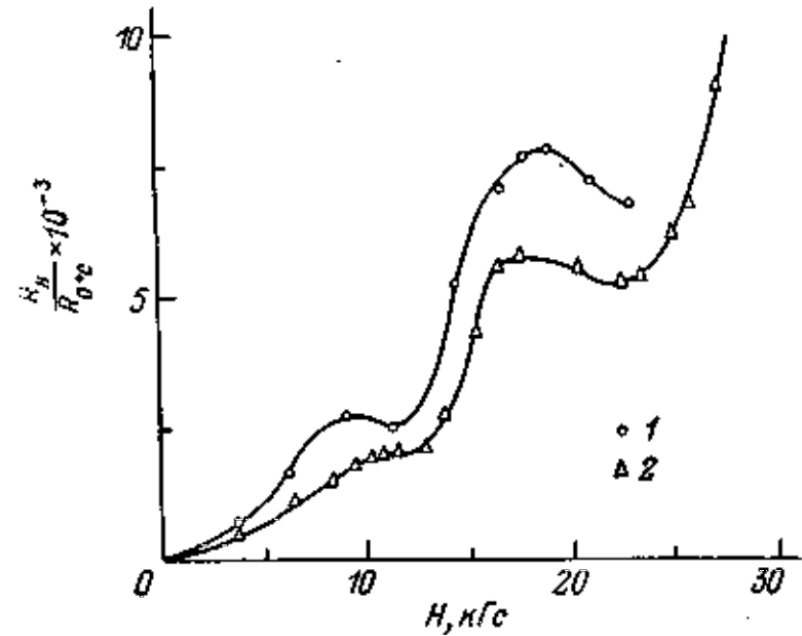
$$\hbar\omega_c = \hbar |e| B / m_c^*$$

Periodic $1/B$ Oscillations in the sample properties:
Electrical Resistivity,
Magnetisation etc



Shubnikov de Haas oscillations

Shubnikov - de Haas oscillations



1930

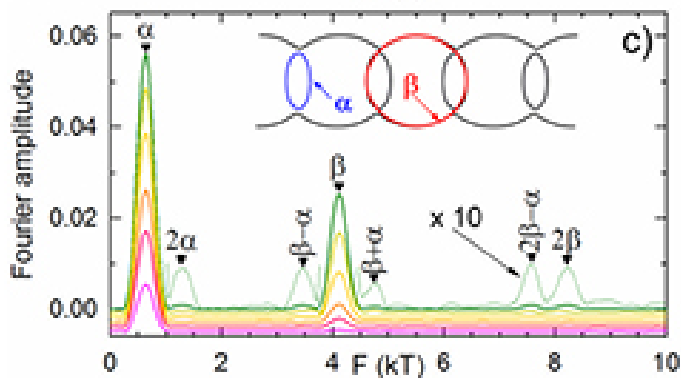
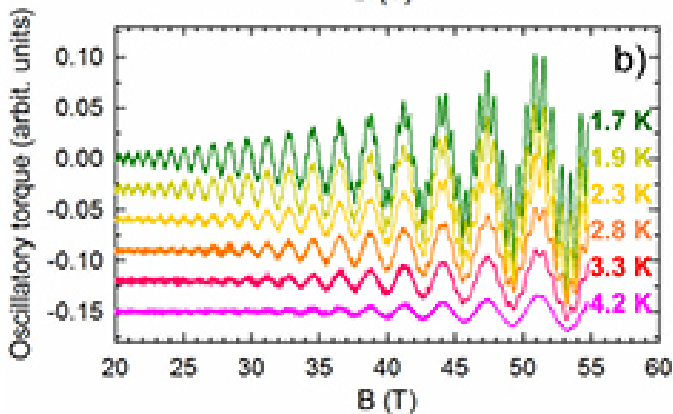
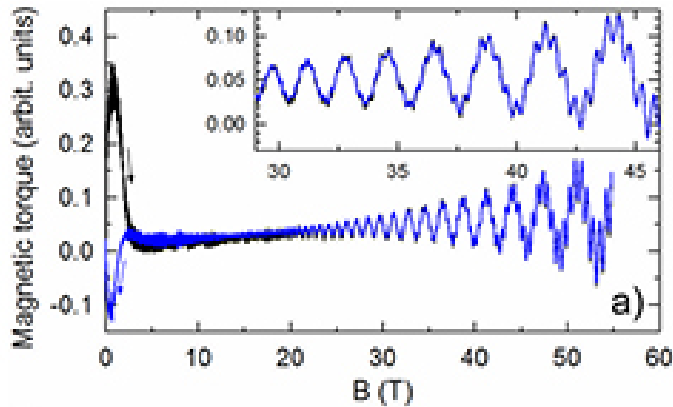
$$\Delta\rho_{xx} = \sum_{i,p} R_{ip} \exp(-\alpha_i T_{D_i}) \frac{\alpha_i T}{\sinh(\alpha_i T)} \sin\left(\frac{2\pi p f_i}{B} + \frac{\pi}{4}\right)$$

Lifshitz – Kosevich formula

Quantum oscillation in organic metal $\kappa\text{-(ET)}_2\text{Cu(SCN)}_2$

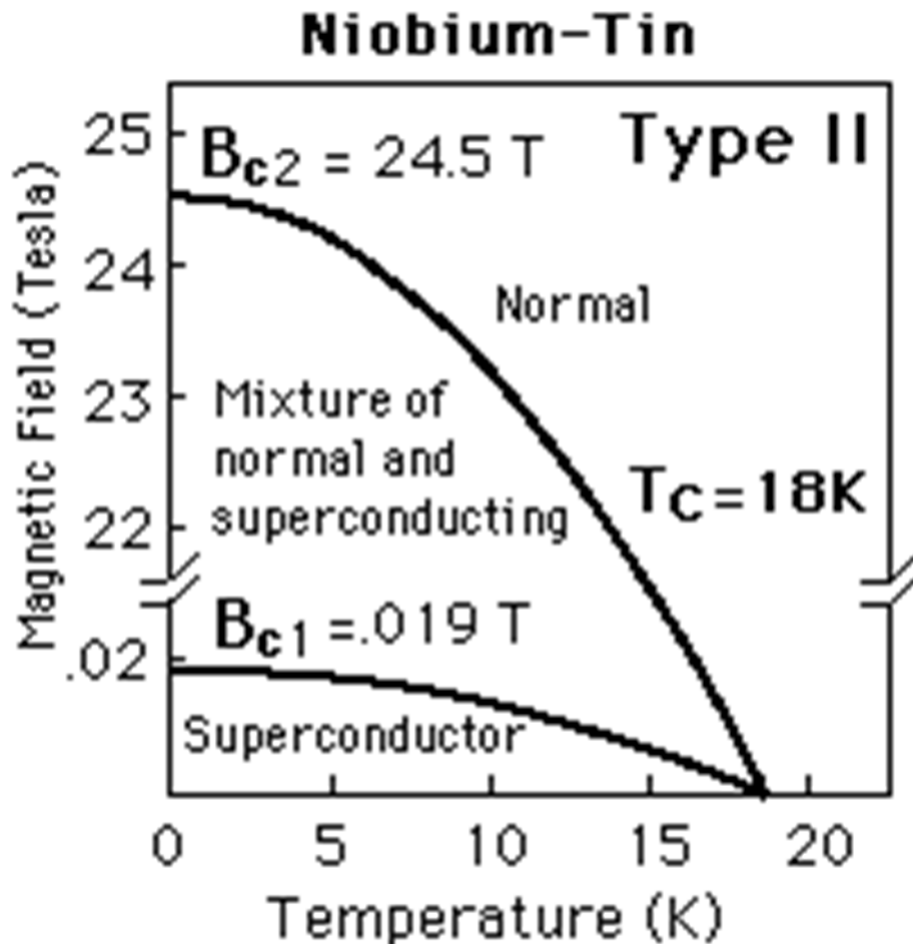
Quantum oscillations provide unique information about Fermi surface of solids

It requires high magnetic fields!



J. Phys.: Condens. Matter **28** (2016) 275702

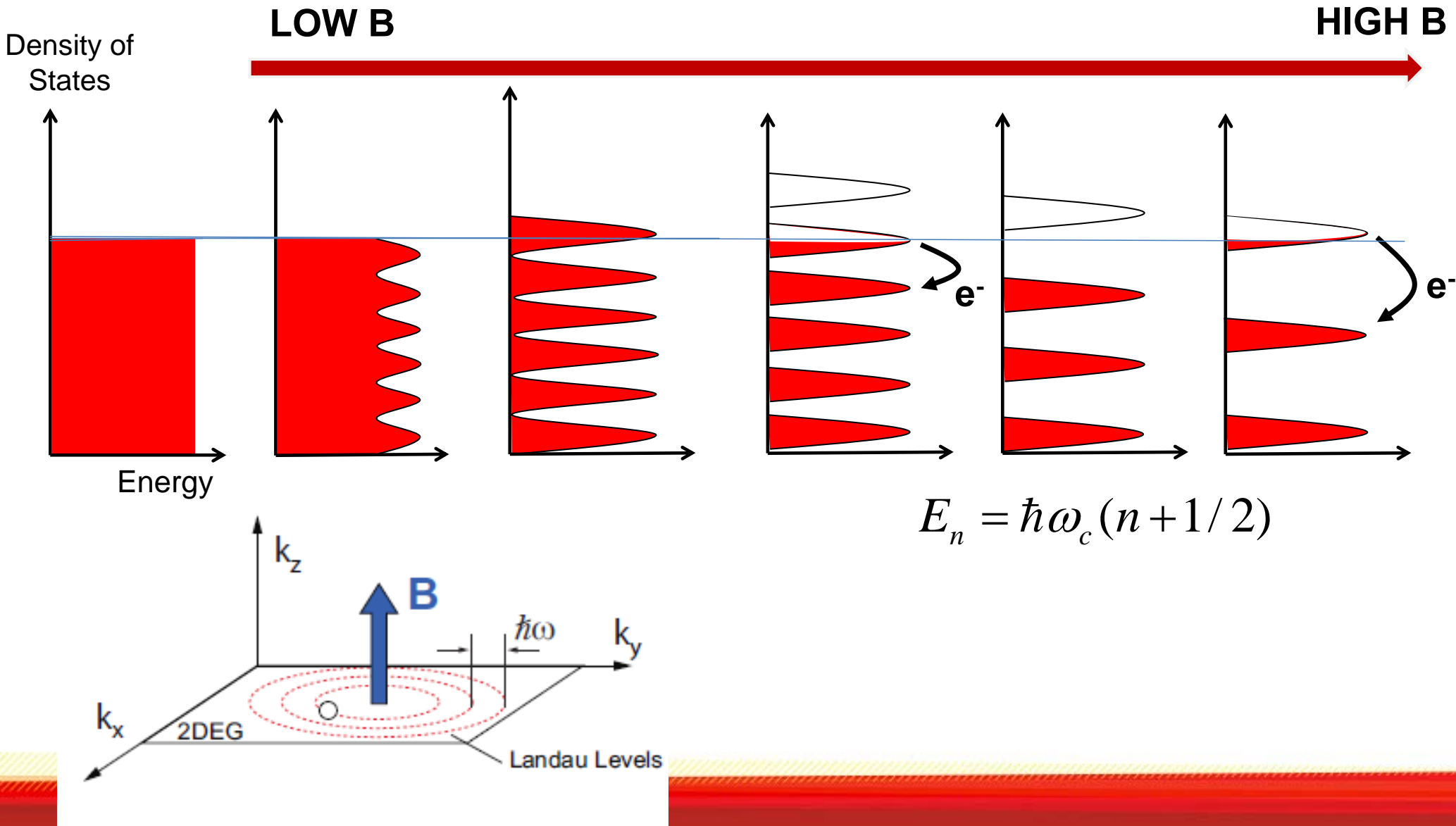
Shubnikov phase in superconductors



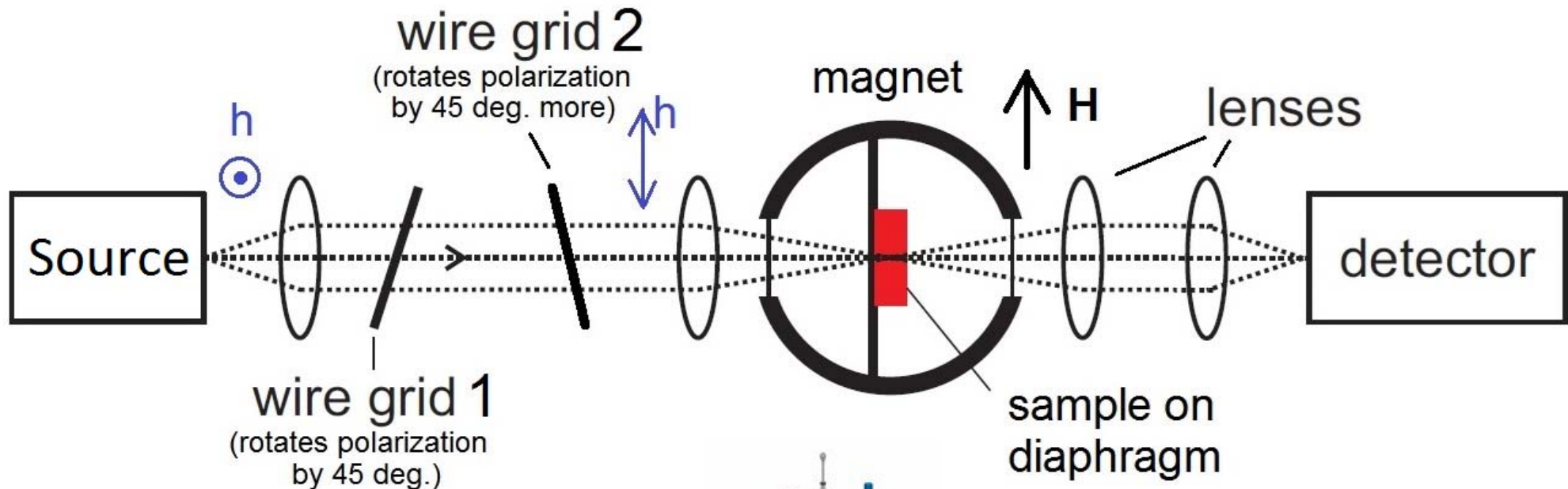
“Thus, in a crystal having 2.5% of thallium at a temperature of 1.92 K, the breakup of the superconductivity occurs in the field interval of 734–1110 Gs. At the same time, H_k only slightly differs from the values for pure lead becoming slightly smaller than these values when the temperature decreases.”

L.V. Shubnikov, V.I. Khotkevich,
Y.D. Shepelev, Y.N. Ryabinin
Zh. Eksper. Teor. Fiz. 7, pp. 221–237 (1937)

Landau quantization



How do we measure CR?



Superconducting
Spectromag
B up to 12 Tesla



Two types of experiments:

- constant B , scan ω
- constant ω , sweep B

Solids from spectroscopy point of view

Metals – a lot of electrons.

Reflects the radiation below plasma frequency:

$$\omega_p = \sqrt{\frac{n_e e^2}{m \epsilon_0}} \text{ – ultraviolet for the metals}$$

Semiconductors – have some charge carriers.

The properties of the carriers could be studied by the cyclotron resonance.

Dielectrics – no charge carriers.

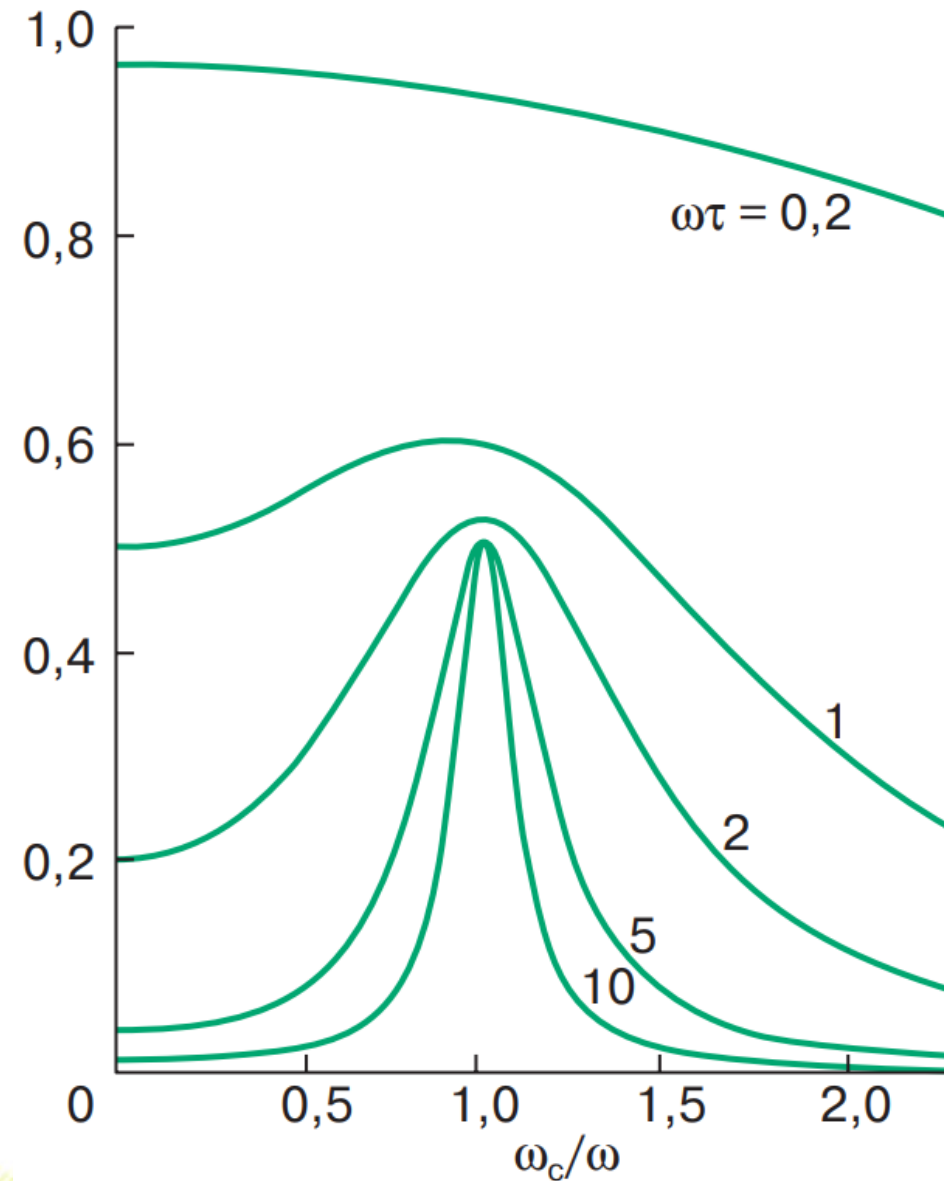
The properties of the magnetic ions or impurities could be studied by means of paramagnetic resonance

Practical requirements

- \mathbf{E}^ω perpendicular to \mathbf{H}
- $\omega_c \tau > 1$
 τ - carrier scattering lifetime
electron can complete its
cyclotron orbit unhindered
- $\hbar \omega_c > k_B T$
 - To have population differences
between levels

$$\tau \sim 10^{-12}$$

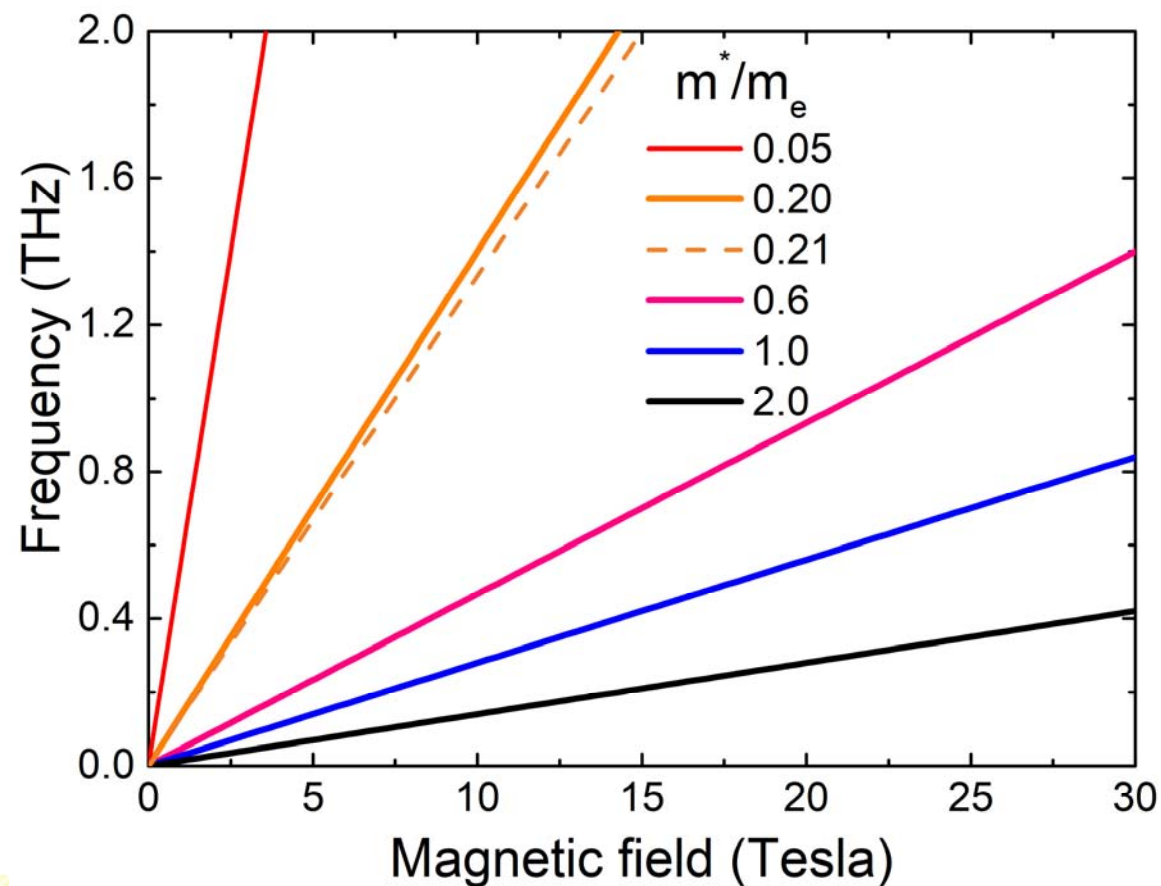
$$\omega_c > 1 \text{ THz}$$



Estimation of energies

ω_c - What frequencies we are talking about?

$$\omega_c = |e| B / m_c^*$$

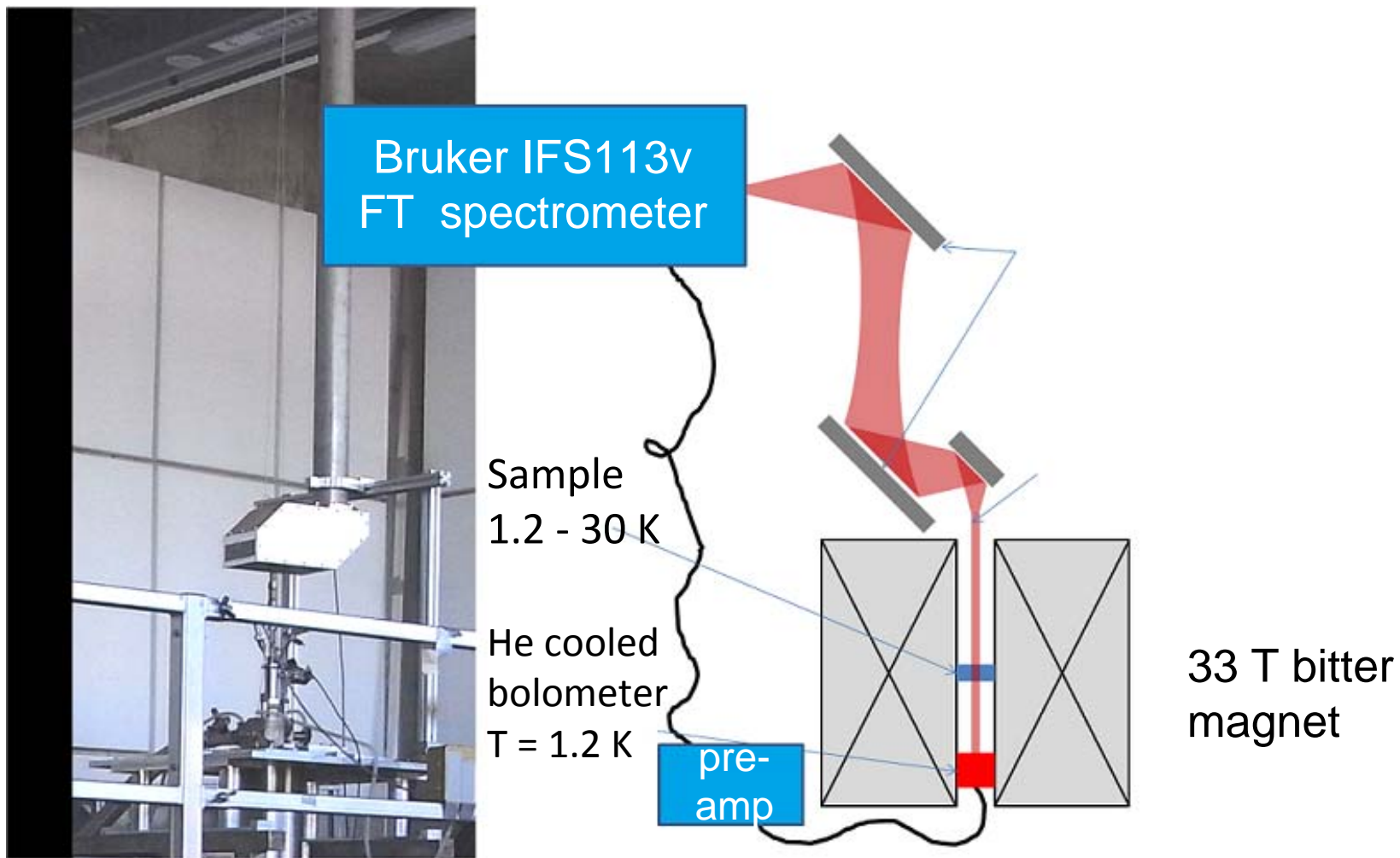


Crystal	Electron	Heavy Hole	Light Hole	Split-off Hole
InSb	0.015	0.39	0.021	0.11
InAs	0.026	0.41	0.025	0.08
InP	0.073	0.4	0.078	0.15
GaSb	0.047	0.3	0.06	0.14
GaAs	0.066	0.5	0.082	0.17
Cu ₂ O	0.99	-	0.58	0.69

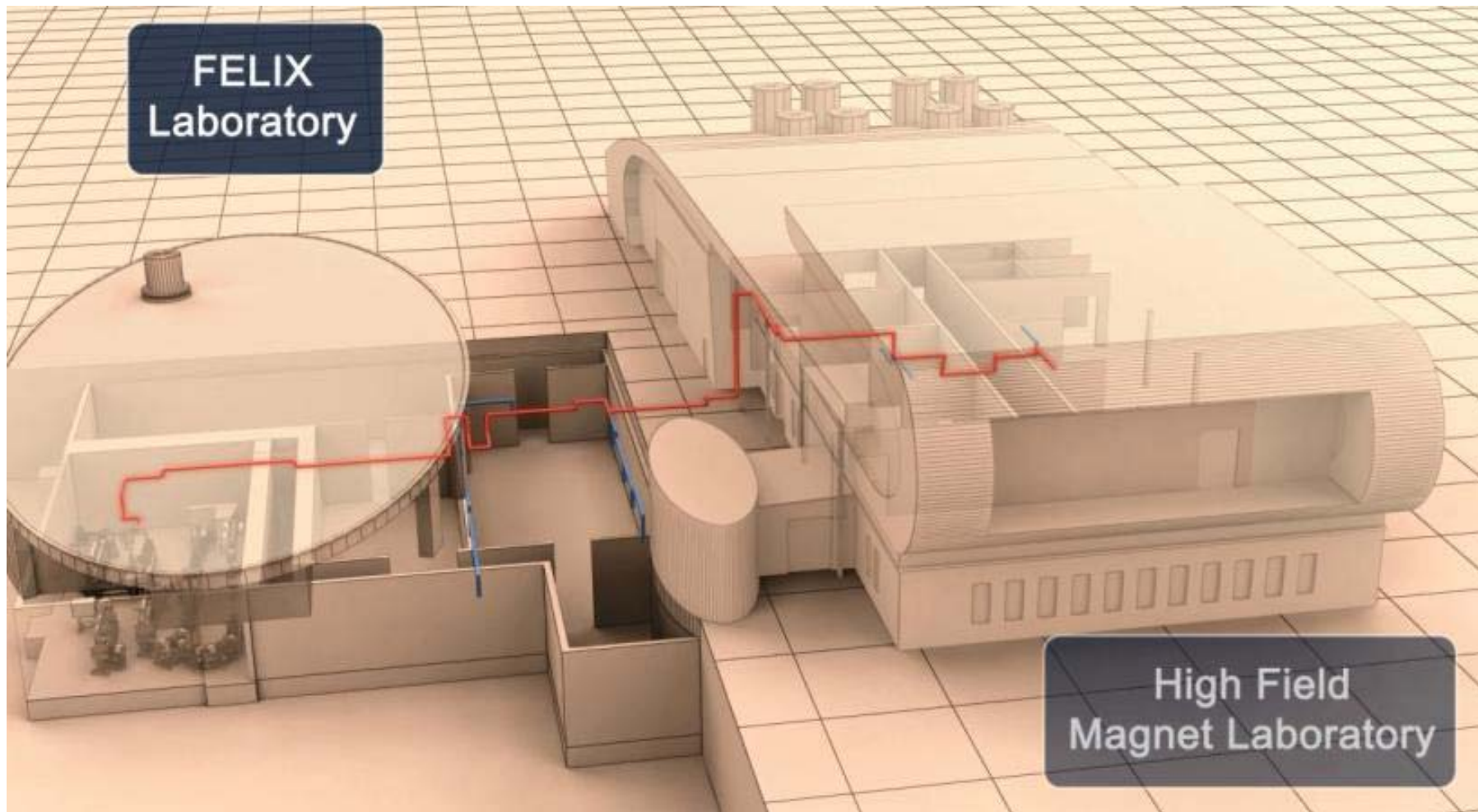
$$\omega_c \tau > 1,$$

$$\tau \approx 10^{-12} \text{ s} \rightarrow \omega_c > 1 \text{ THz}$$

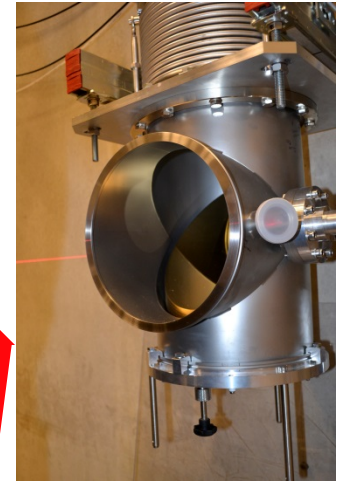
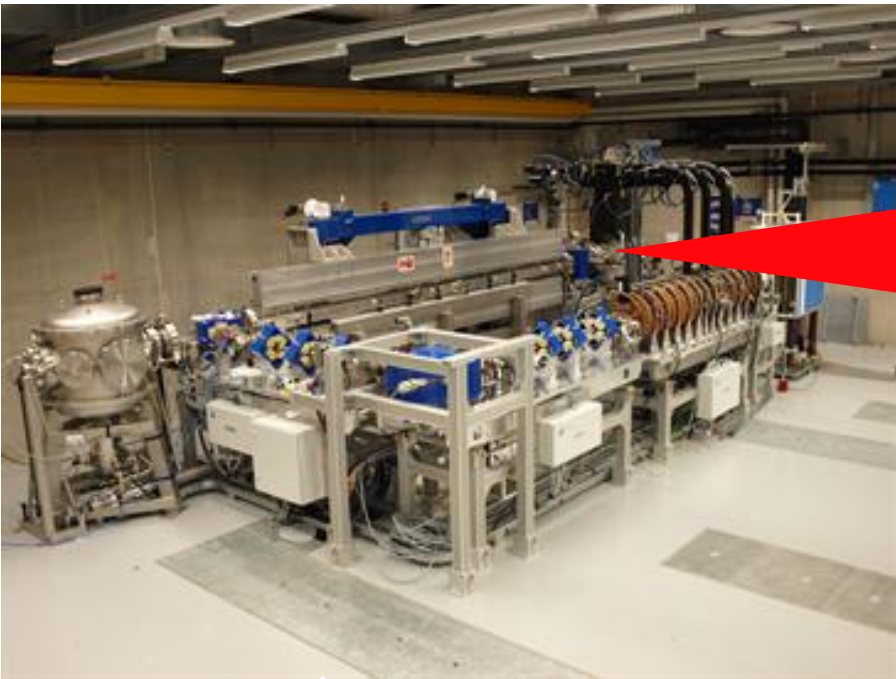
High field FT-FIR transmission setup



Coupling between FELIX and HFML

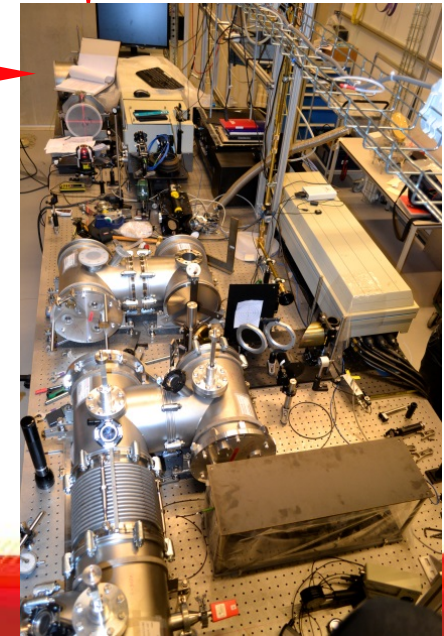


FELIX – HFML beamline

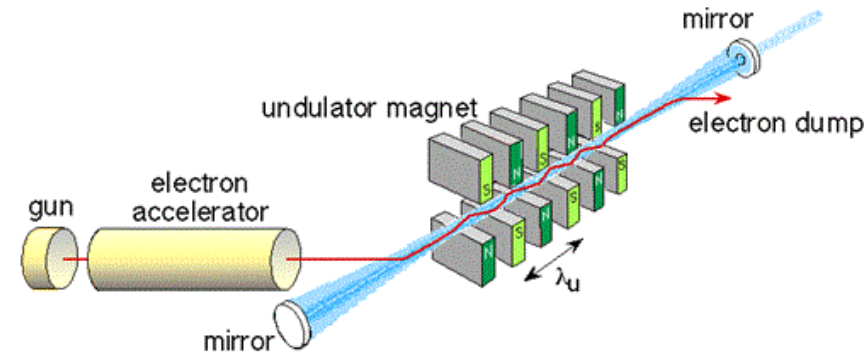
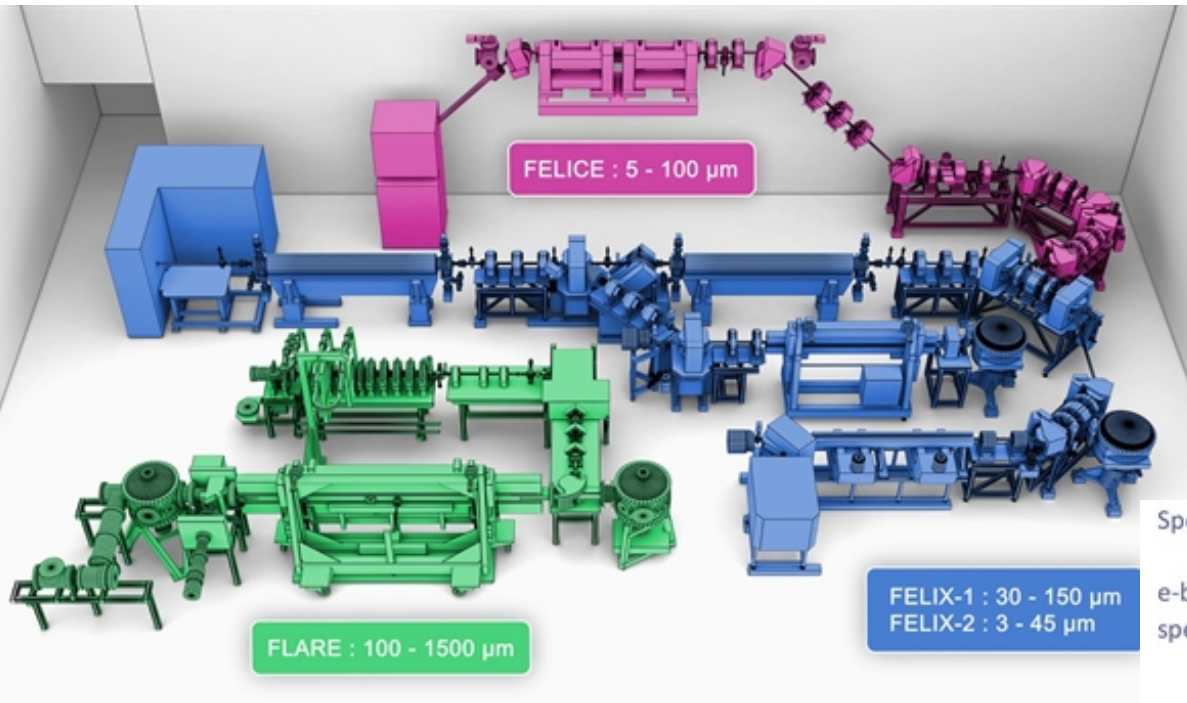


42 mirrors
90 m beamline

220 mWatt



FELIX facility



Specs:

e-beam energy
spectral range

pulse structure
rep. rate
micropulse energy
macropulse energy
peak power
polarisation
spectral bandwidth
(FWHM)
continuous tunability

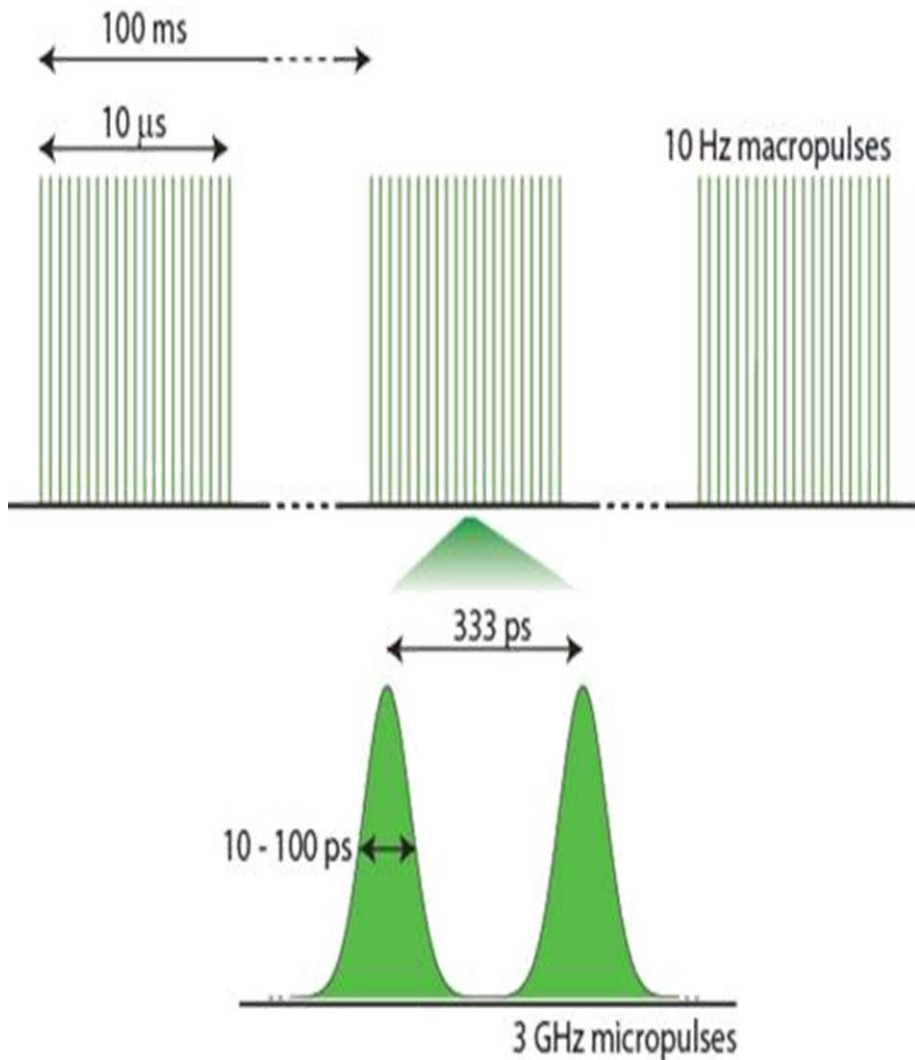
FELIX:

50/45 – 15 MeV
2.7 – 150 micron
3600 - 66 cm⁻¹
120 – 2 THz
450 – 8 meV
micro / macropulse
25 MHz/1 GHz@10 Hz
1-20 μJ
≤ 100 mJ @ 1 GHz
≤ 100 MW
linear
0.2 - 5%
200 - 300 %

FLARE:

15 – 10 MeV
100 - 1500 micron
100 - 6 cm⁻¹
3 – 0.25 THz
12 – 0.75 meV
micro / macropulse
3 GHz@10 Hz
≈ 5 μJ
≤ 100 mJ @ 3 GHz
≤ 10 MW
linear
≤ 1%
spectral mode ≤ 10⁻⁴
?

FLARE radiation profile



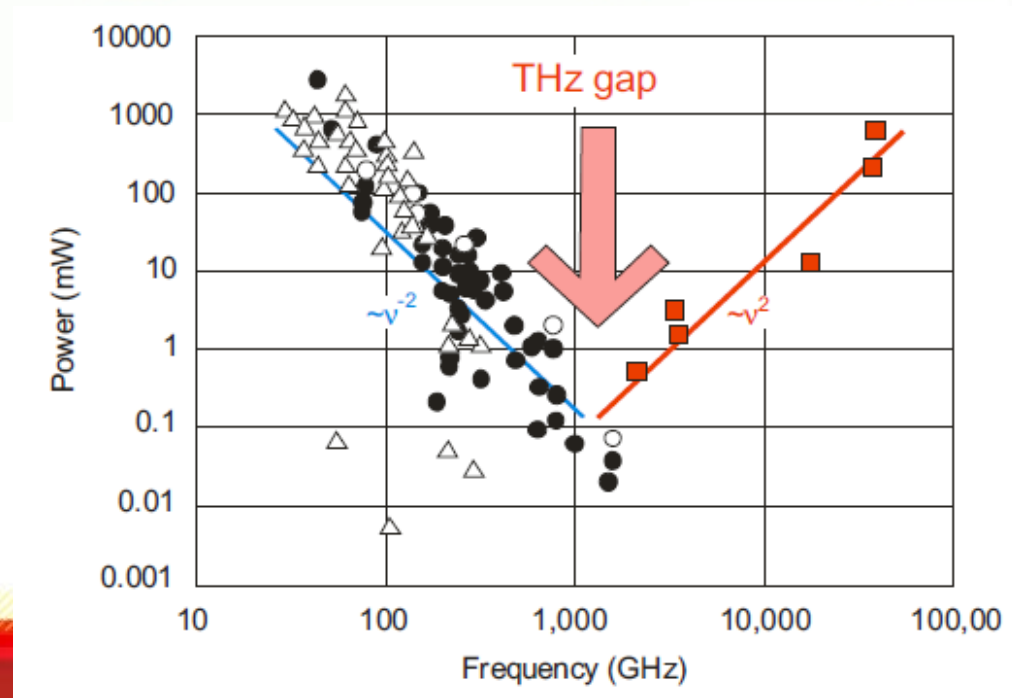
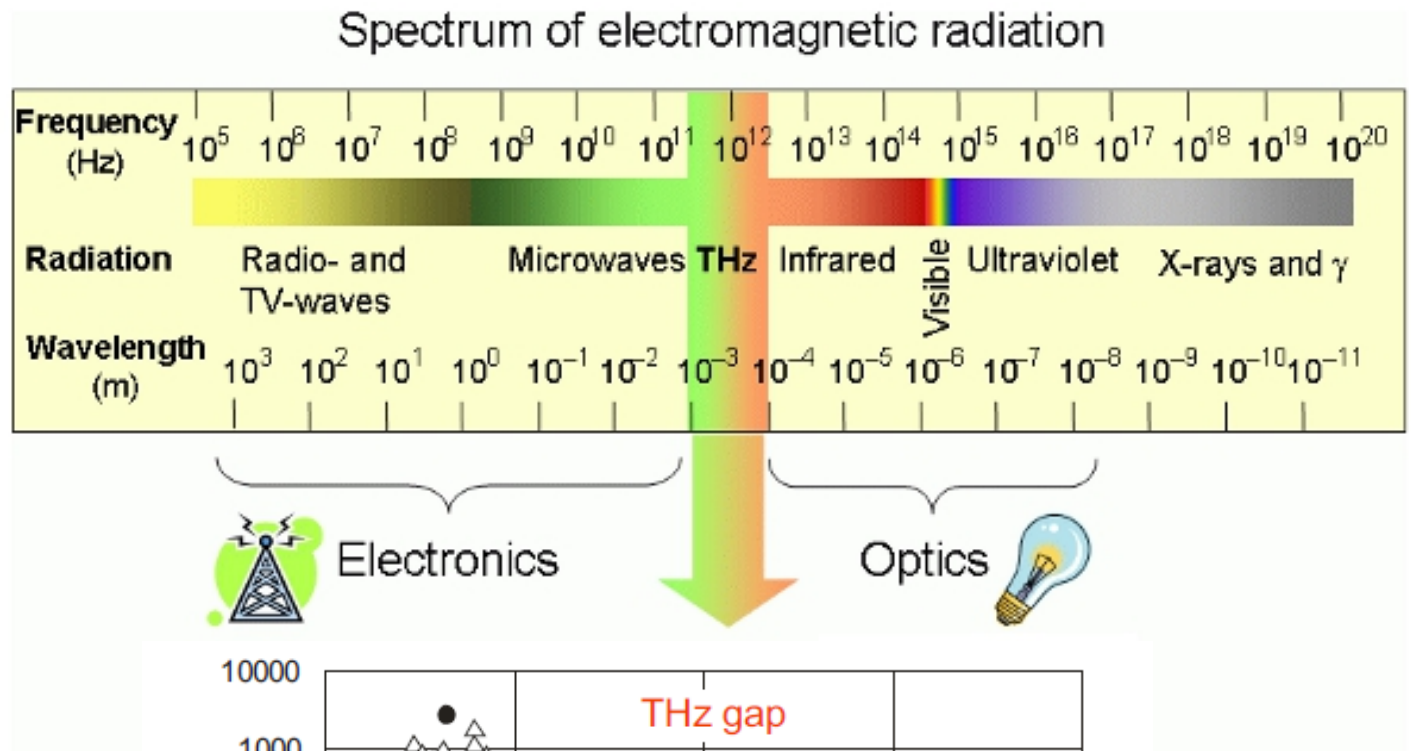
- Frequency range: 7 – 100 cm⁻¹ (0.2 – 3 THz)
- Average power up to 1W
- photon flux ~ 10²¹ photons/second

Macropulse ≈ 90mJ

Duty cycle 10 μs/100 ms = 10⁻⁴
10kW during macropulse
i.e. 10²⁰ photons

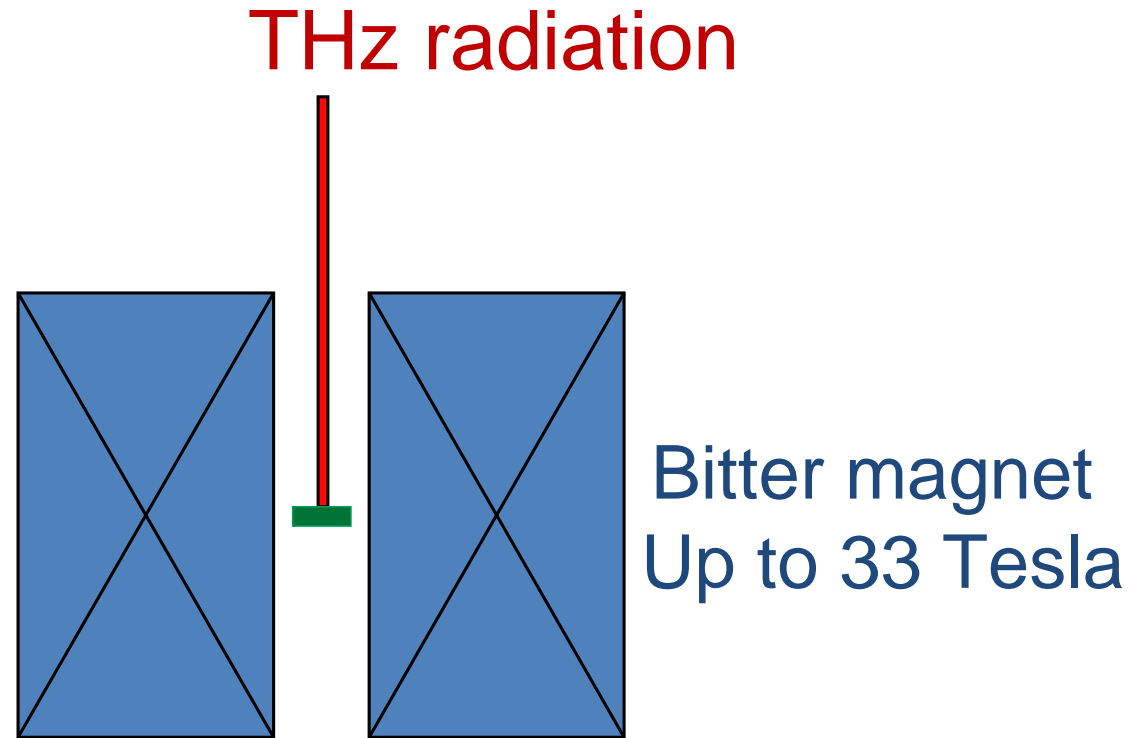
Micropulse ~50ps, 3μJ
~10¹⁶ photons per pulse
~100kW during a
micropulse

THz gap



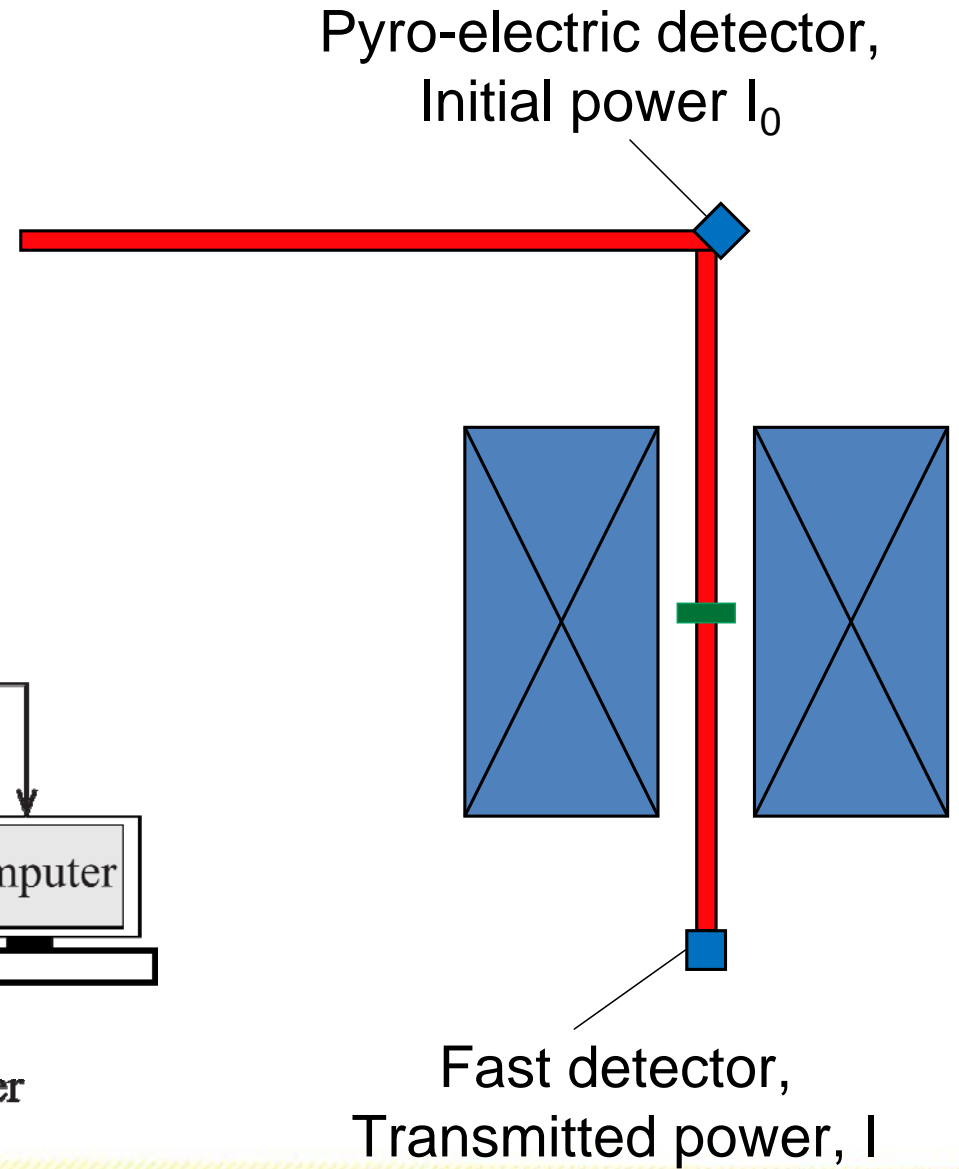
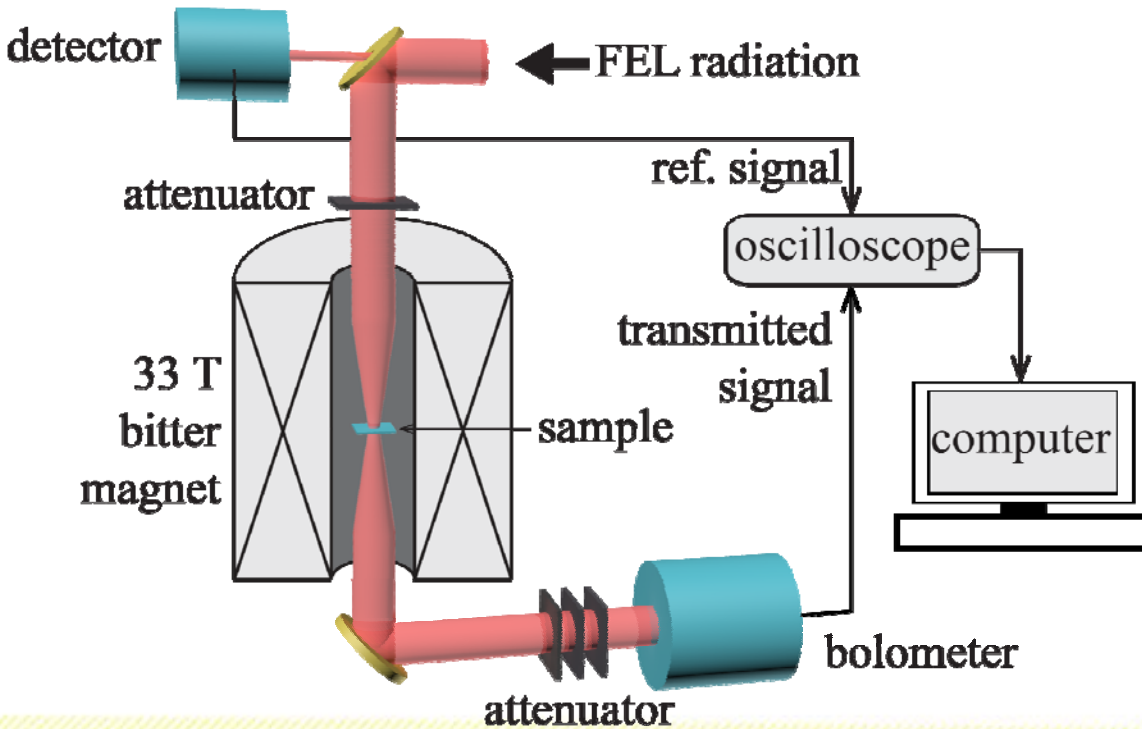
FELIX + HFML = unique opportunities!

- Magnetotransmission over unique ranges of fields, frequencies, powers
- Pump-probe experiments to study materials in unique environment

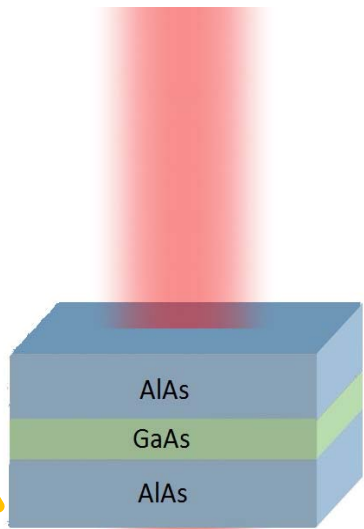


Transmission measurement

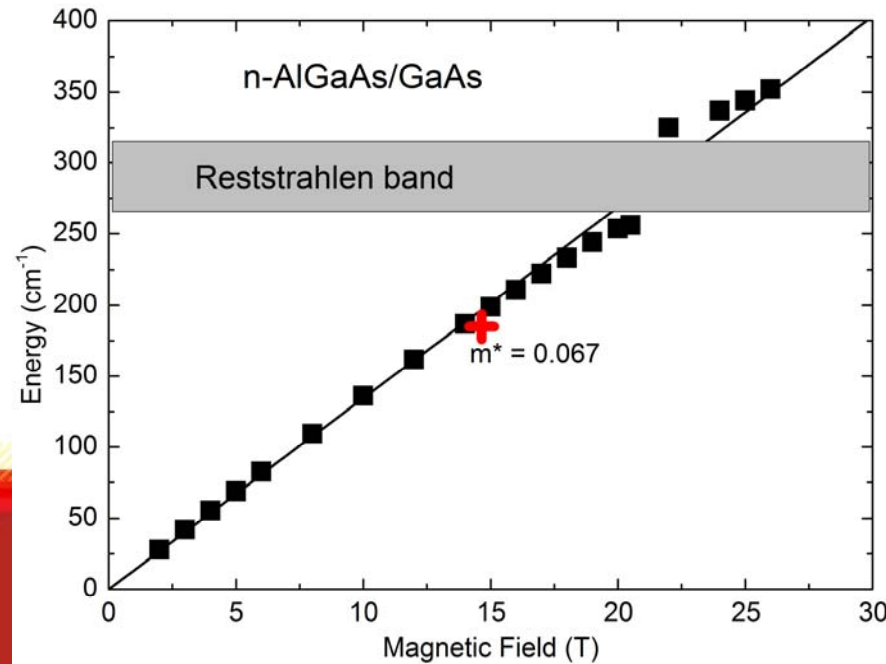
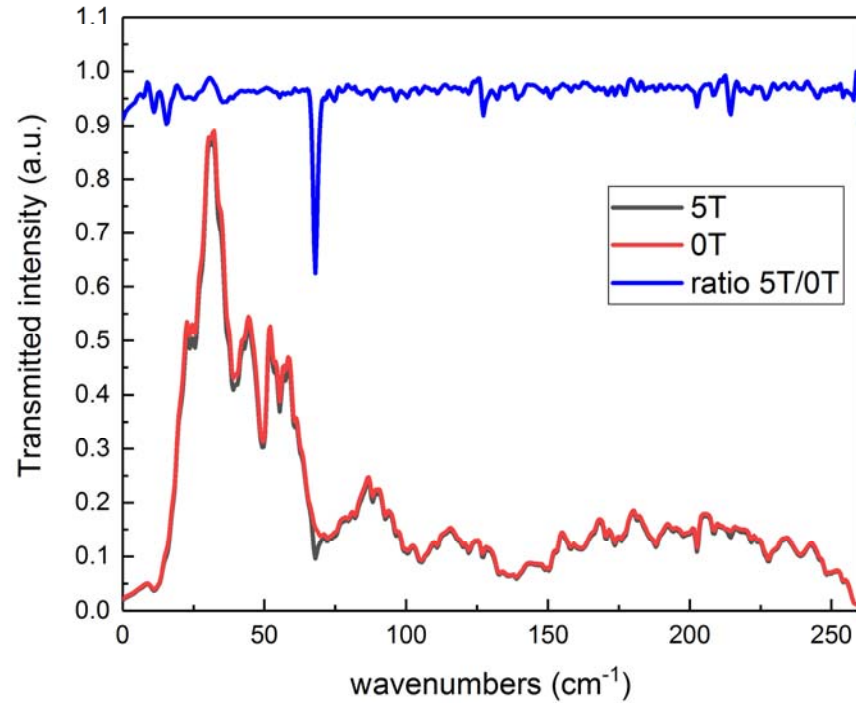
- 1 - 10 Hz macropulse,
- 10 μ s duration:
- Duty cycle $\sim 10^{-4}$
- Lock-in no good:
Use oscilloscope



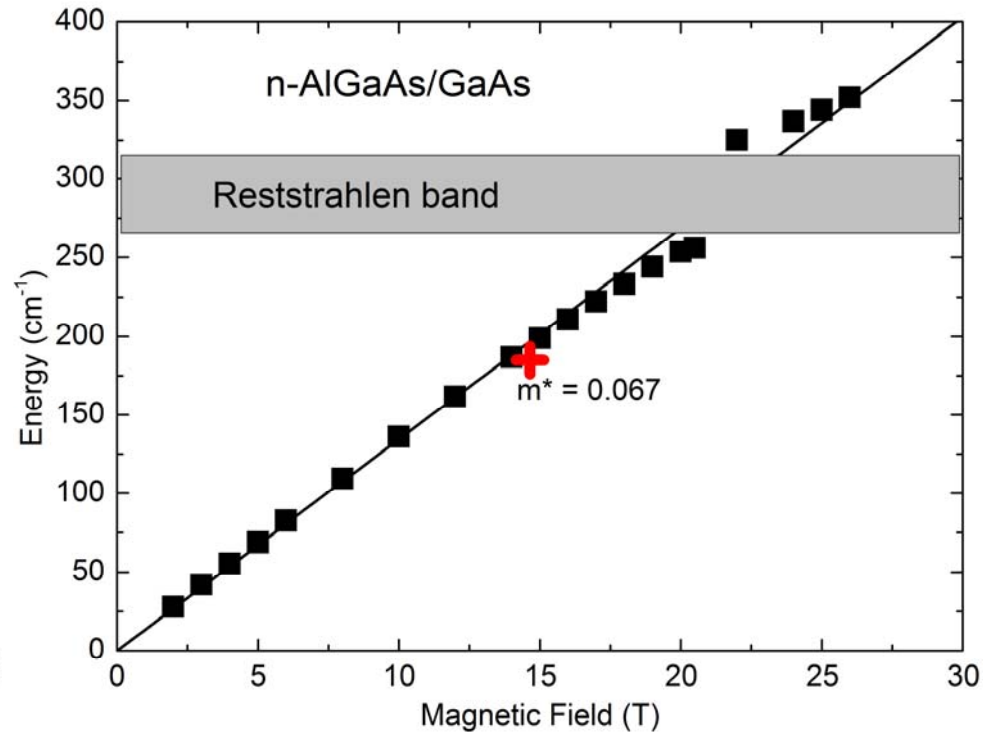
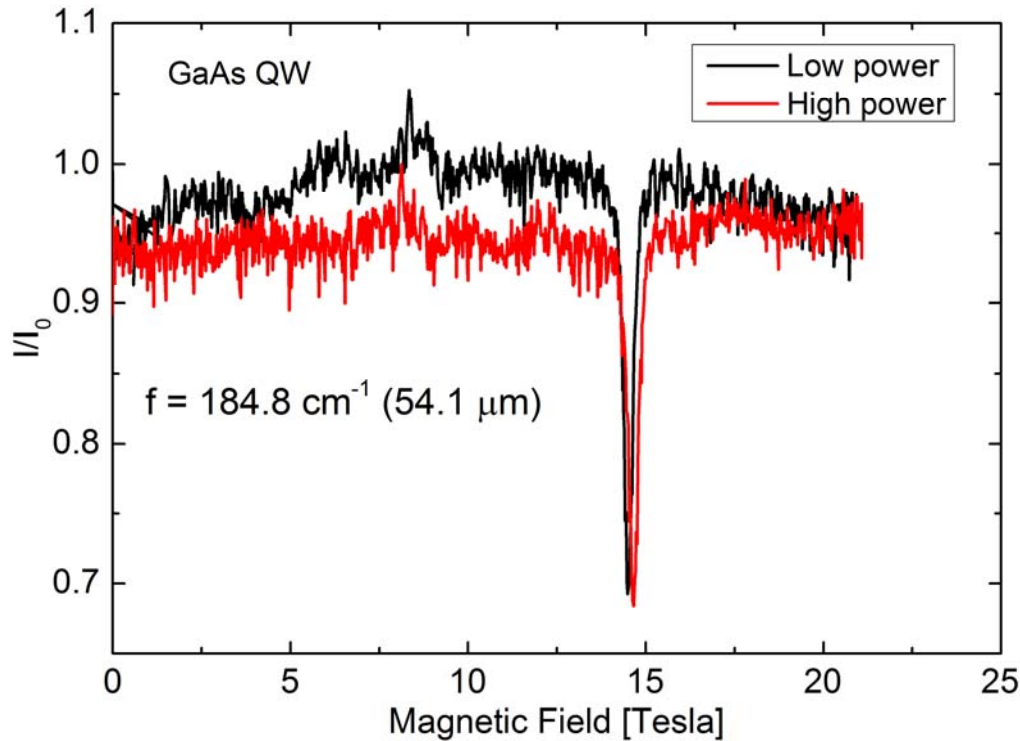
CR experiments on GaAs QW



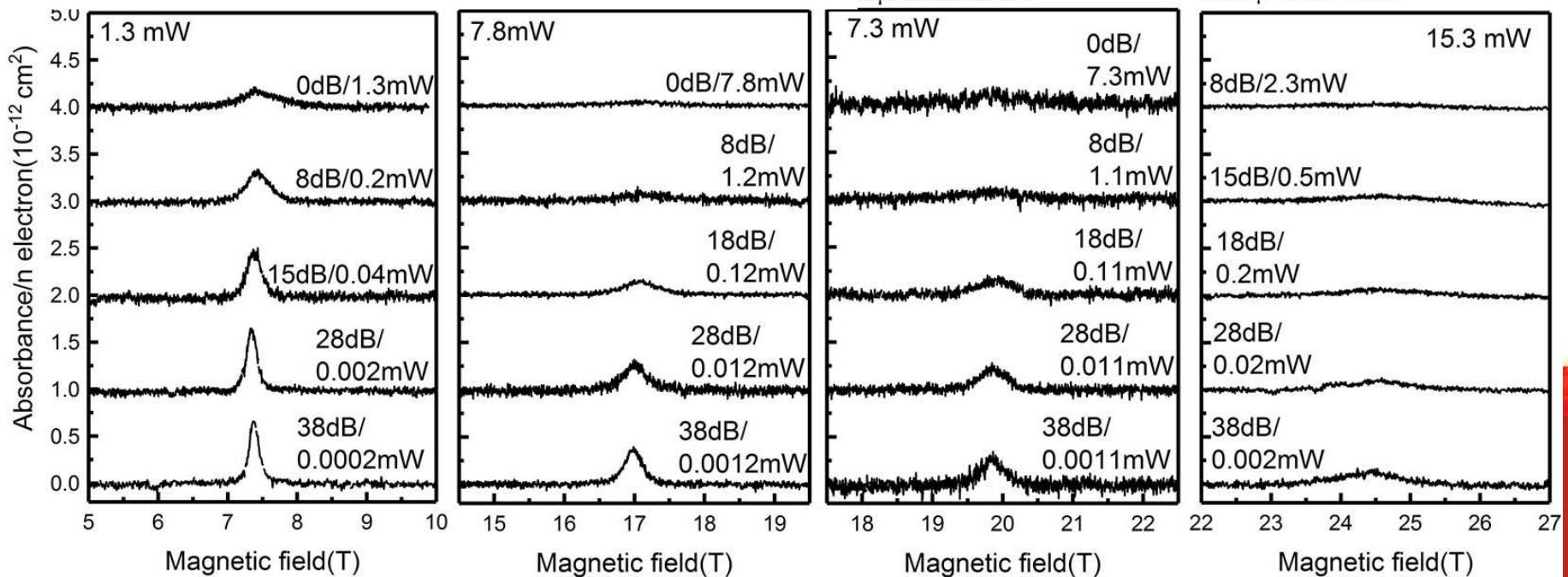
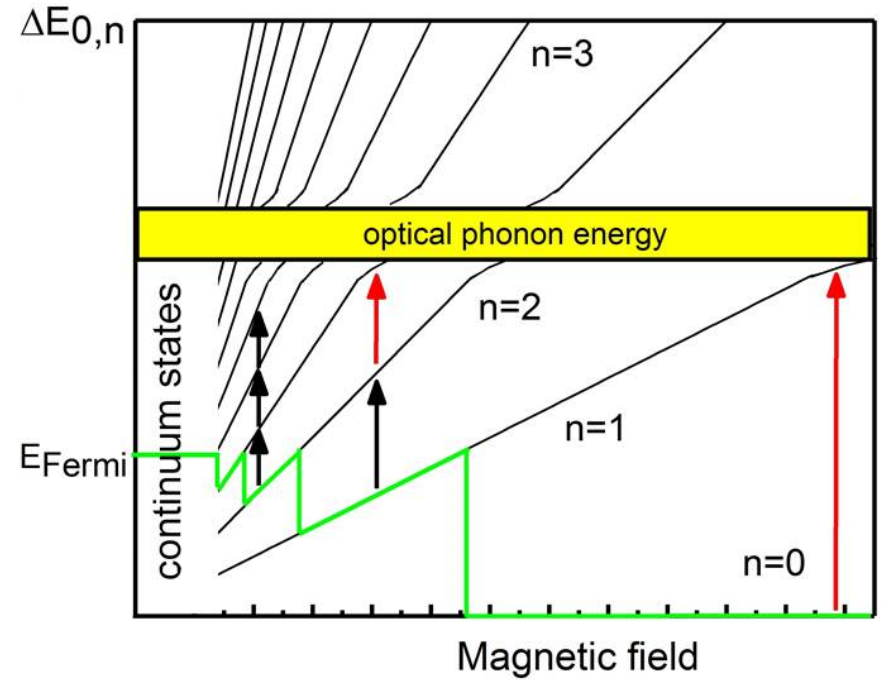
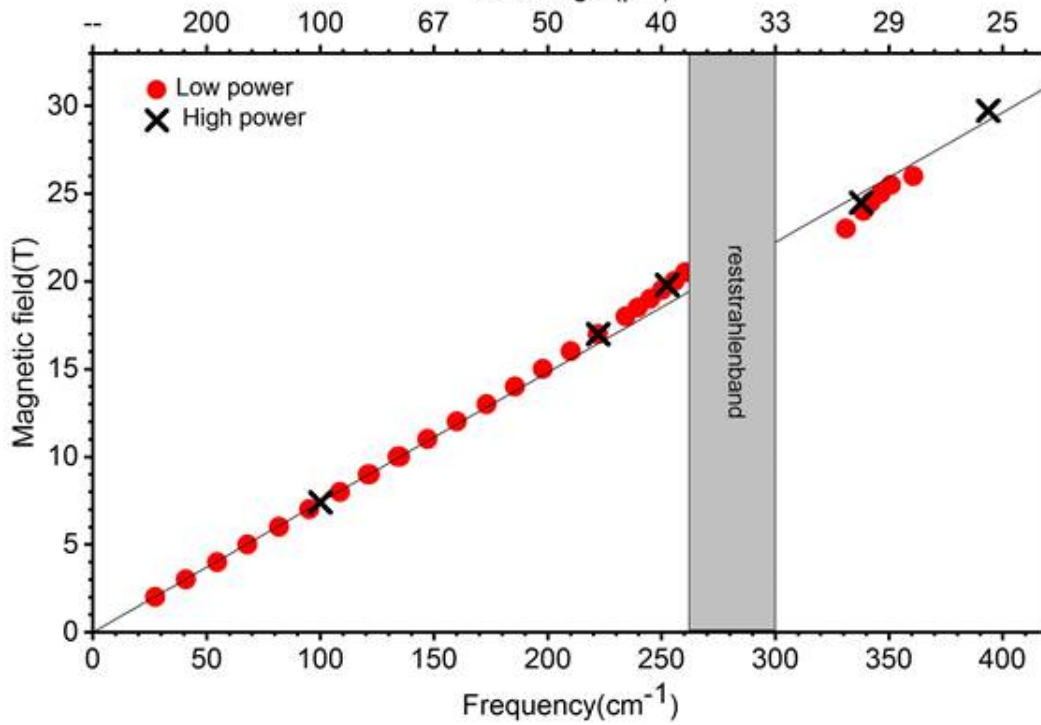
$$\omega_c = |e| B / m_c^*$$



Experiments with free electron laser

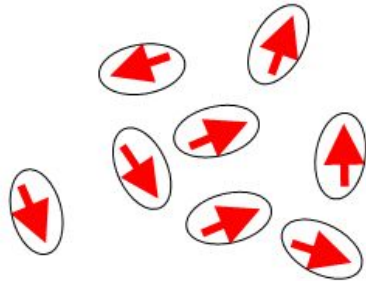


High power behavior

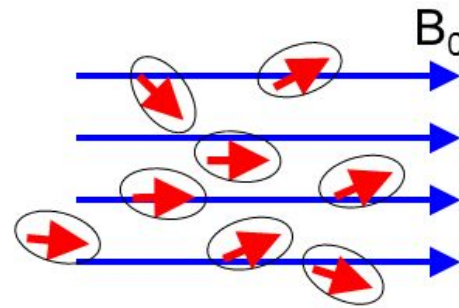


Магнітний момент (спін) у магнітному полі

$B = 0$



$B = B_0$

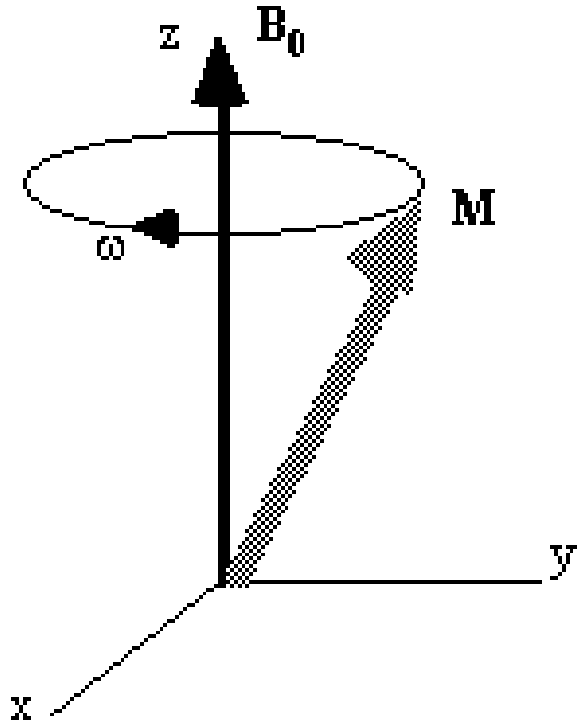


Magnetic particles in magnetic field

$$\frac{d\vec{\mu}}{dt} = \gamma \vec{\mu} \times \vec{B}_0$$

$$\vec{\tau} = \vec{\mu} \times \vec{B}_0$$

$$\left. \begin{aligned} \frac{d\mu_x}{dt} &= \gamma \mu_y B_0 \\ \frac{d\mu_y}{dt} &= -\gamma \mu_x B_0 \\ \frac{d\mu_z}{dt} &= 0 \end{aligned} \right\} \begin{aligned} \mu_x &= A \cos(\omega t + \phi) \\ \mu_y &= -A \sin(\omega t + \phi) \end{aligned}$$



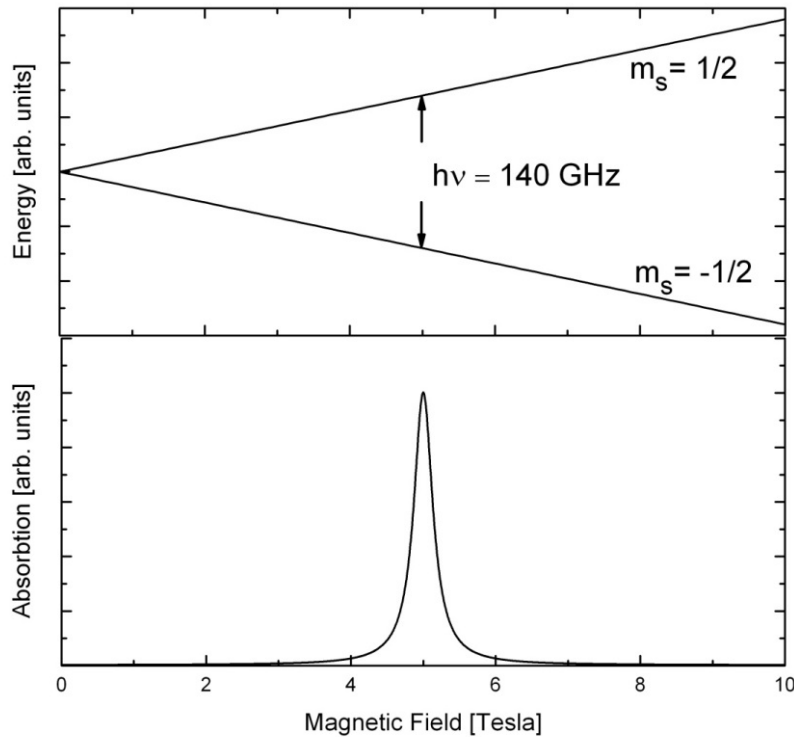
$$\omega = \gamma B_0$$

$$\gamma = \frac{eg}{2m}$$

$$\mu_B = \frac{e\hbar}{4\pi m}$$

$$h\nu = g\mu_B B_0$$

How is it in quantum terms?

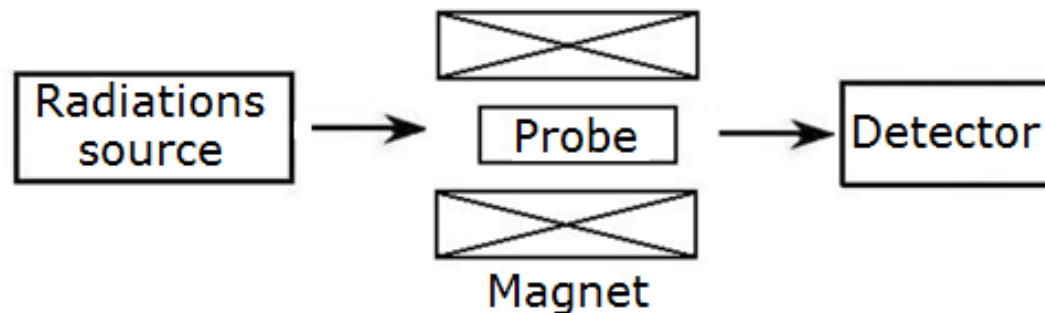


$$\Delta E = \hbar\omega = g\mu_B H$$

μ_B – Bohr magneton.
 g – factor Lande.

Studied parameters:
resonance field,
resonance linewidth,
lineshape.

140 GHz = 5 Tesla



Where are magnetic moments?

- Radicals, spin traps
- Nitroxides, spin labels
- Paramagnetic Transition Metal Ions (*d* and *f* ions)

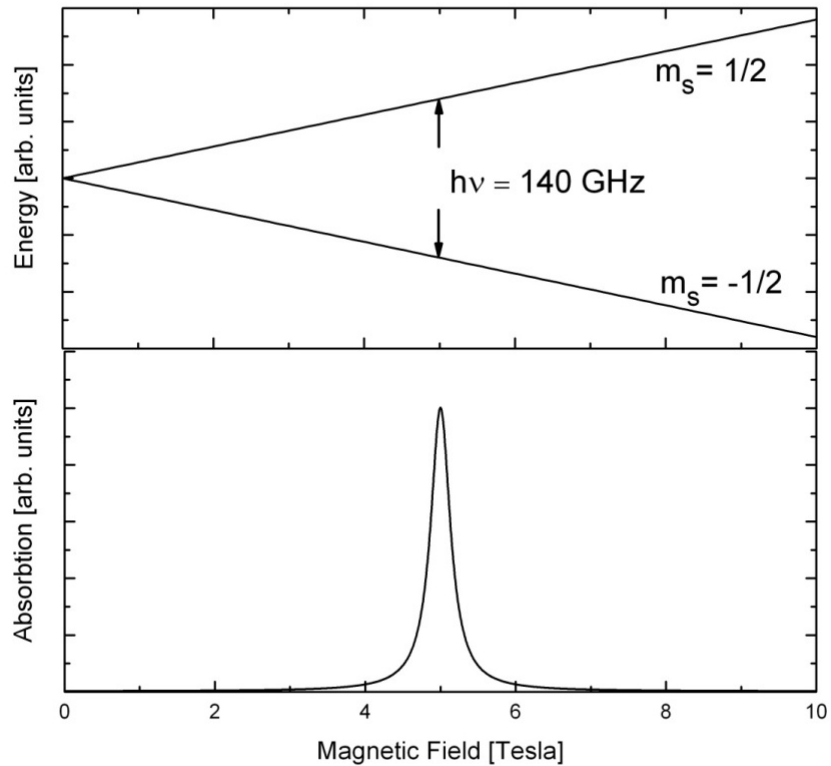
EPR or ESR

- Nuclei with magnetic moments (^1H , ^{14}N , ^{15}N , ^{17}O *etc*)

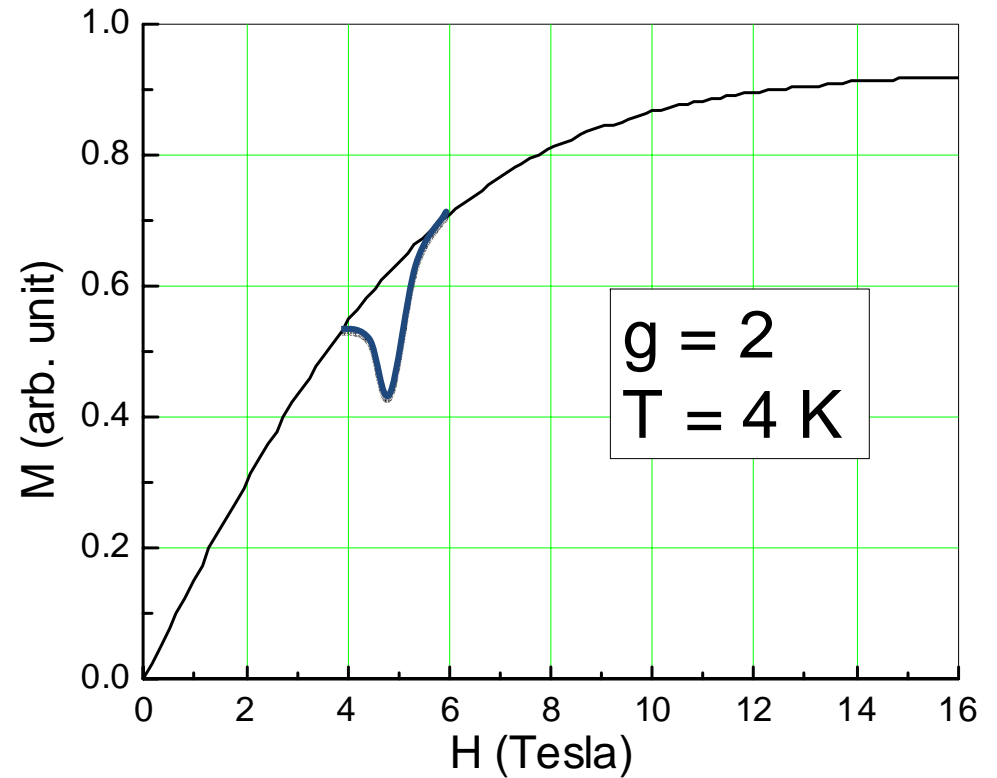
NMR

Magnetically detected EPR

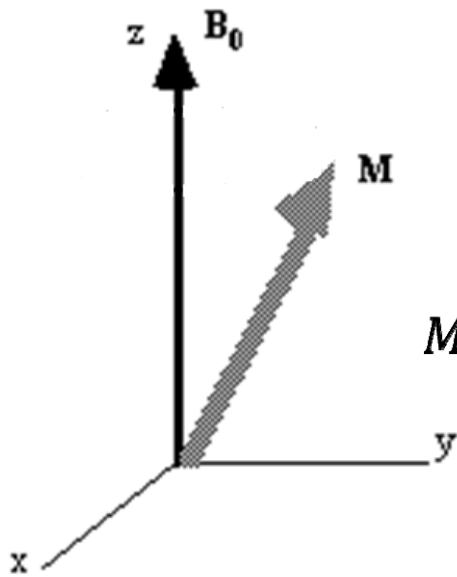
Energy levels



Magnetization

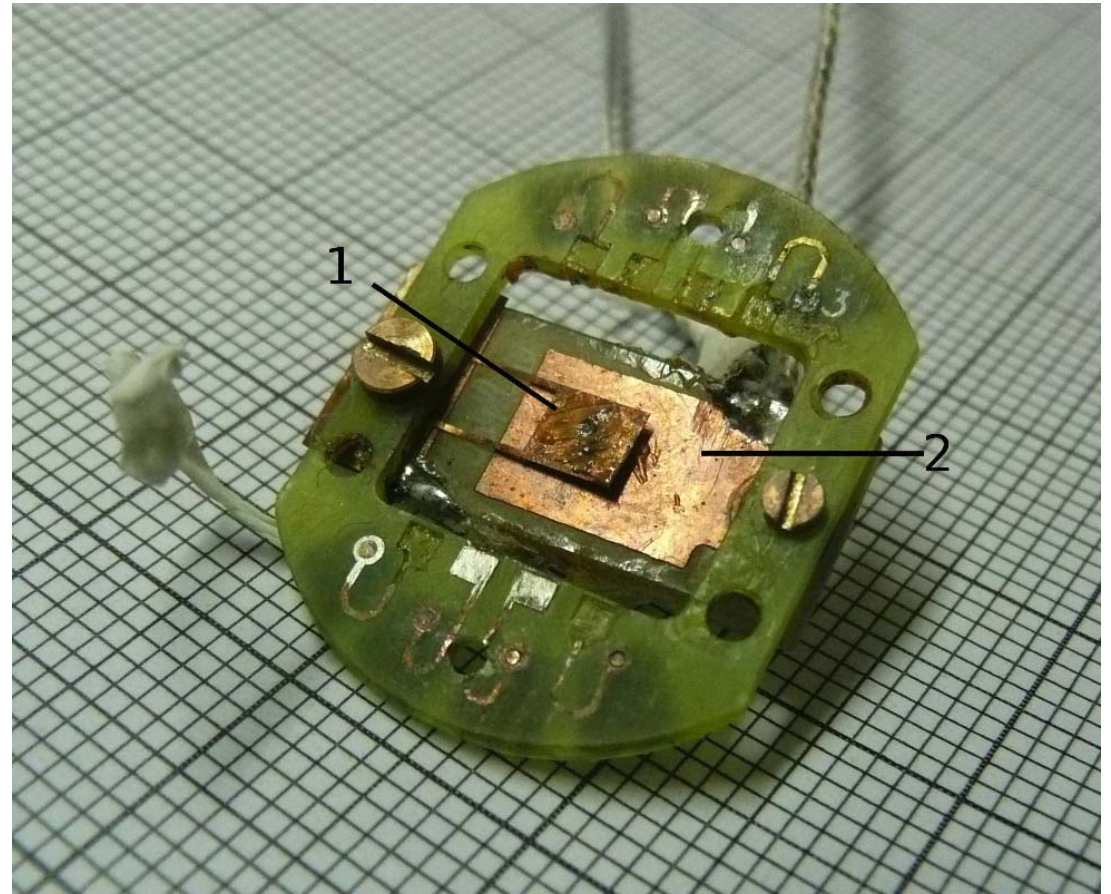


How to measure magnetization



$$\vec{\tau} = \vec{M} \times \vec{B}_0$$

$$M \sim B \rightarrow \tau \sim B^2$$



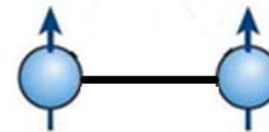
- + Sensitivity increase with B fields
- + Transparency does not matter!

- Anisotropy required!
- Relatively low sensitivity @ fields about 2-3 Tesla

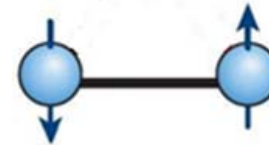
Spin – spin exchange interaction

$$E = J \vec{S}_1 \cdot \vec{S}_2$$

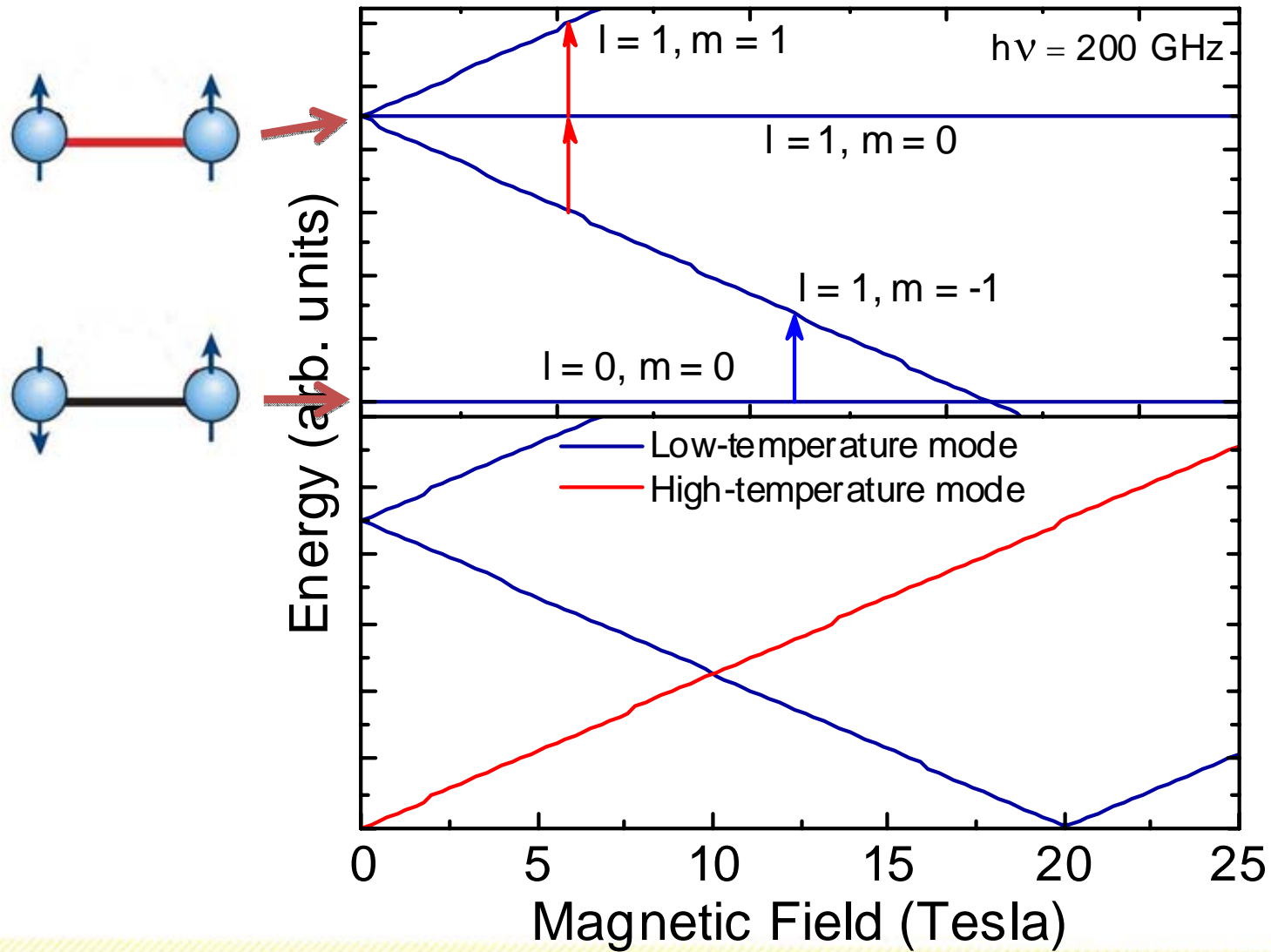
$J > 0$ Ferromagnetic exchange



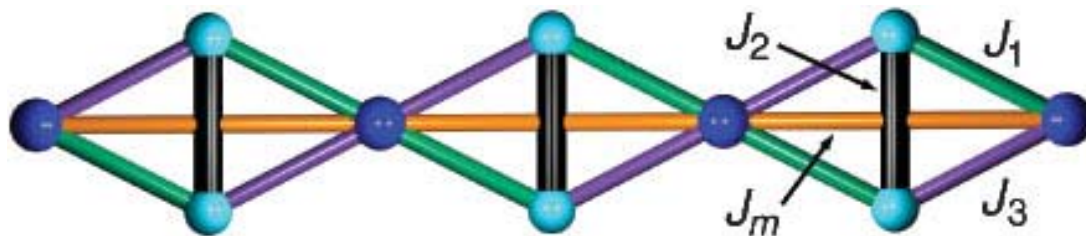
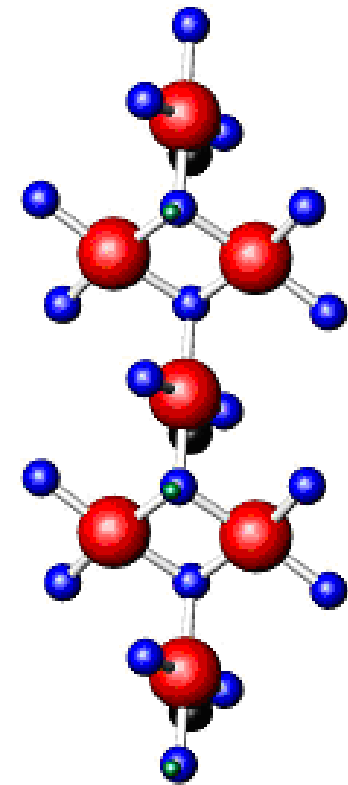
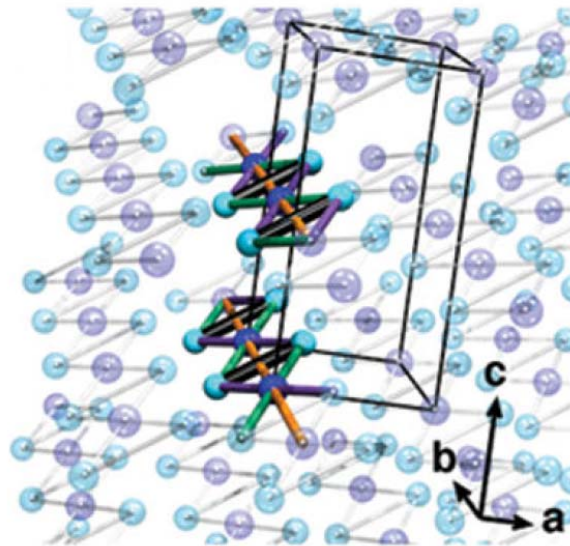
$J < 0$ Antiferromagnetic exchange



Energy diagram of spin-dimer system

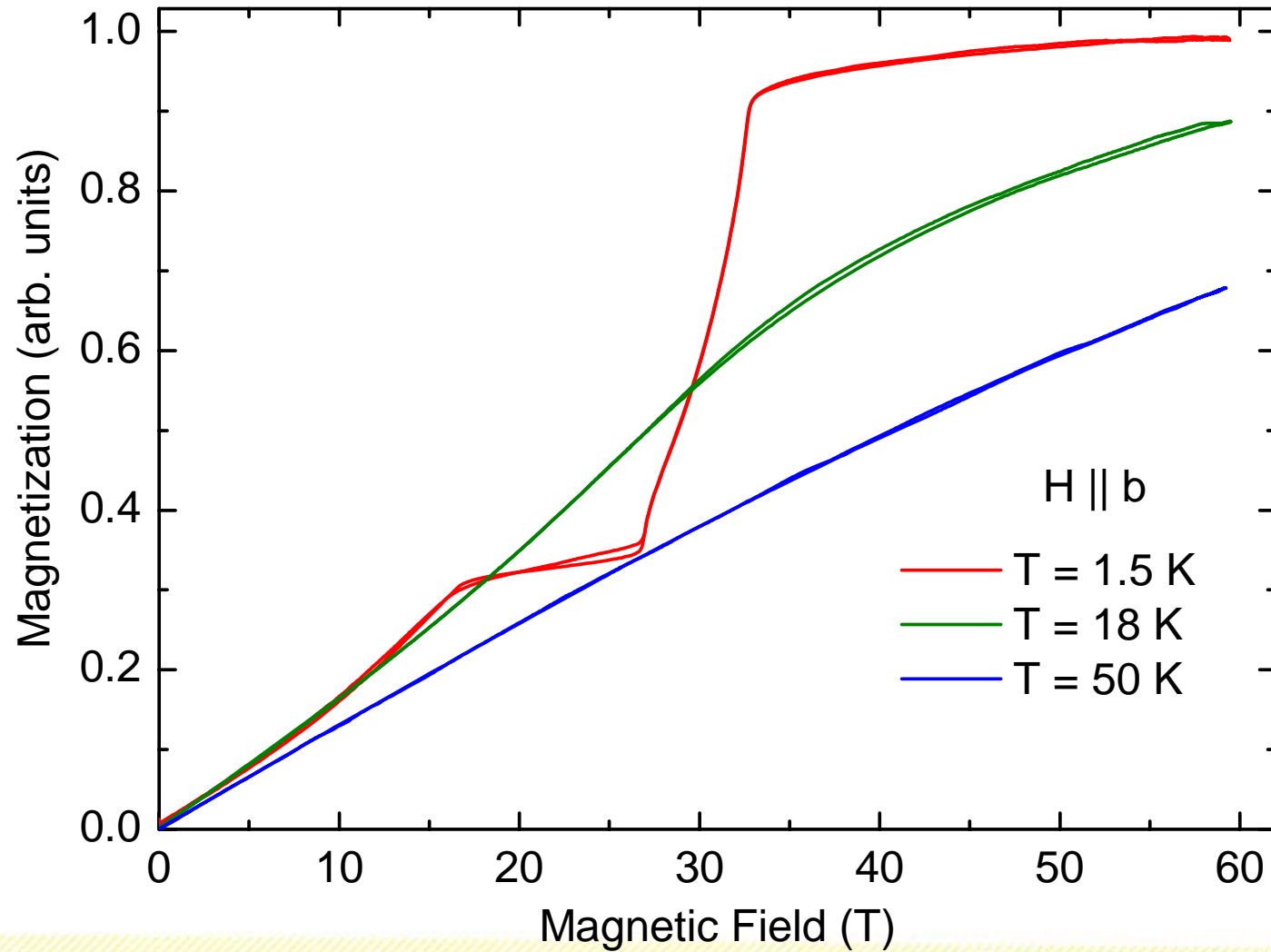


Azurite ($\text{Cu}_3(\text{CO}_3)_2(\text{OH})$)

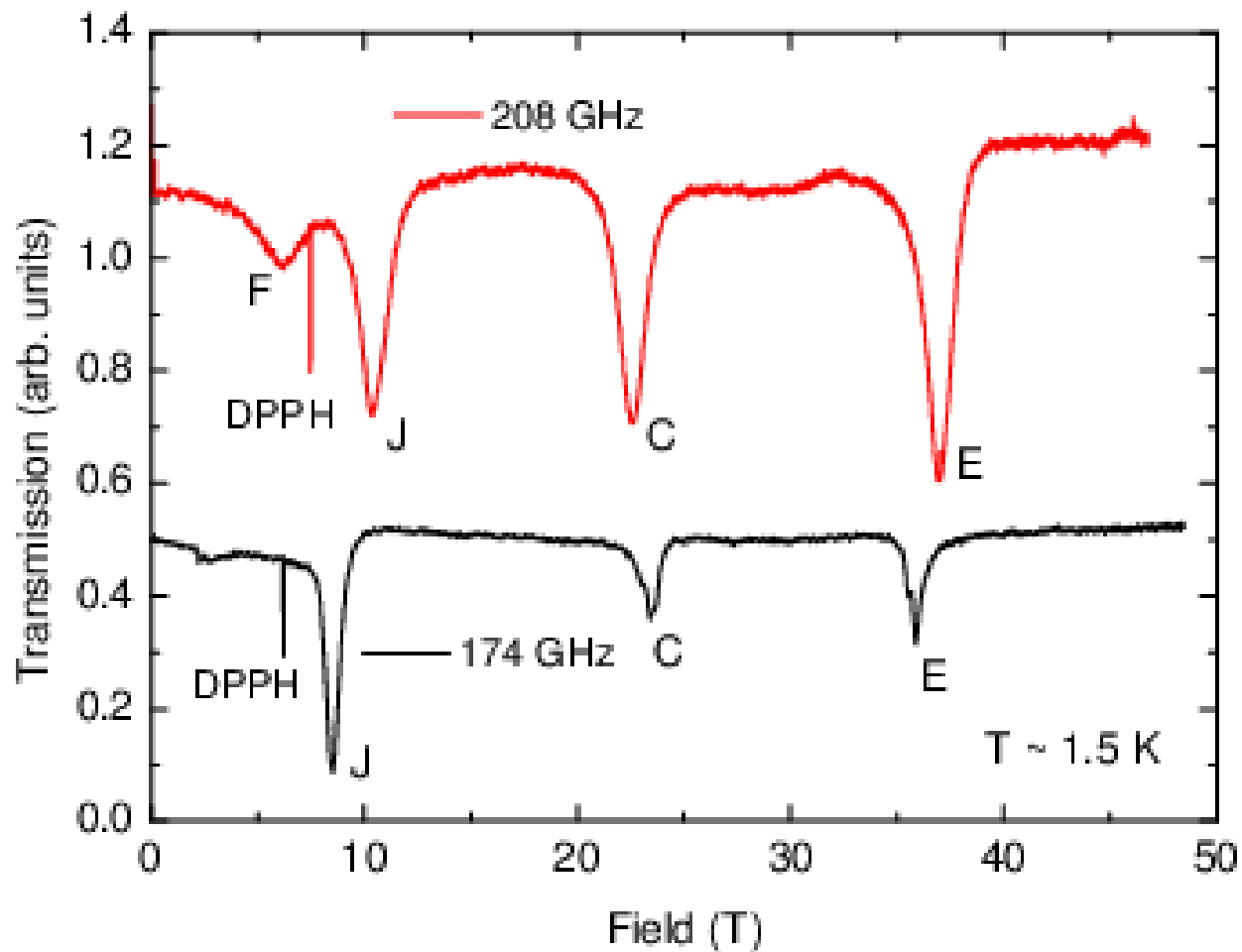


H. Jeschke *et al*, Phys. Rev. Lett 106, 217201 (2011).

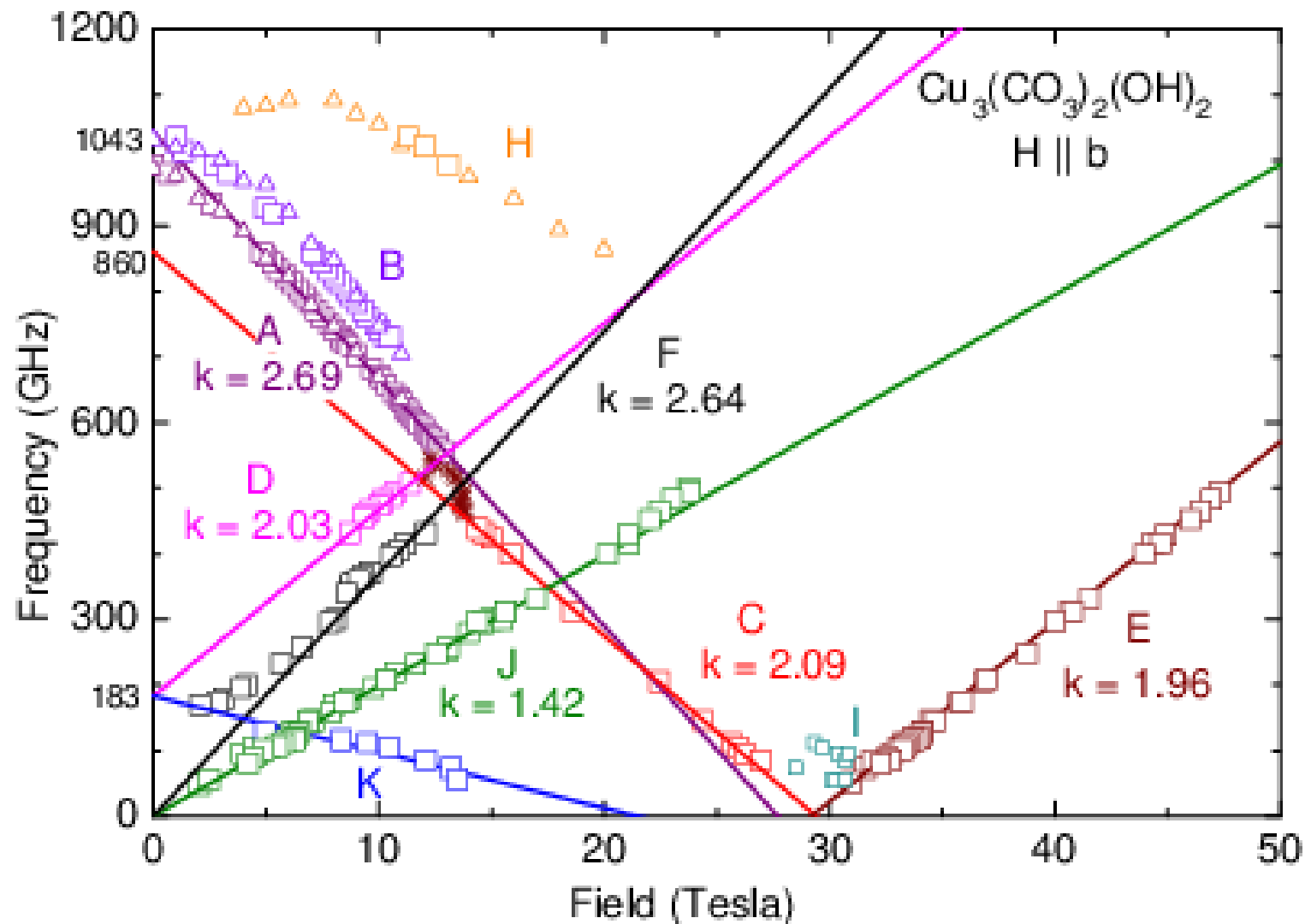
Magnetisation plateau



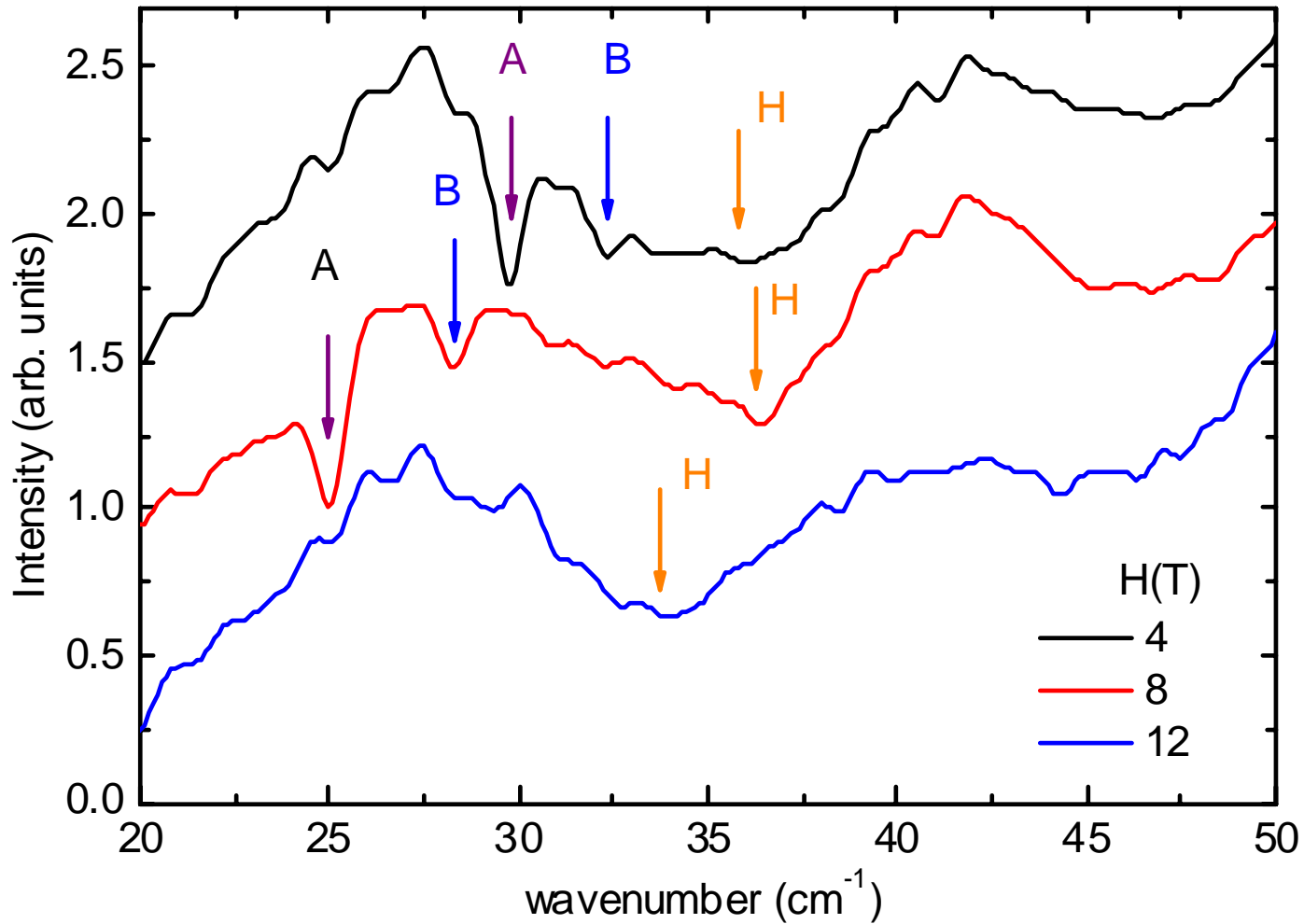
ESR spectra in pulsed fields



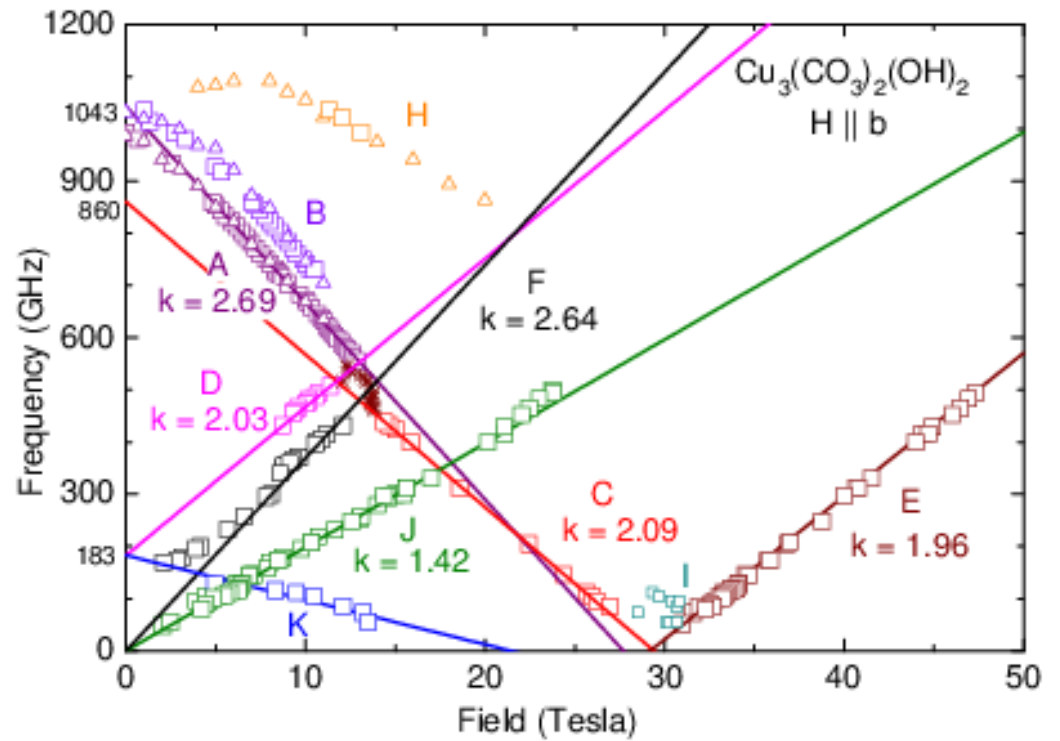
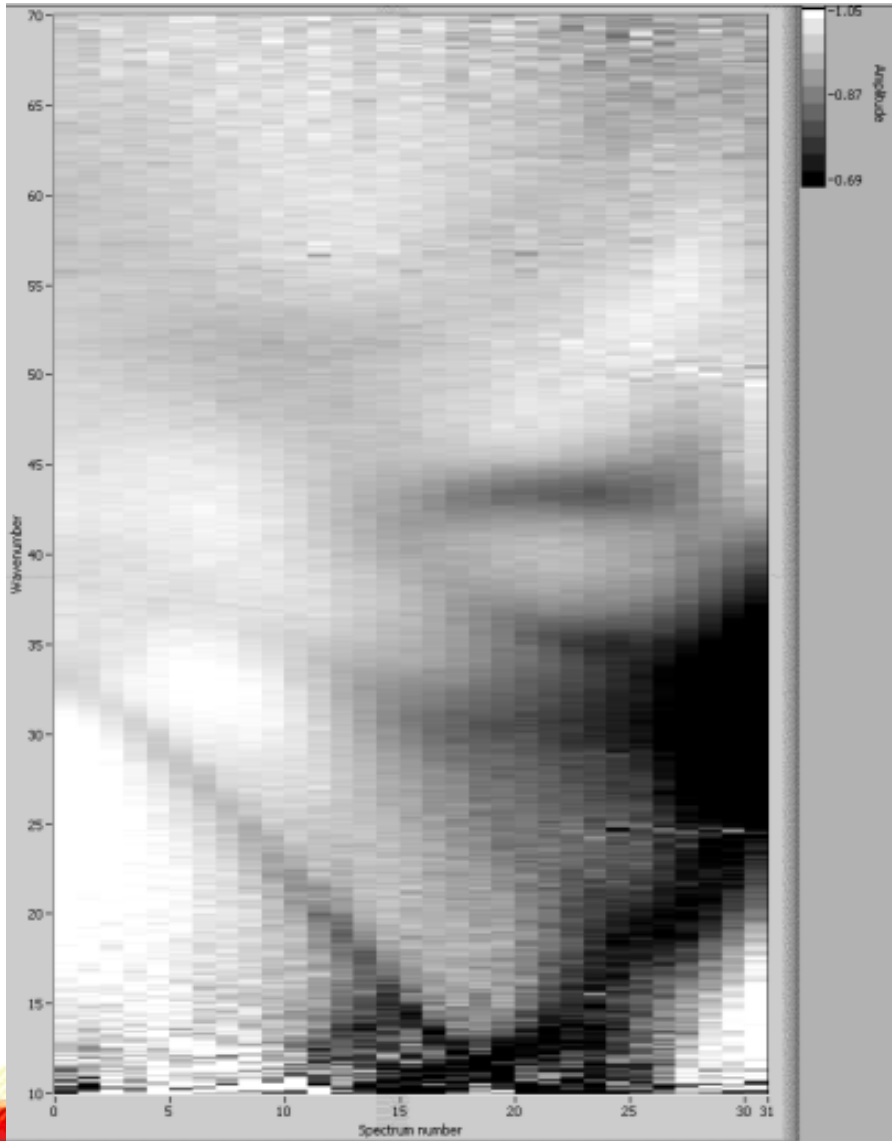
Frequency-field dependence



FIR spectra in magnetic field

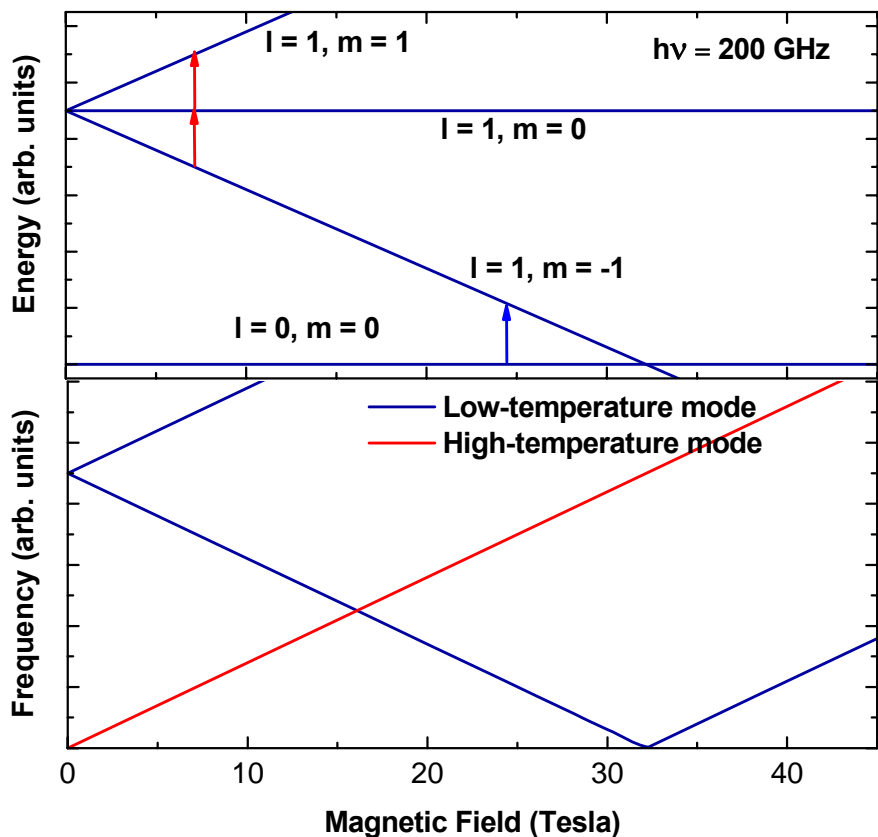


FIR and ESR

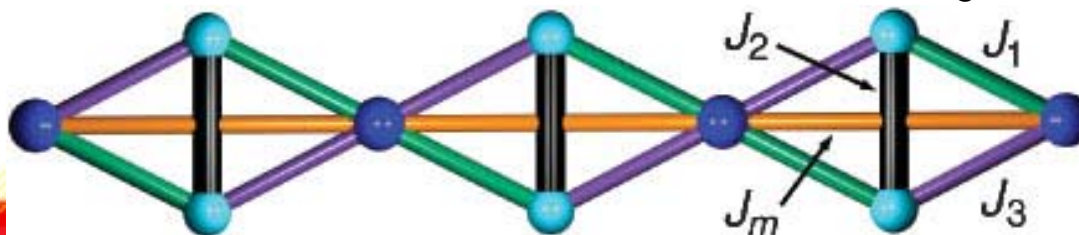
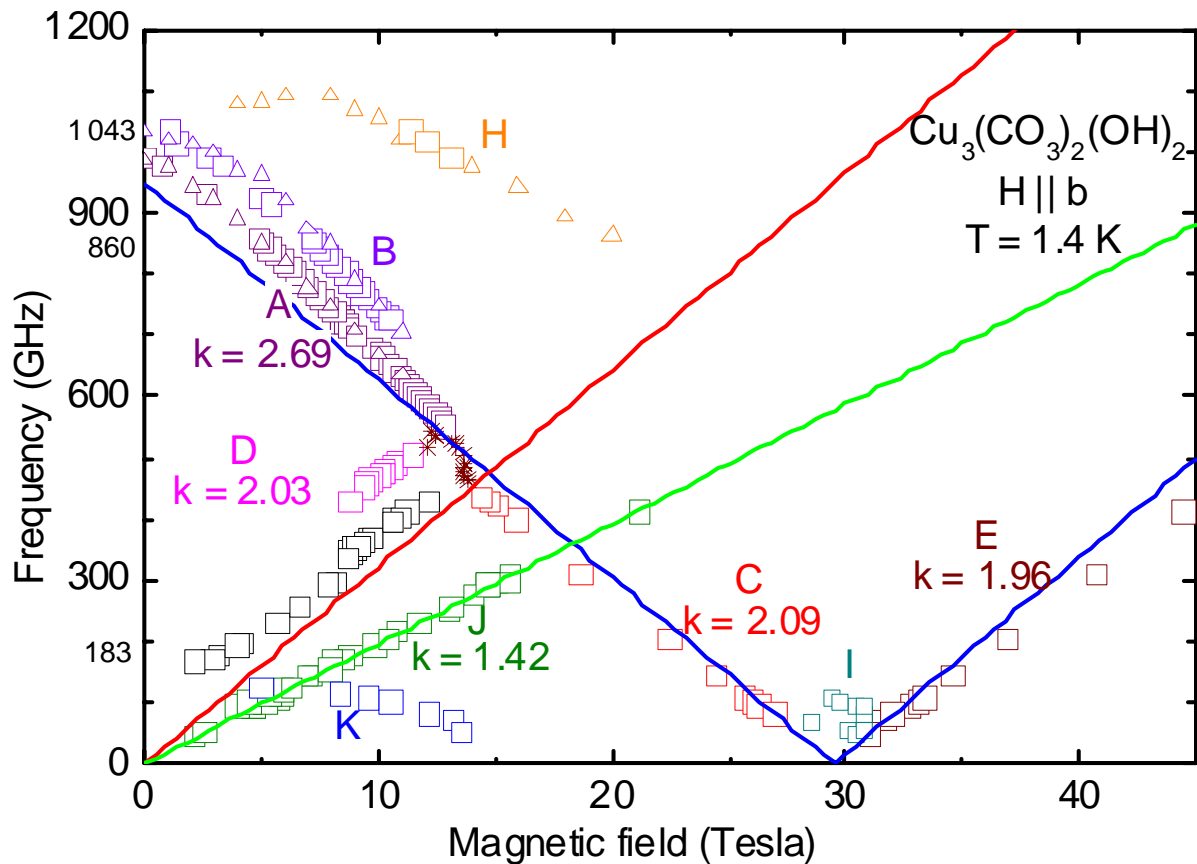


Dimer-monomer model

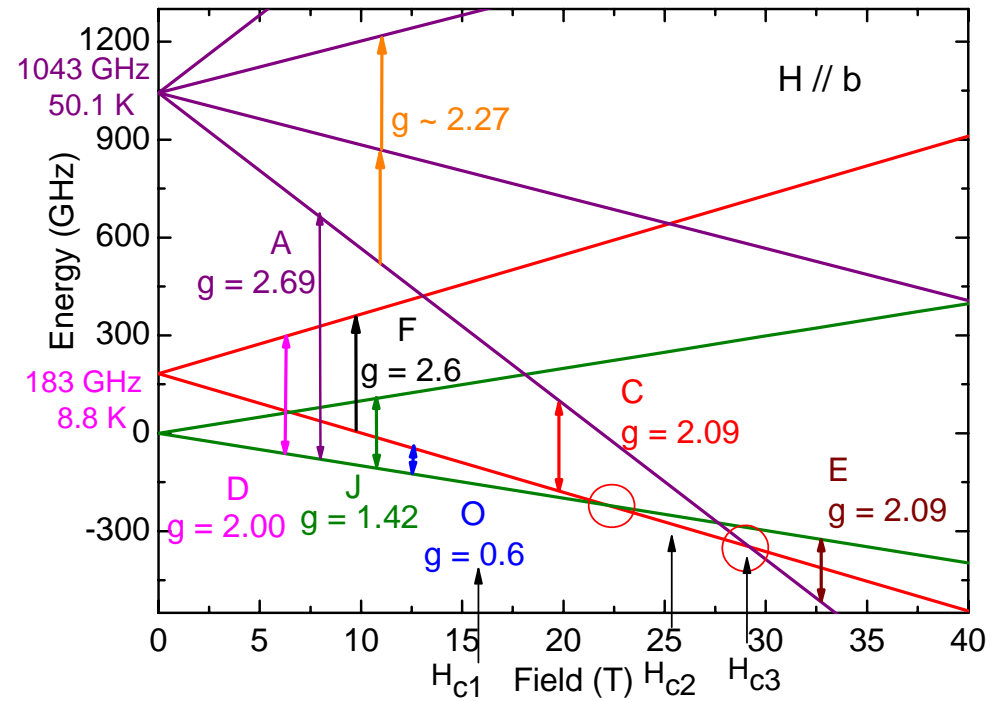
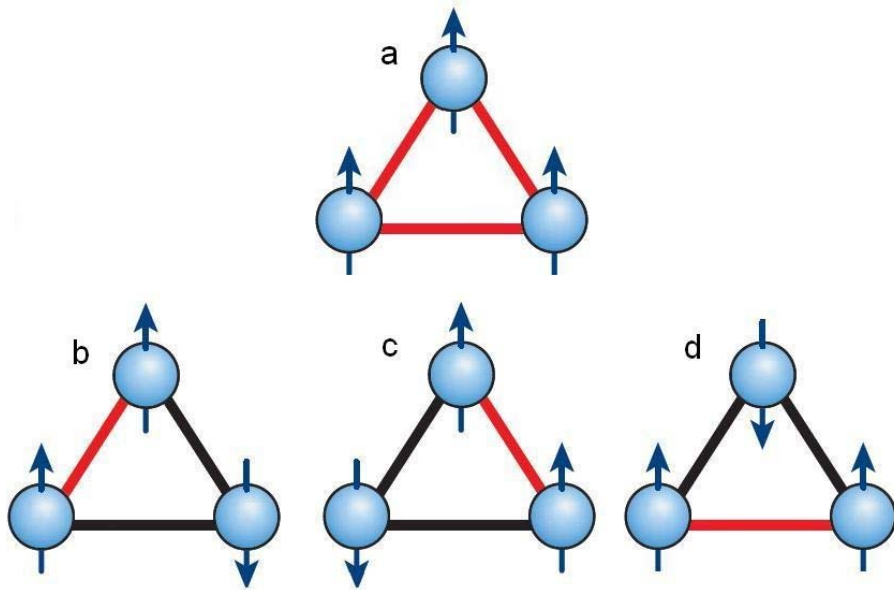
Dimers subsystem



Experimental results

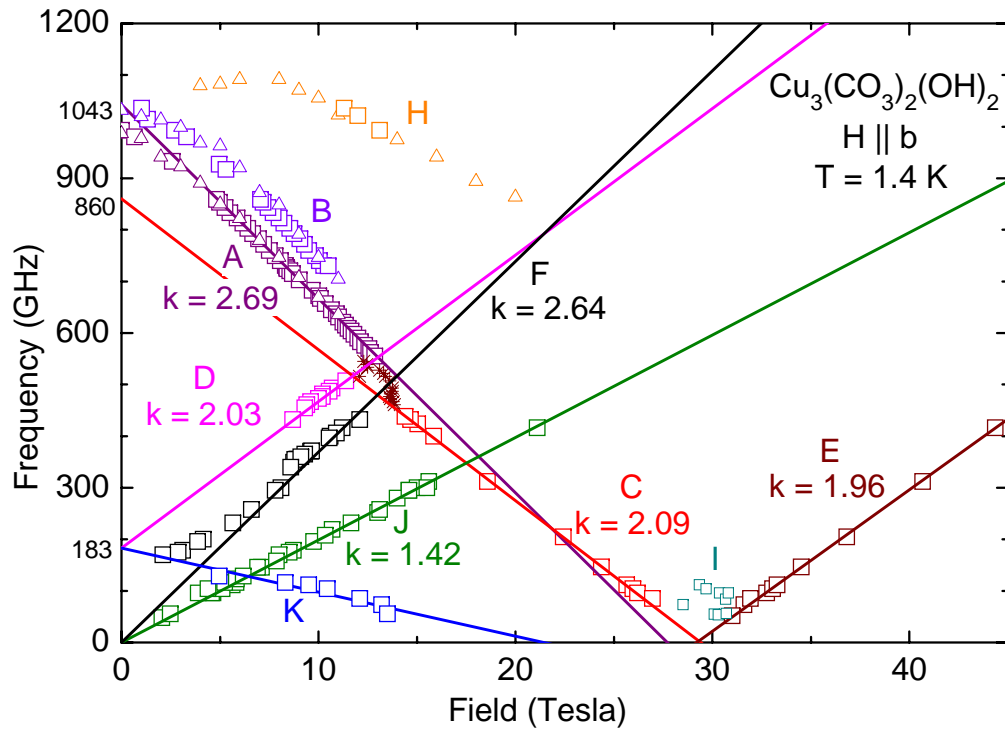


“Triangular” model

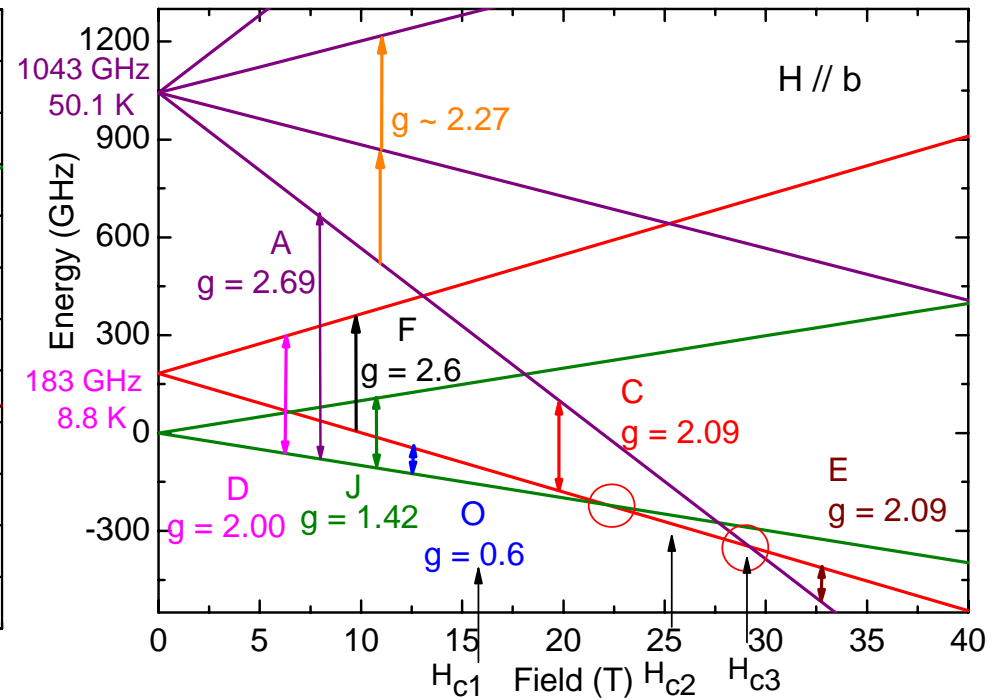


1 level with $S = 3/2$
 3 levels with $S = 1/2$

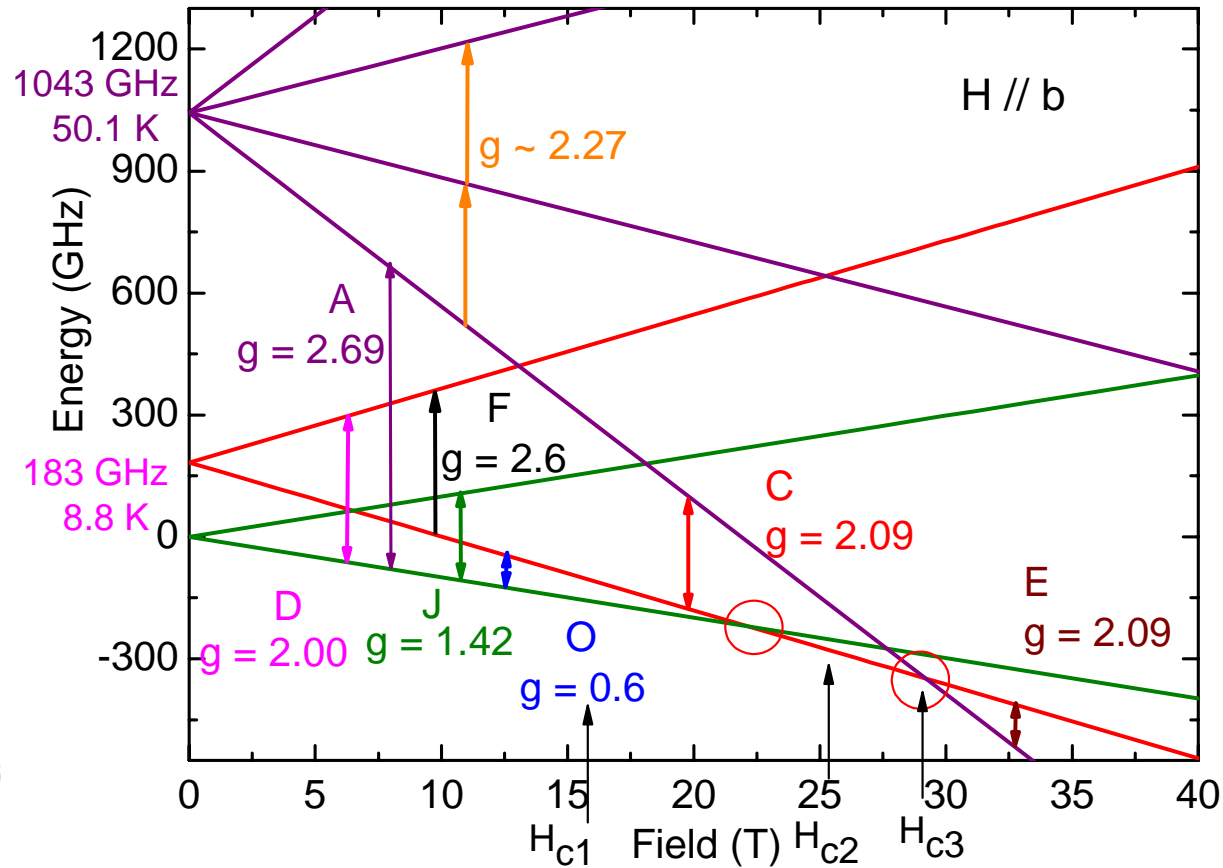
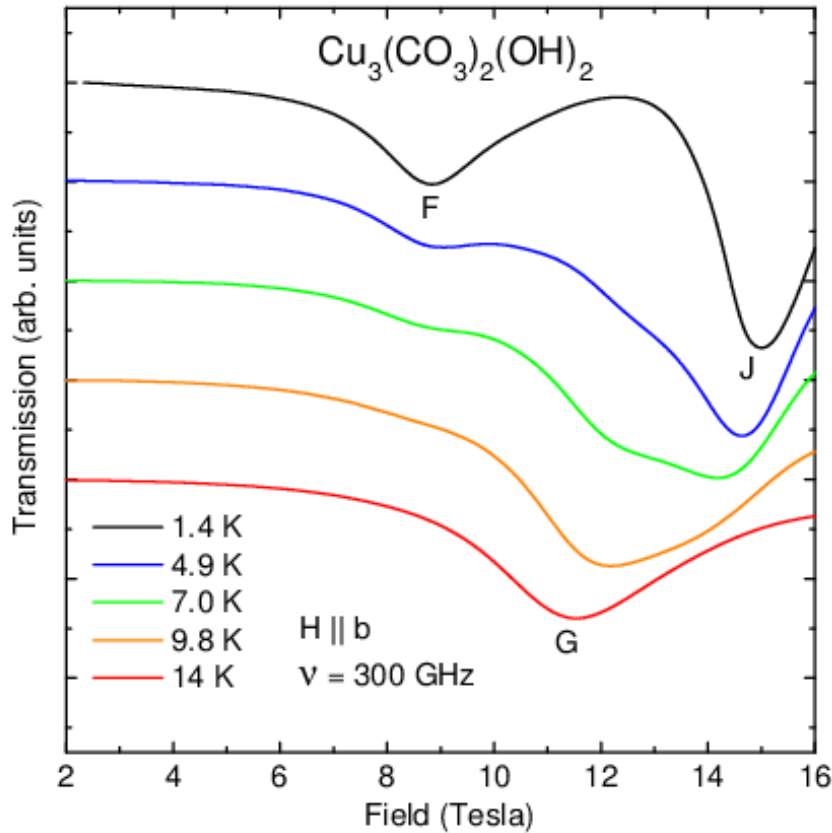
Frequency-field dependance



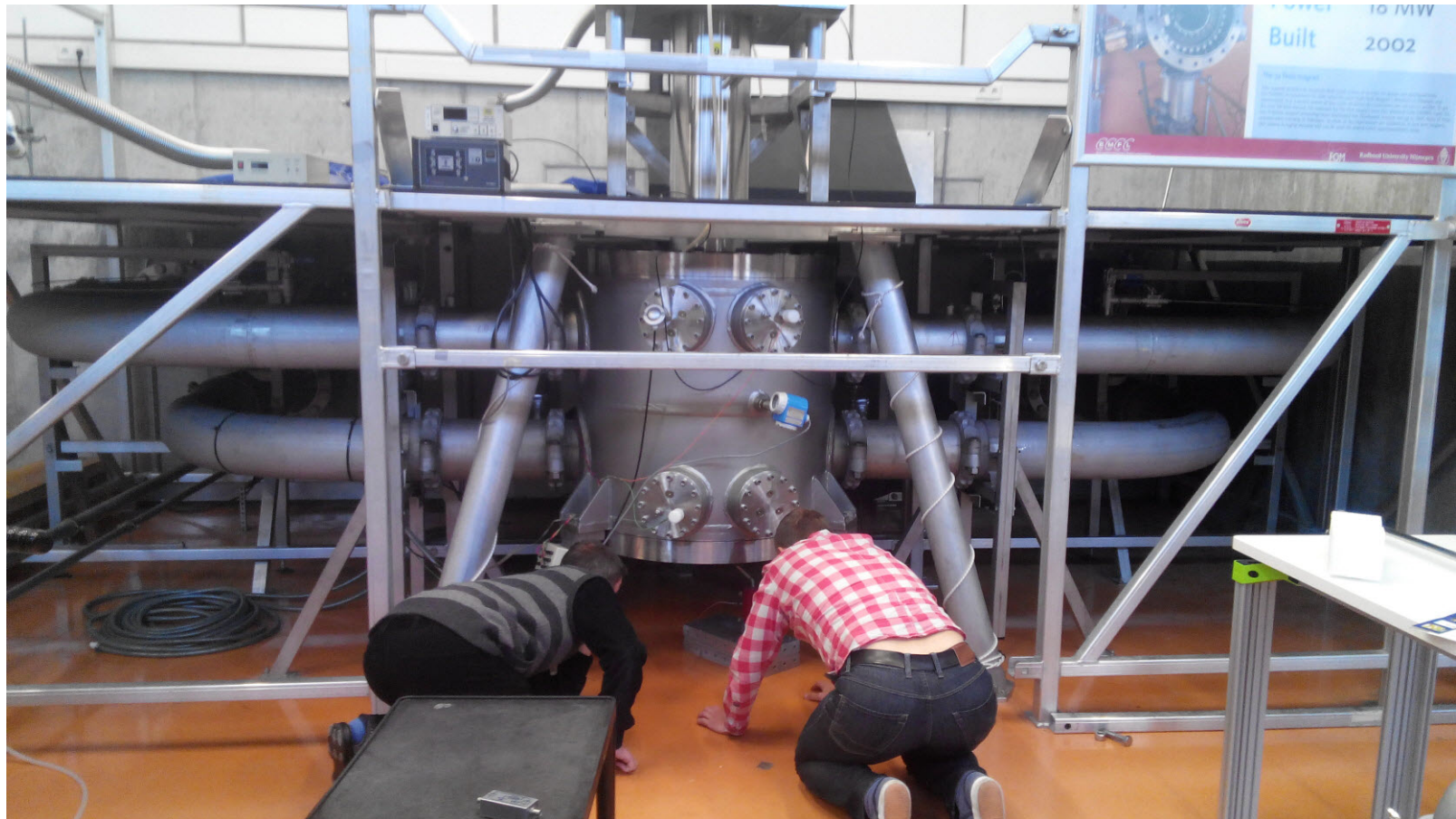
Energy diagram



Temperature evolution of the spectrum



Thank you for your time



Summary

The connection between FELIX and HFML facilities .

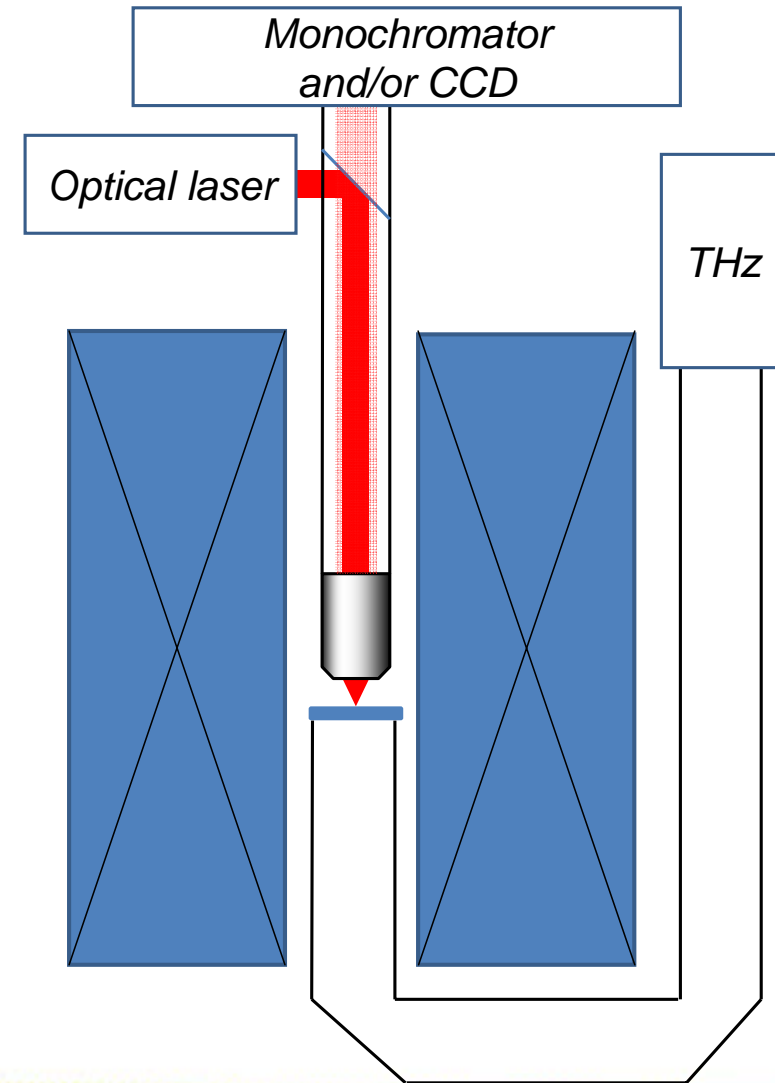
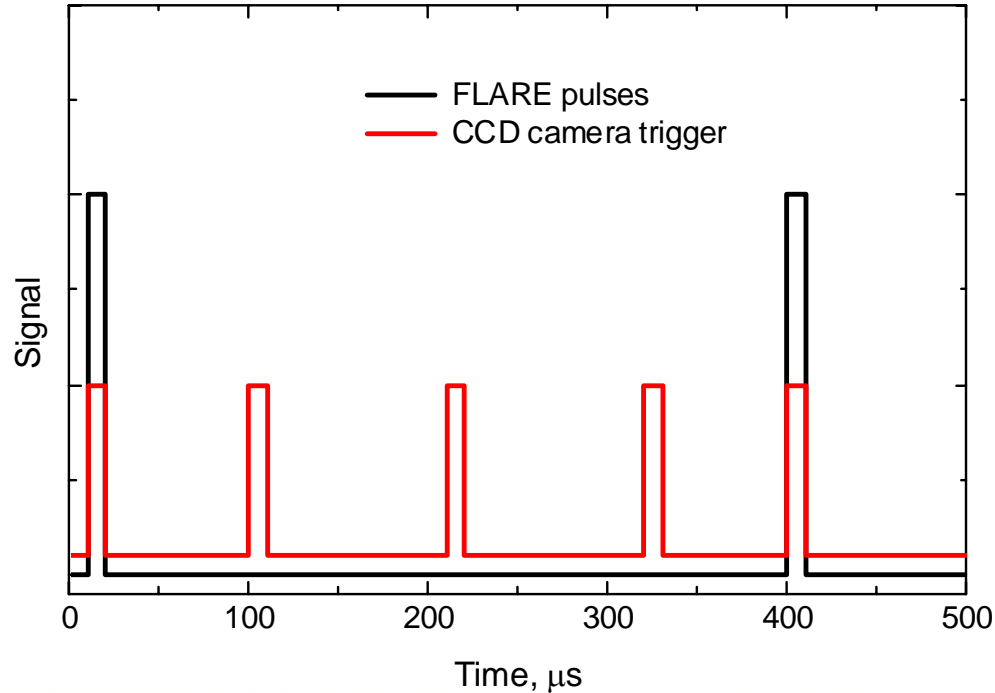
Magneto-transmission setup is working routinely. Magnetic and cyclotron resonance studies are possible over unique ranges of fields, frequencies, radiation powers.

First photoluminescence measurements are successfully performed.

The development of others pump-probe setups in progress

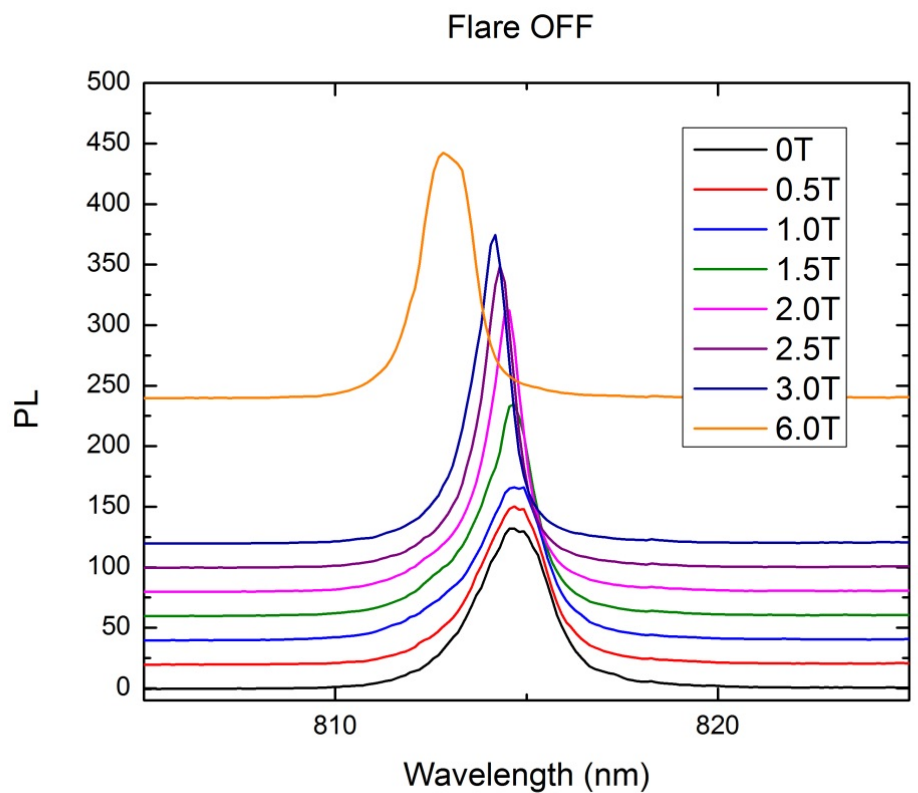
THz pump, PL probe

- Optical photoluminescence spectra as a function of magnetic field and THz radiation (frequency and power).
- THz spectroscopy at optical resolution

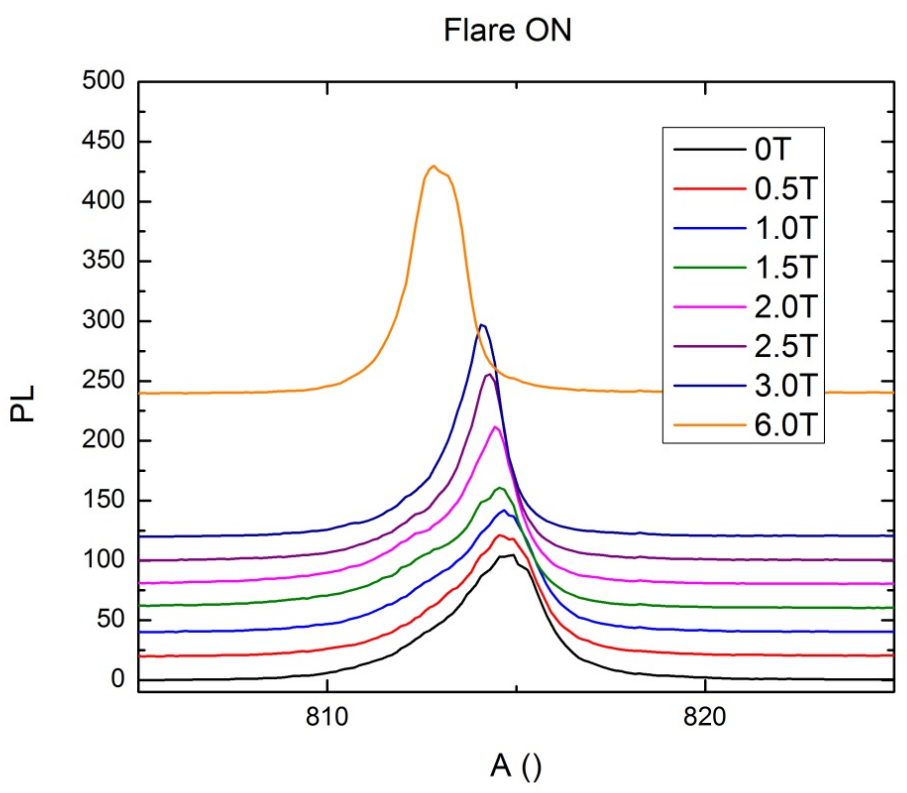


GaAs QW Photoluminescence spectra

PL Spectra

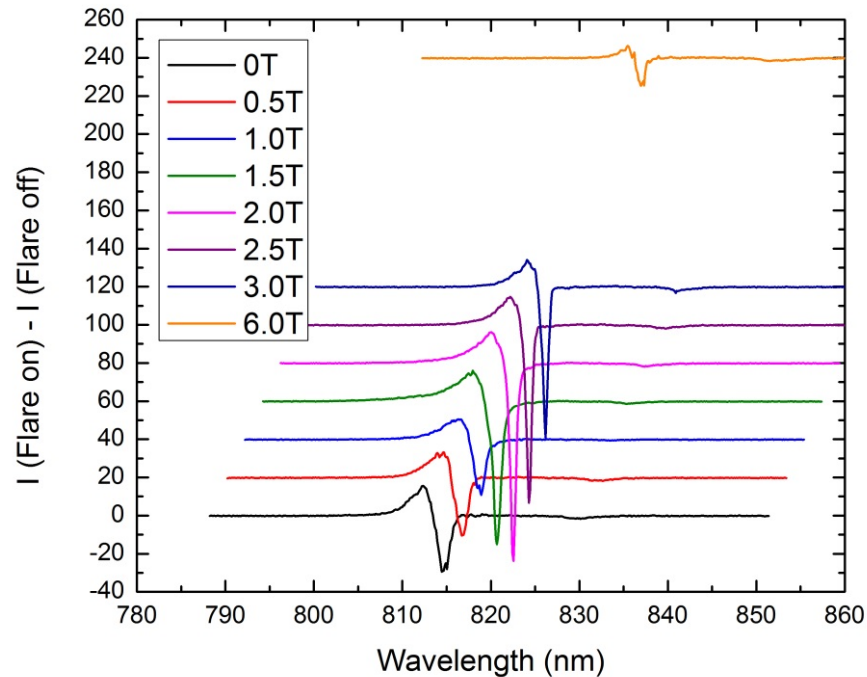


PL Spectra + High THz power (34 cm^{-1})



GaAs QW

Difference between PL Spectra
with and without FLARE



Cyclotron frequency:

$$h\nu = eB/m^*$$

$$m^* = 0.07m_e$$

